

Heavy Ion Physics



Part 2 Raimond Snellings



Content

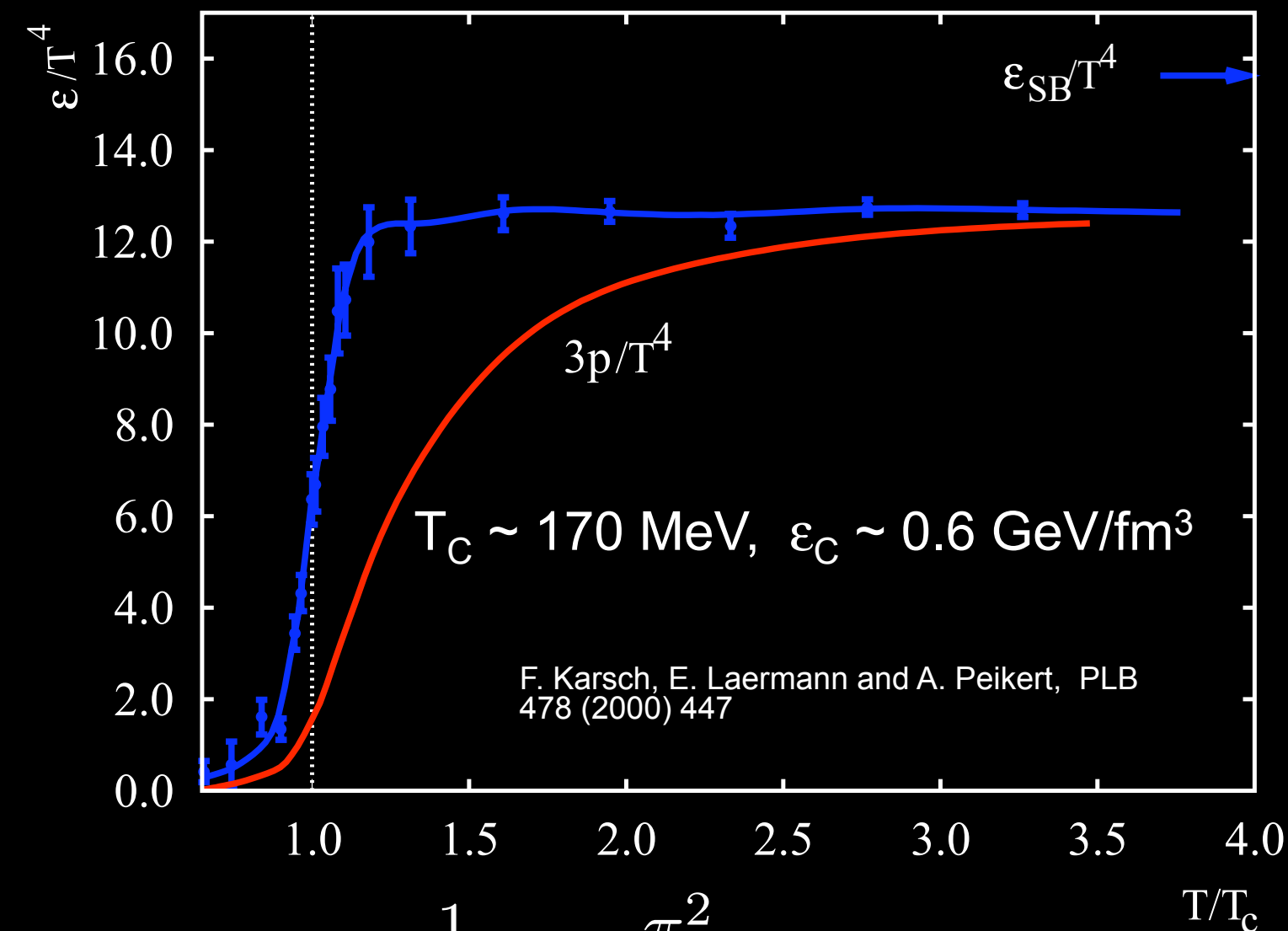
QCD at high density and temperature

heavy-ion accelerators, experiments, global collision
characterization

QGP observables and experimental probes

the LHC heavy-ion program and ALICE the dedicated
heavy-ion detector

QCD on the Lattice



at the critical temperature a strong increase in the degrees of freedom

✓ gluons, quarks & color!

not an ideal gas!

✓ residual interactions

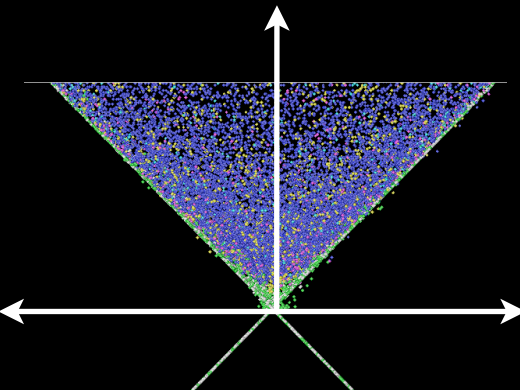
at the phase transition $dp/d\epsilon$ decreases rapidly

$dp/d\epsilon$ drives the collective expansion of the system

$$p = \frac{1}{3}\epsilon = g \frac{\pi^2}{90} T^4$$

$$g_H \approx 3 \quad g_{QGP} \approx 37$$

$$g = 2_{\text{spin}} \times 8_{\text{gluons}} + \frac{7}{8} \times 2_{\text{flavors}} \times 2_{q\bar{q}} \times 2_{\text{spin}} \times 3_{\text{color}}$$

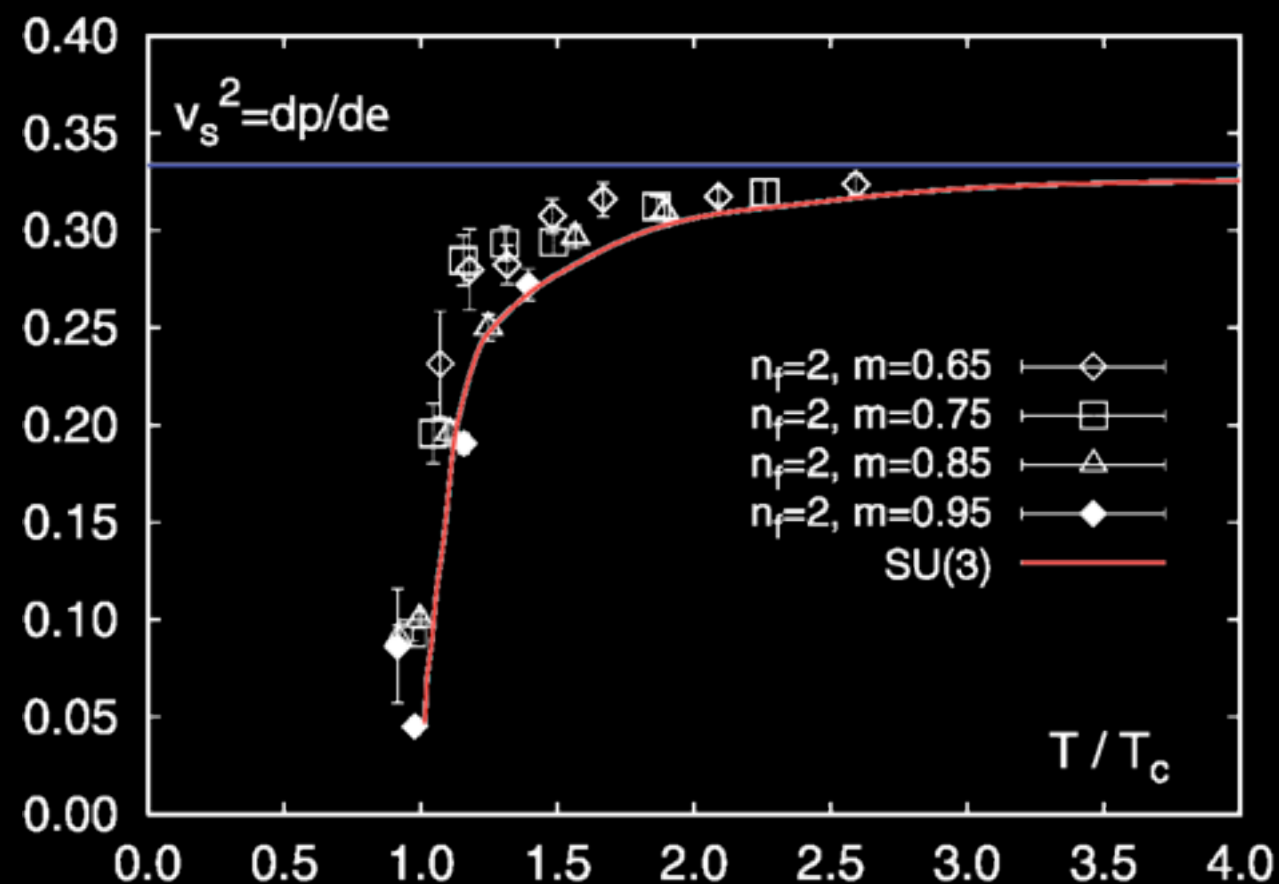
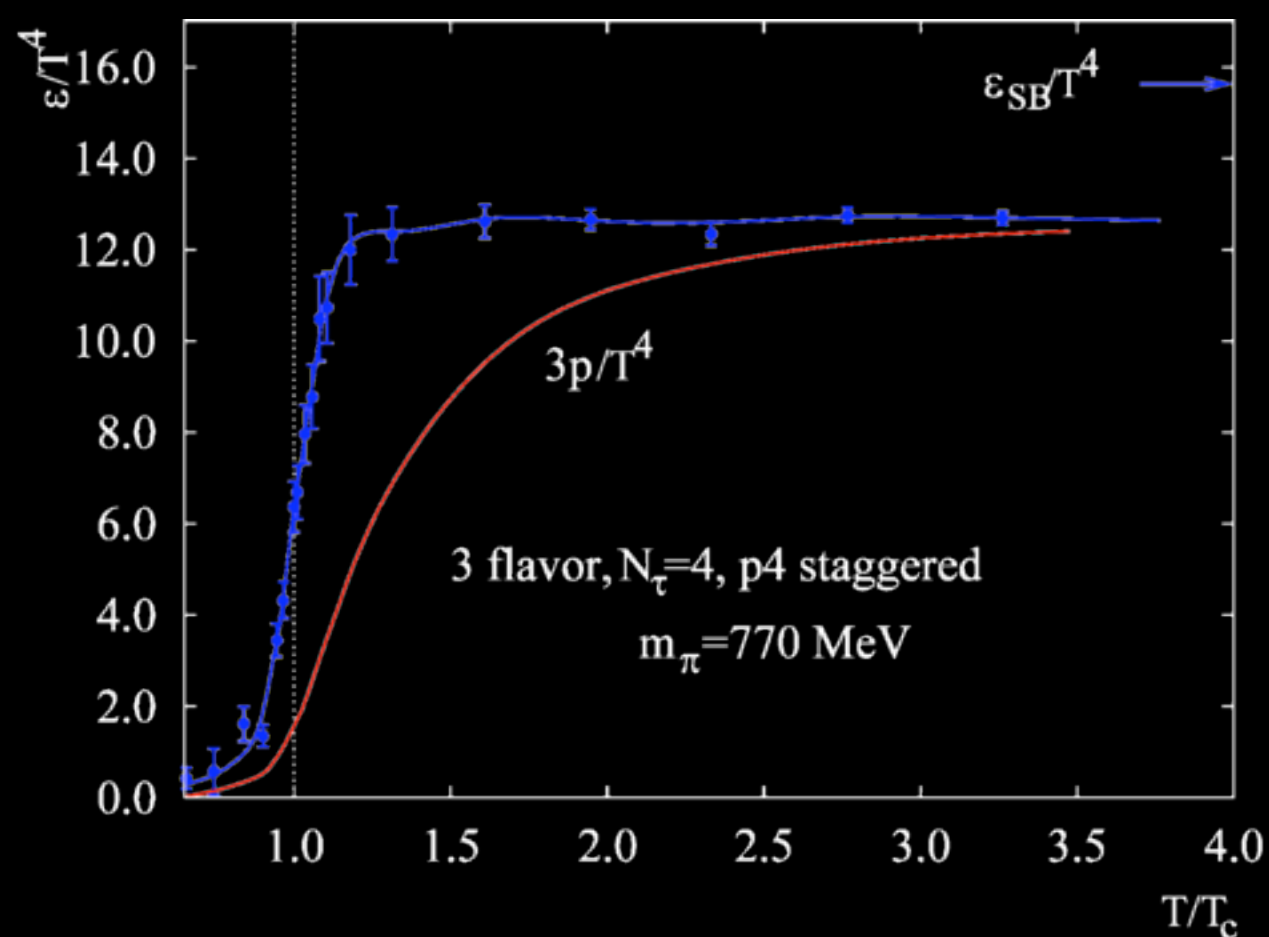


Velocity of Sound

F. Karsch and E. Laermann, arXiv:hep-lat/0305025

$$P_{\text{QGP}} = \frac{1}{3} \epsilon_{\text{QGP}} = g \frac{\pi^2}{90} T^4$$

$$\text{velocity of sound } C_s = \sqrt{\frac{dP}{d\epsilon}}$$



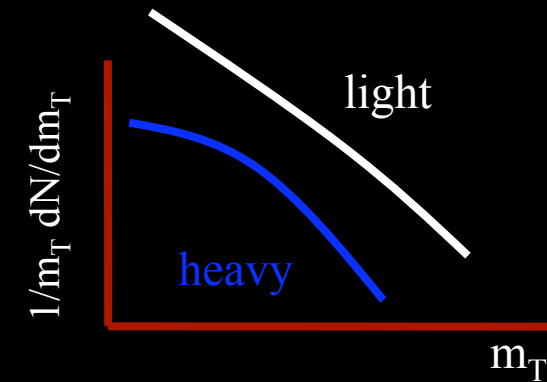
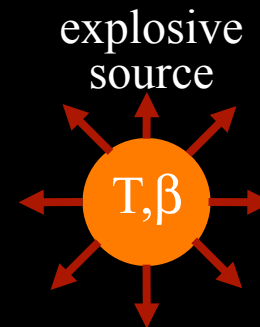
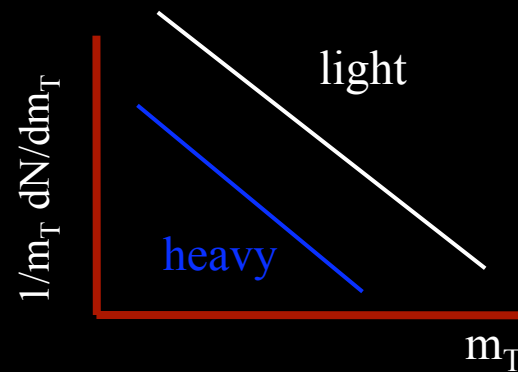
the magnitude of the collective motion is proportional to the velocity of sound

Collective Motion

$$m_T = \sqrt{(m^2 + p_t^2)}$$

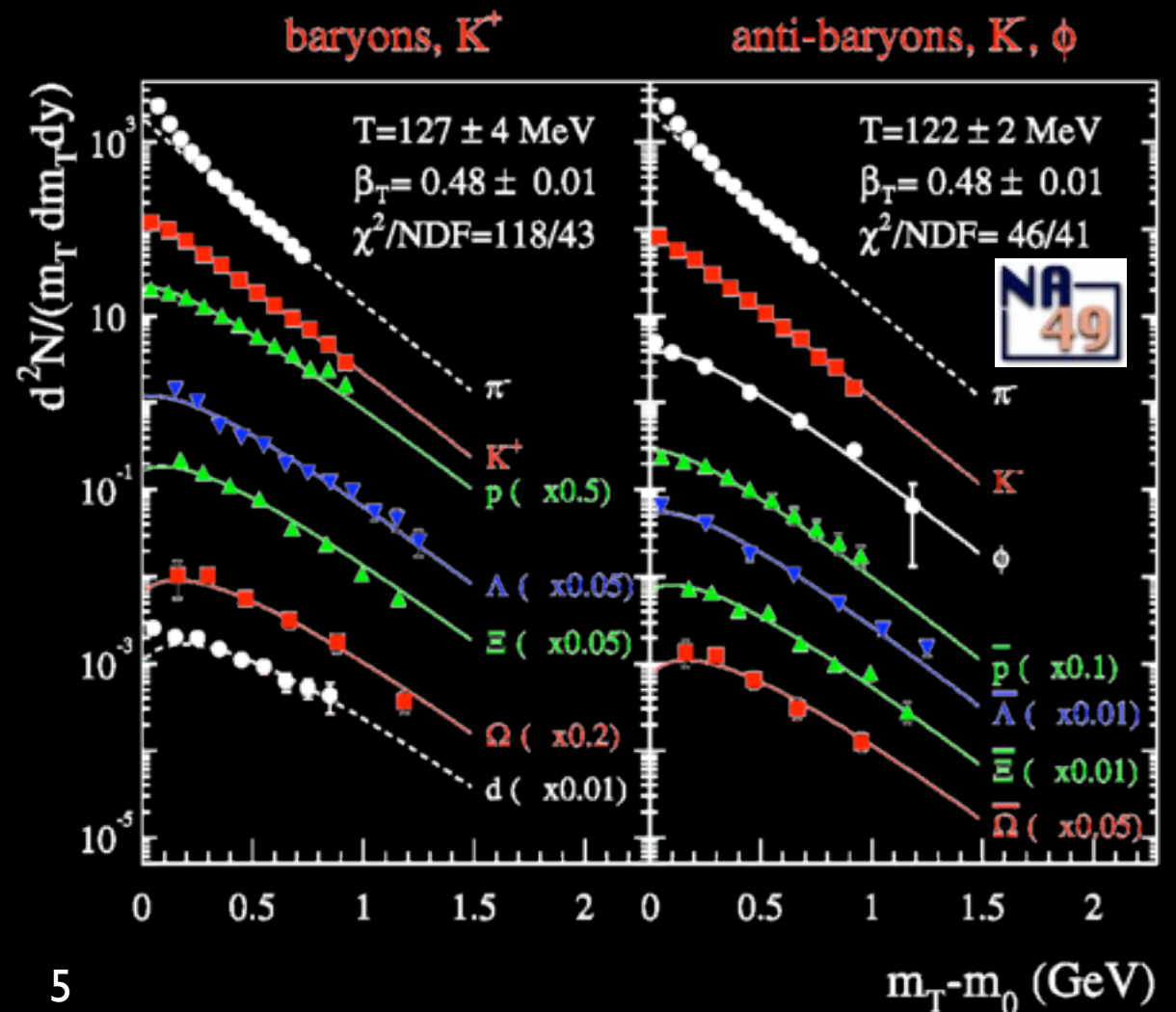
$$\frac{dN}{m_T dm_T} \propto e^{-m_T/T}$$

purely thermal
source



in p-p at low transverse momenta the particle yields are well described by thermal spectra (m_T scaling)

boosted thermal spectra give a very good description of the particle distributions measured in heavy-ion collisions

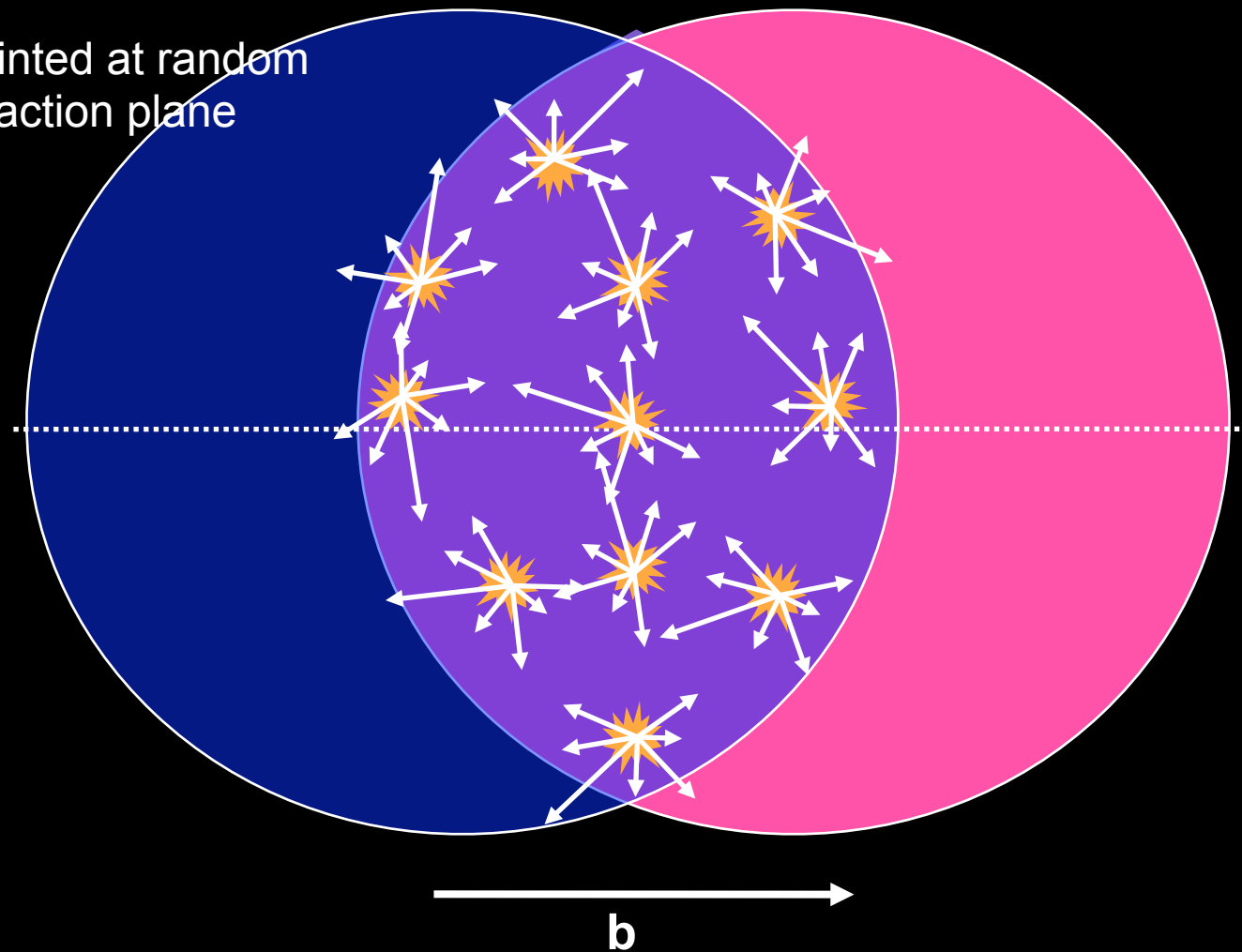


Elliptic Flow

Animation: Mike Lisa

1) superposition of independent p+p:

momenta pointed at random
relative to reaction plane

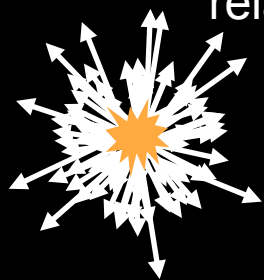


$$\varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

Elliptic Flow

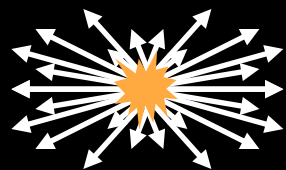
1) superposition of independent p+p:

momenta pointed at random
relative to reaction plane

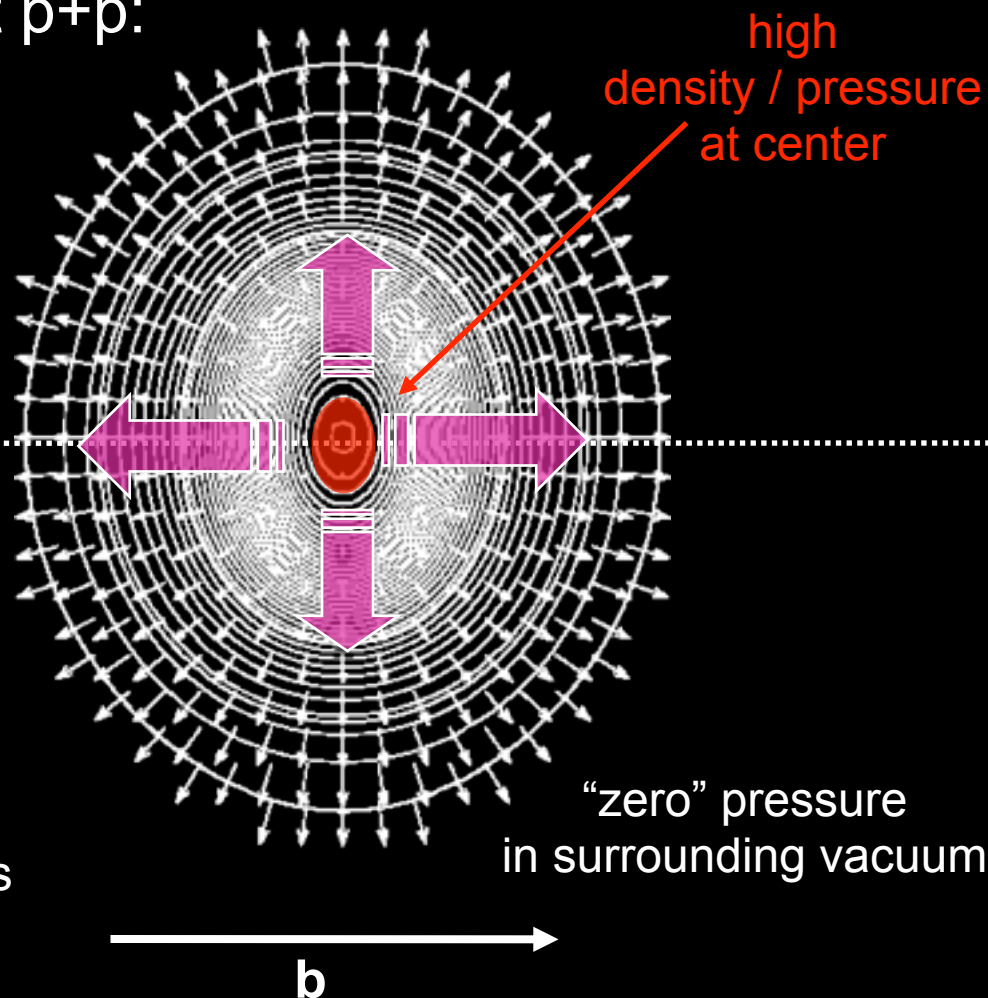


2) evolution as a **bulk system**

pressure gradients (larger in-plane)
push bulk “out” → “flow”



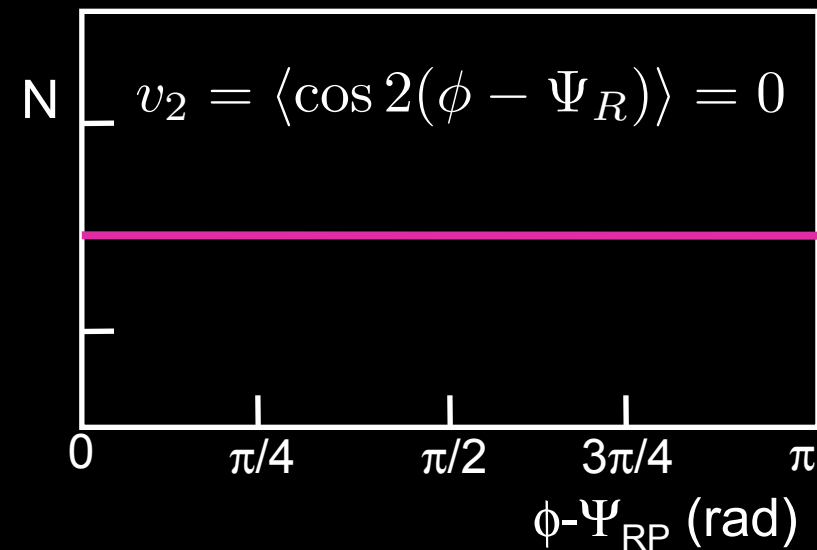
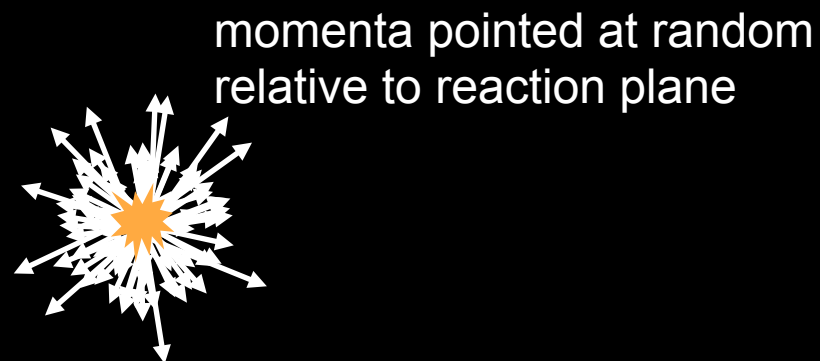
more, faster particles
seen in-plane



$$\varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

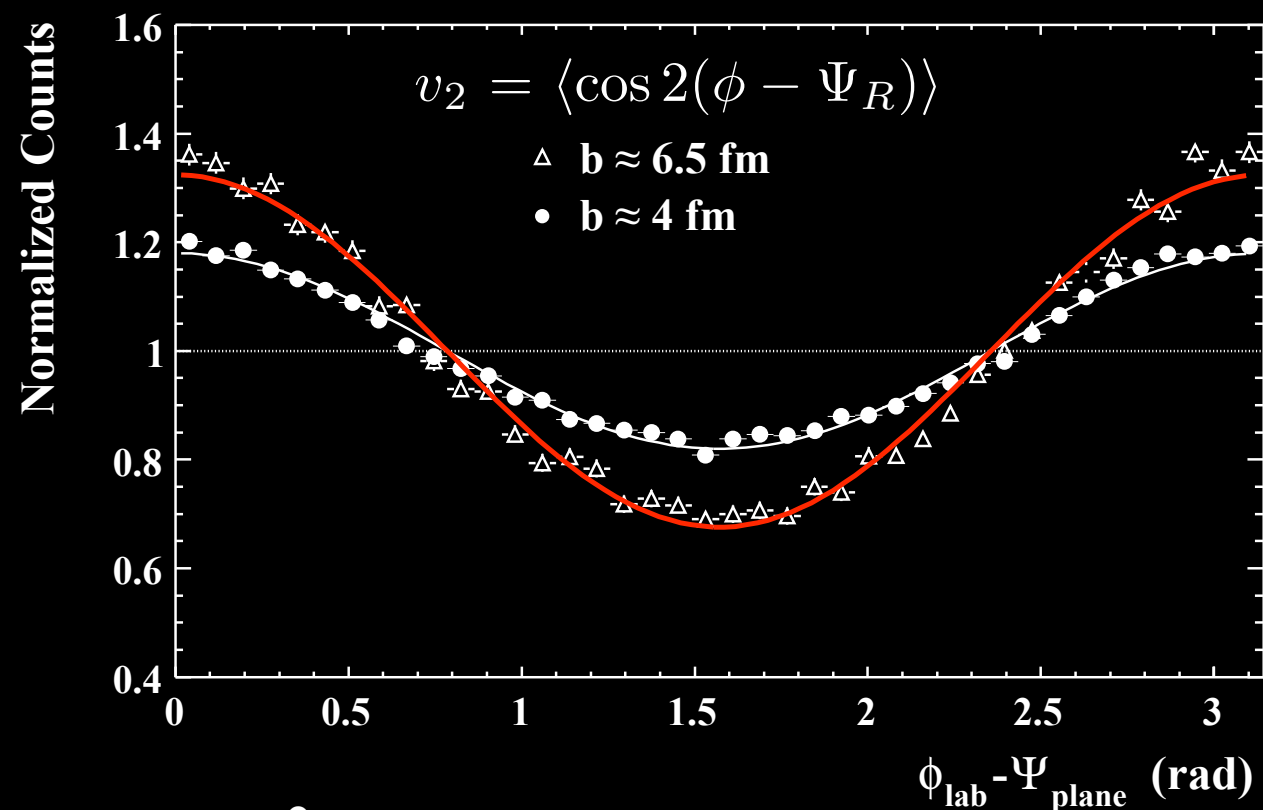
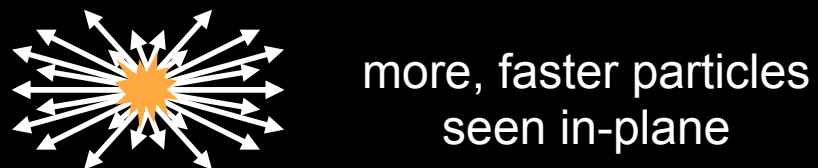
Elliptic Flow

1) superposition of independent p+p:



2) evolution as a **bulk system**

pressure gradients (larger in-plane)
push bulk “out” → “flow”

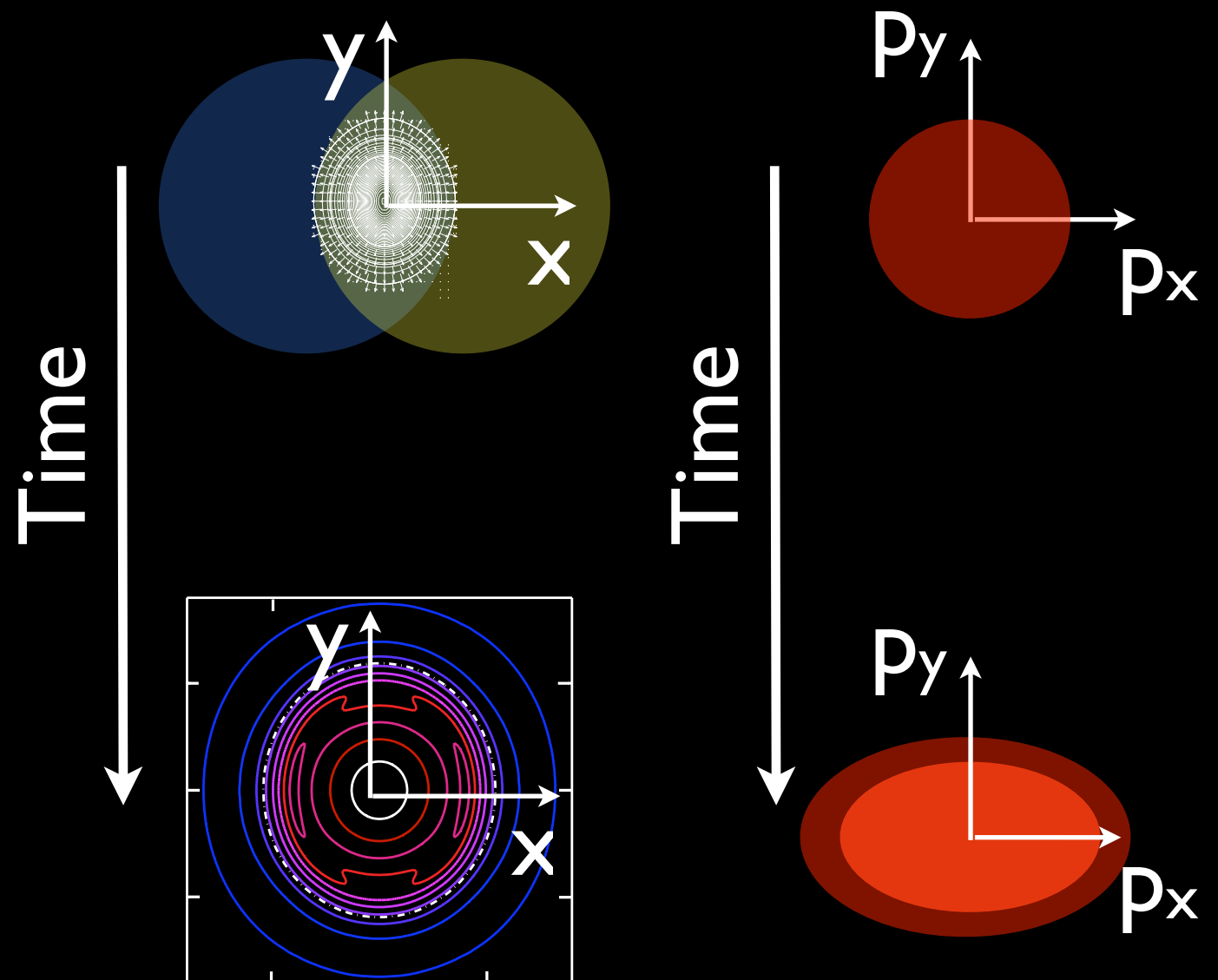


Elliptic Flow

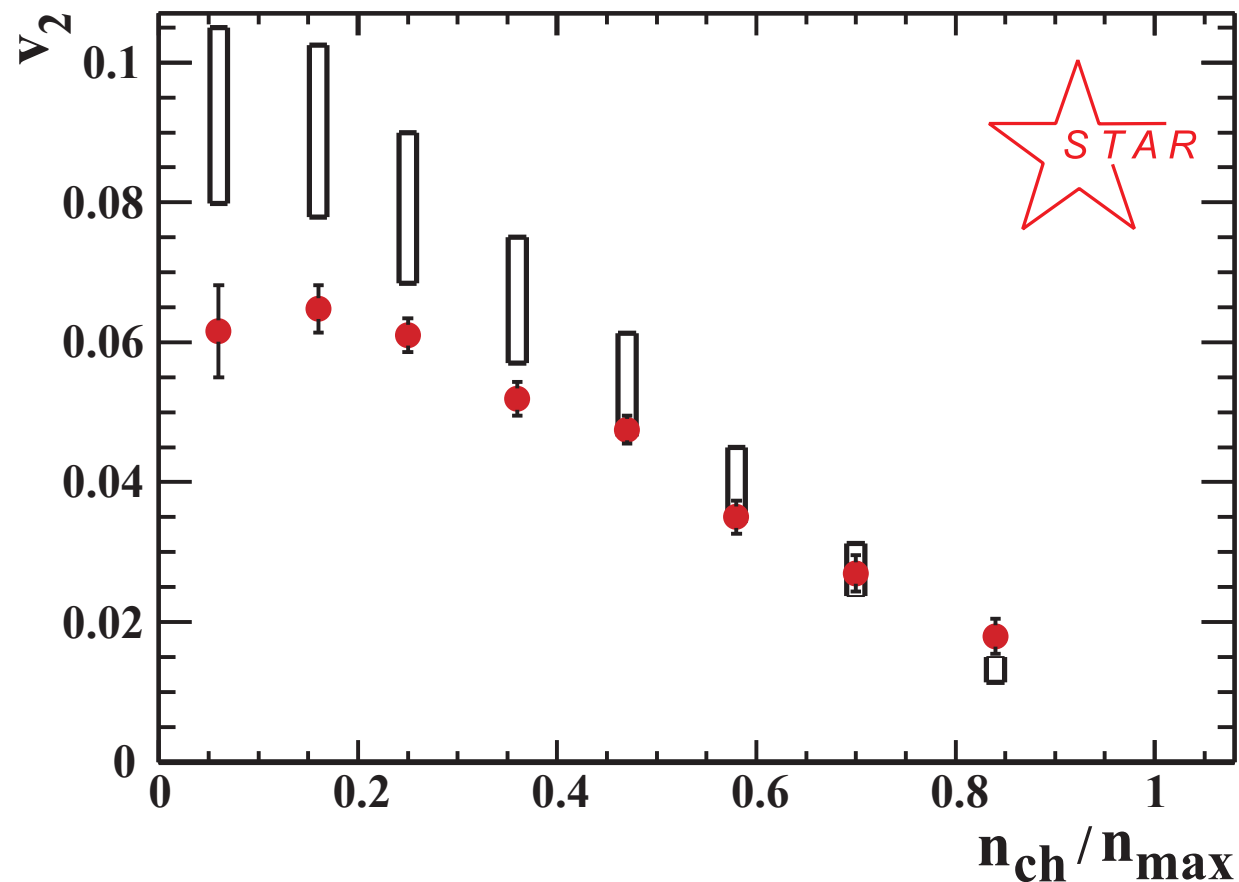
- in non central collisions coordinate space configuration is anisotropic (almond shape). However, initial momentum distribution isotropic (spherically symmetric)
 - interactions among constituents generate a pressure gradient which transforms the initial coordinate space anisotropy into the observed momentum space anisotropy → anisotropic flow
 - self-quenching → sensitive to early stage
- a unique hadronic probe of the early stage

$$\varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

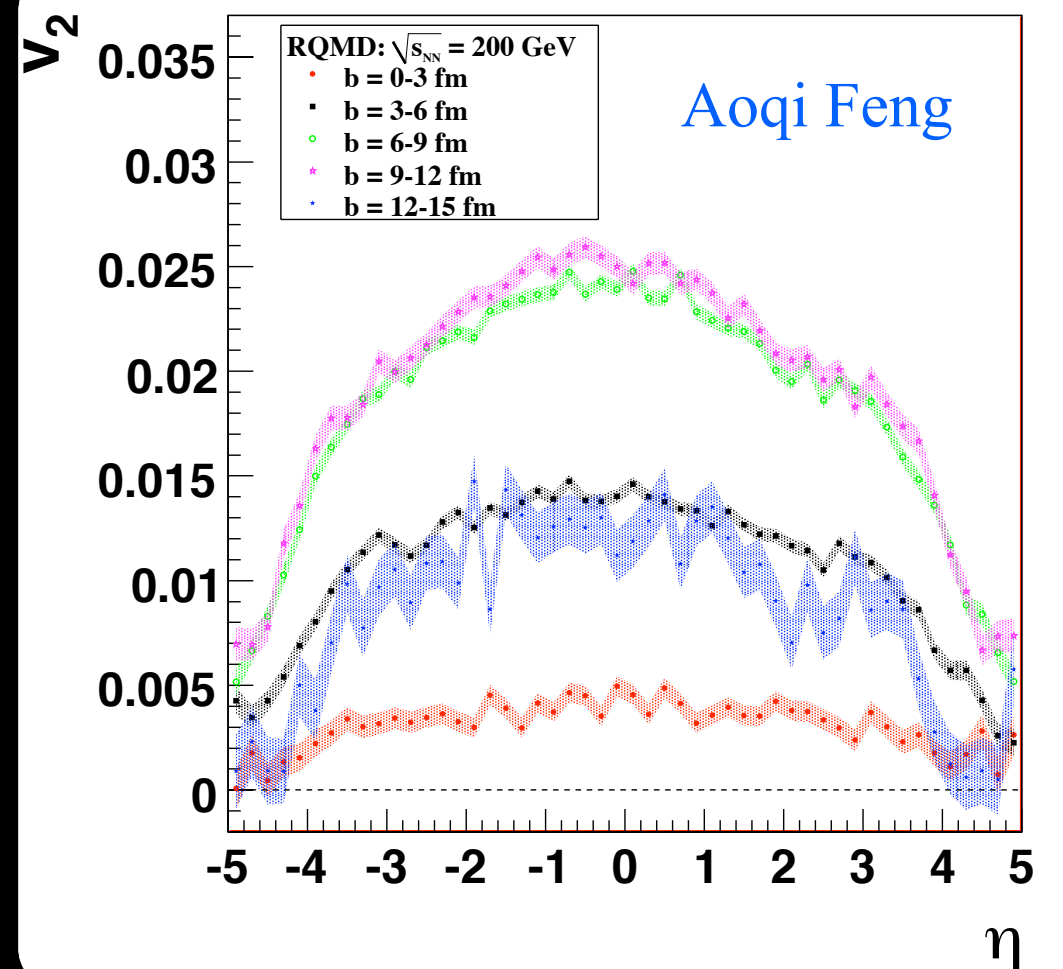
$$v_2 = \langle \cos 2\phi \rangle$$



Flow at RHIC



STAR Phys. Rev. Lett. 86, 402–407 (2001)



ideal hydro gets the magnitude for more central collisions
hadron transport calculations are factors 2-3 off

Hydro Motivated Fit

$$v_2(p_t) = \frac{\int_0^{2\pi} d\phi_b \cos(2\phi_b) I_2(\alpha_t) K_1(\beta_t) (1 + 2s_2 \cos(2\phi_b))}{\int_0^{2\pi} d\phi_b I_0(\alpha_t) K_1(\beta_t) (1 + 2s_2 \cos(2\phi_b))}$$

$$\alpha_t(\phi_b) = \left(\frac{p_t}{T_f}\right) \sinh(\rho(\phi_b)) \quad \beta_t(\phi_b) = \left(\frac{m_t}{T_f}\right) \cosh(\rho(\phi_b))$$

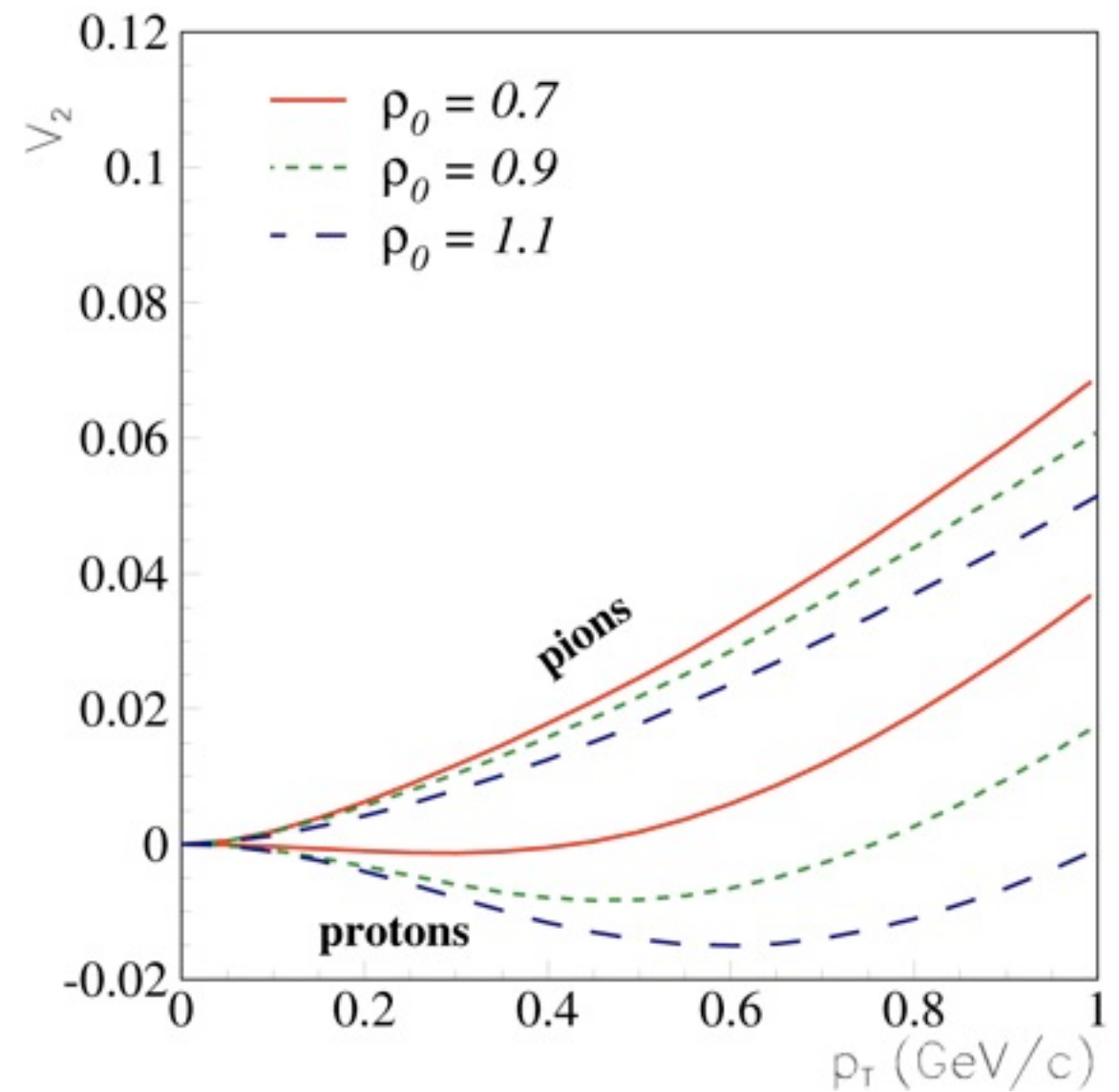
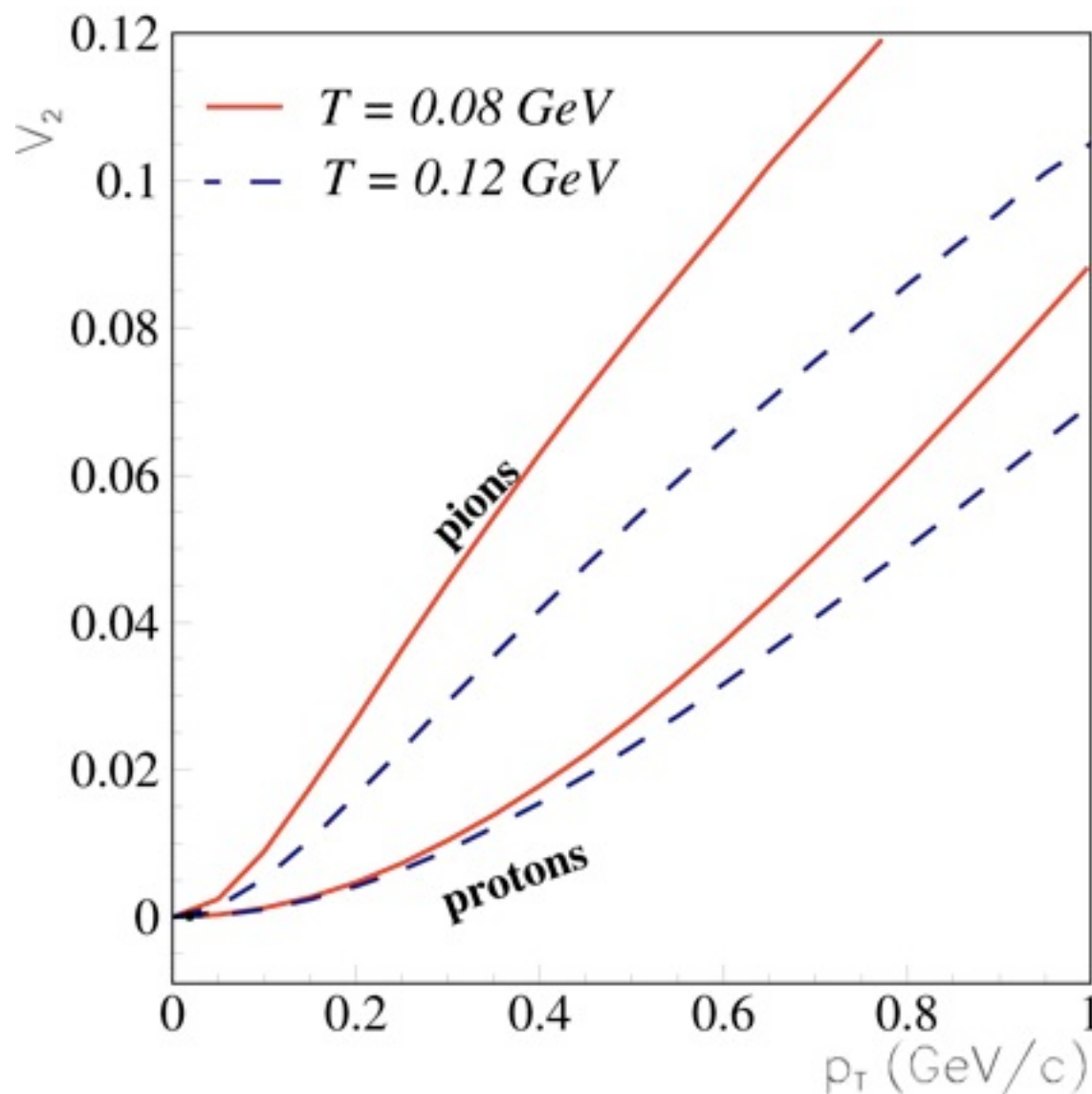
$$\rho(\phi_b) = \rho_0 + \rho_a \cos(2\phi_b)$$

STAR Phys. Rev. Lett. 87, 182301 (2001)

$v_2(p_t)$ and particle mass

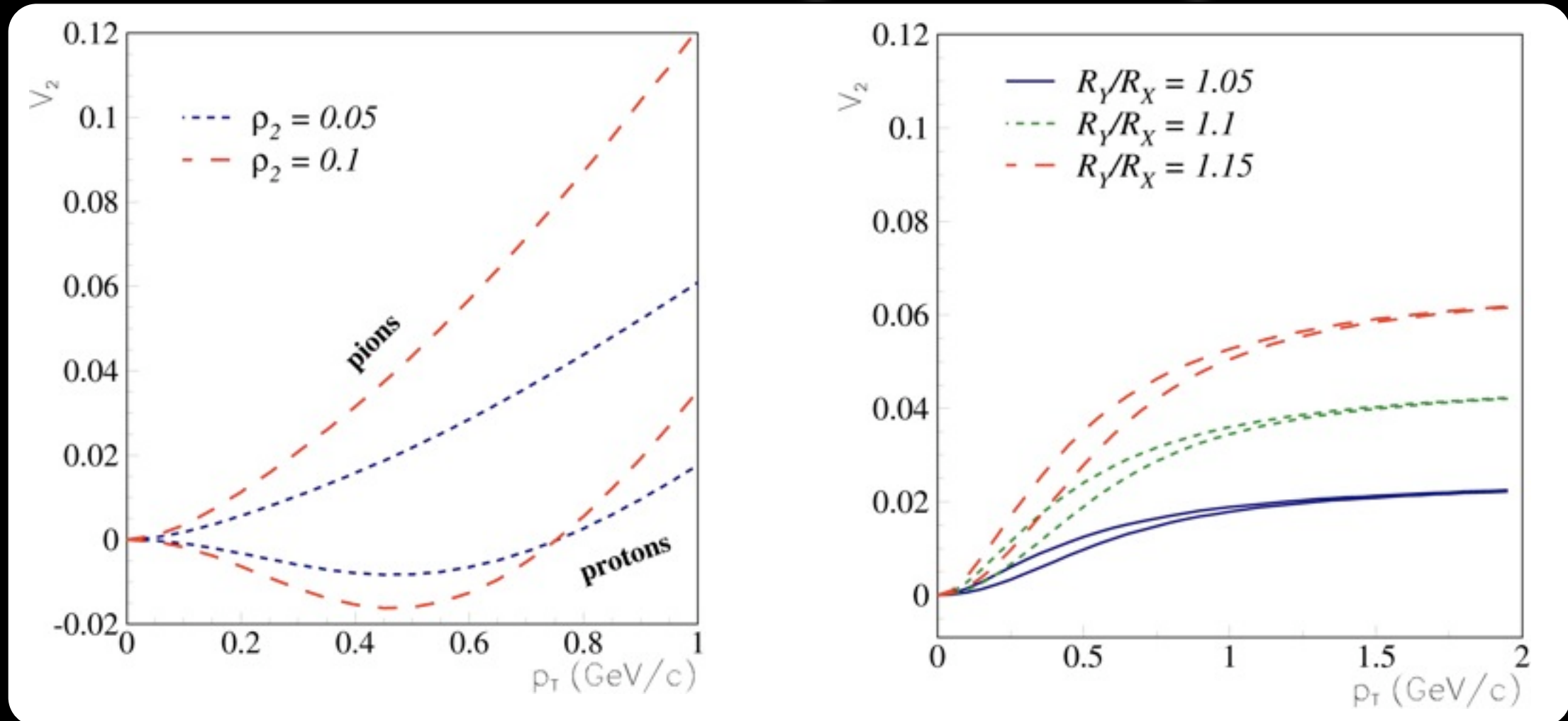
- on what **freeze-out** variables does it depend (simplification)?
- the average velocity difference in and out of plane (due to Δp)
- but also
 - the average freeze-out temperature
 - the average transverse flow
 - the average spatial eccentricity

The effect of freeze-out temperature and radial flow on v_2



- light particle $v_2(p_t)$ very sensitive to temperature
- heavier particles $v_2(p_t)$ more sensitive to transverse flow

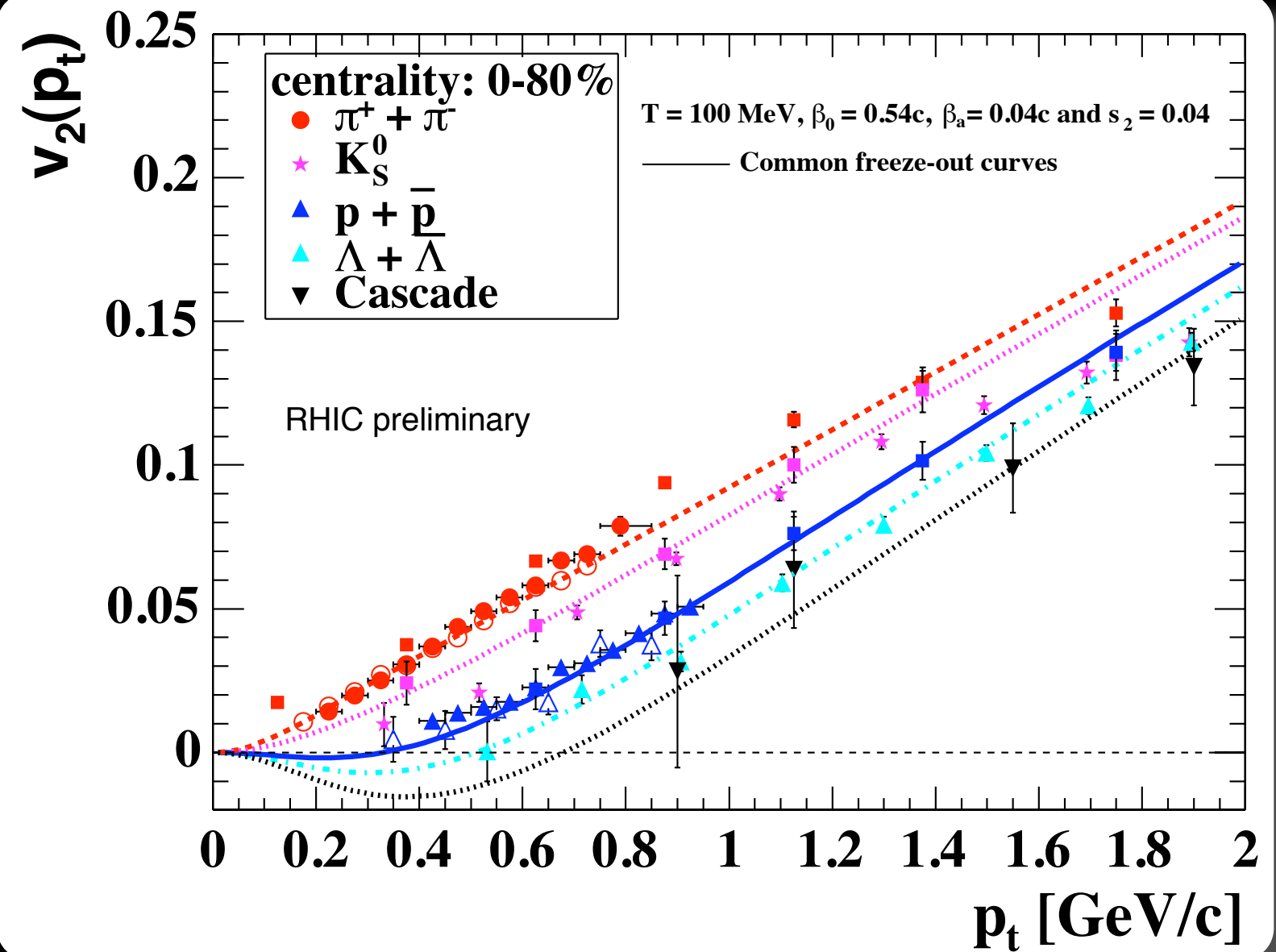
The effect of the azimuthal asymmetric flow velocity and shape



- larger value of the difference in collective velocity in and out of the reaction plane leads to larger slope of $v_2(p_T)$ above $\sim \langle p_T \rangle$ of the particle
- larger spatial anisotropy leads to larger v_2 with little mass dependence (transverse flow boosts more particles in the reaction plane)

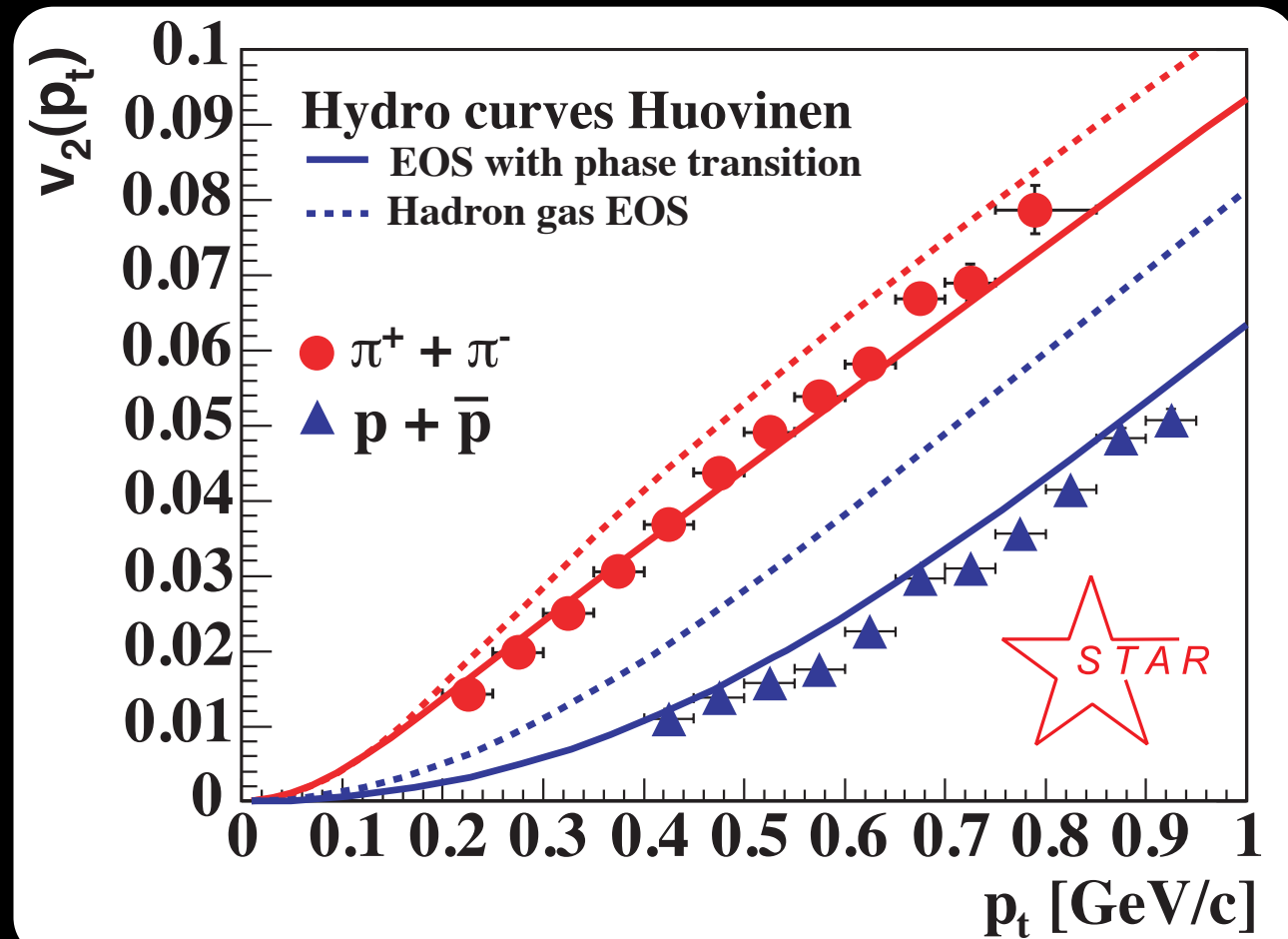
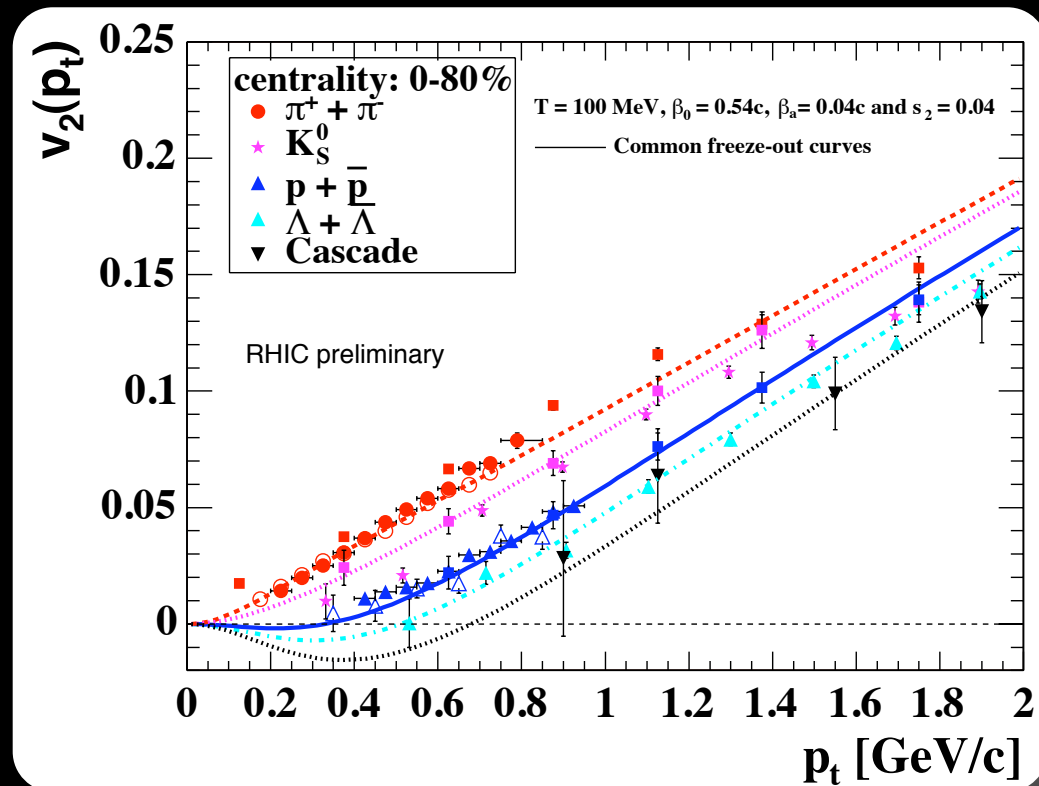
boosted thermal spectra

the observed particles are characterized by a single freeze-out temperature and a common azimuthal dependent boost velocity



Fits from STAR Phys. Rev. Lett. 87, 182301 (2001)

The EoS

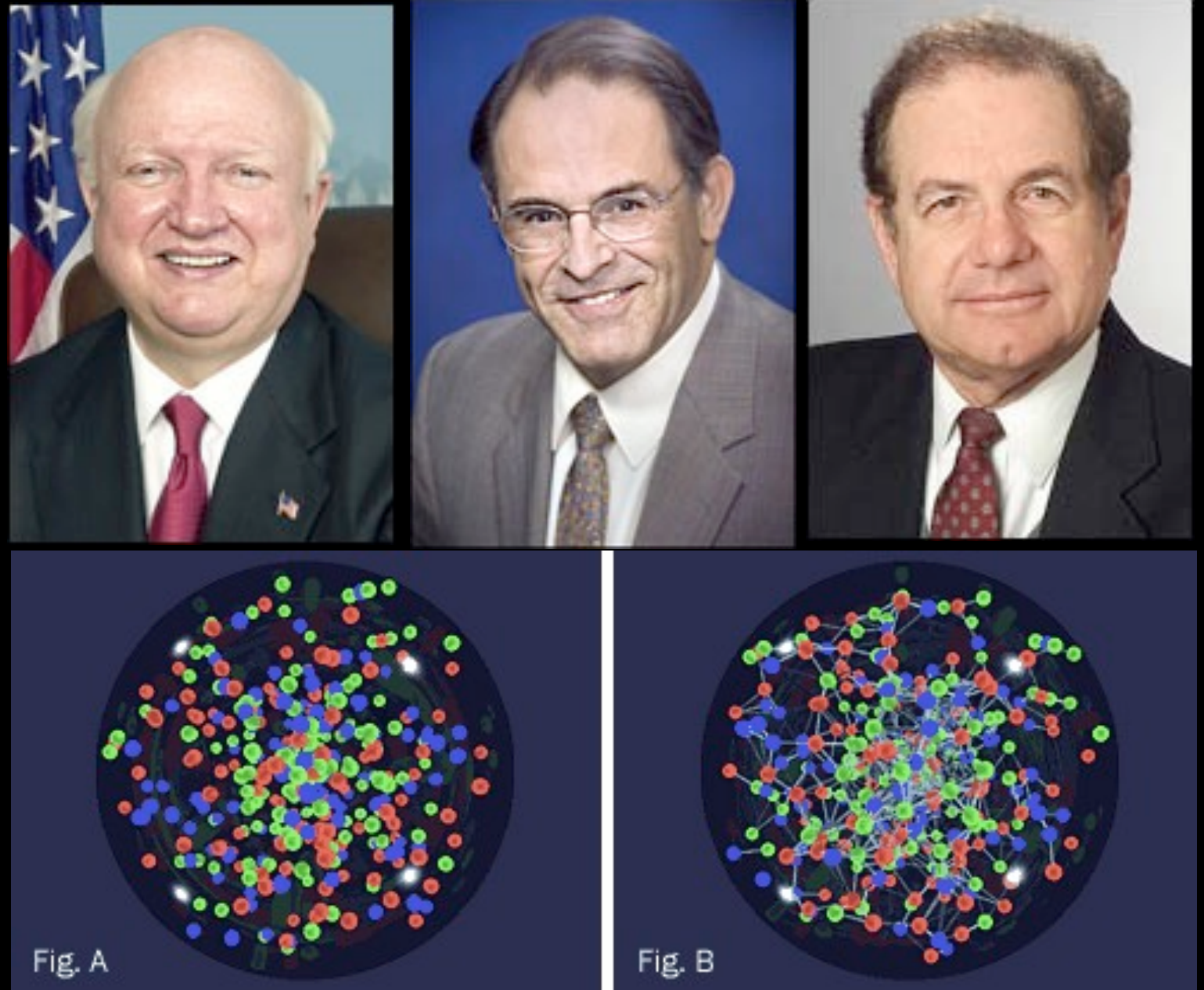
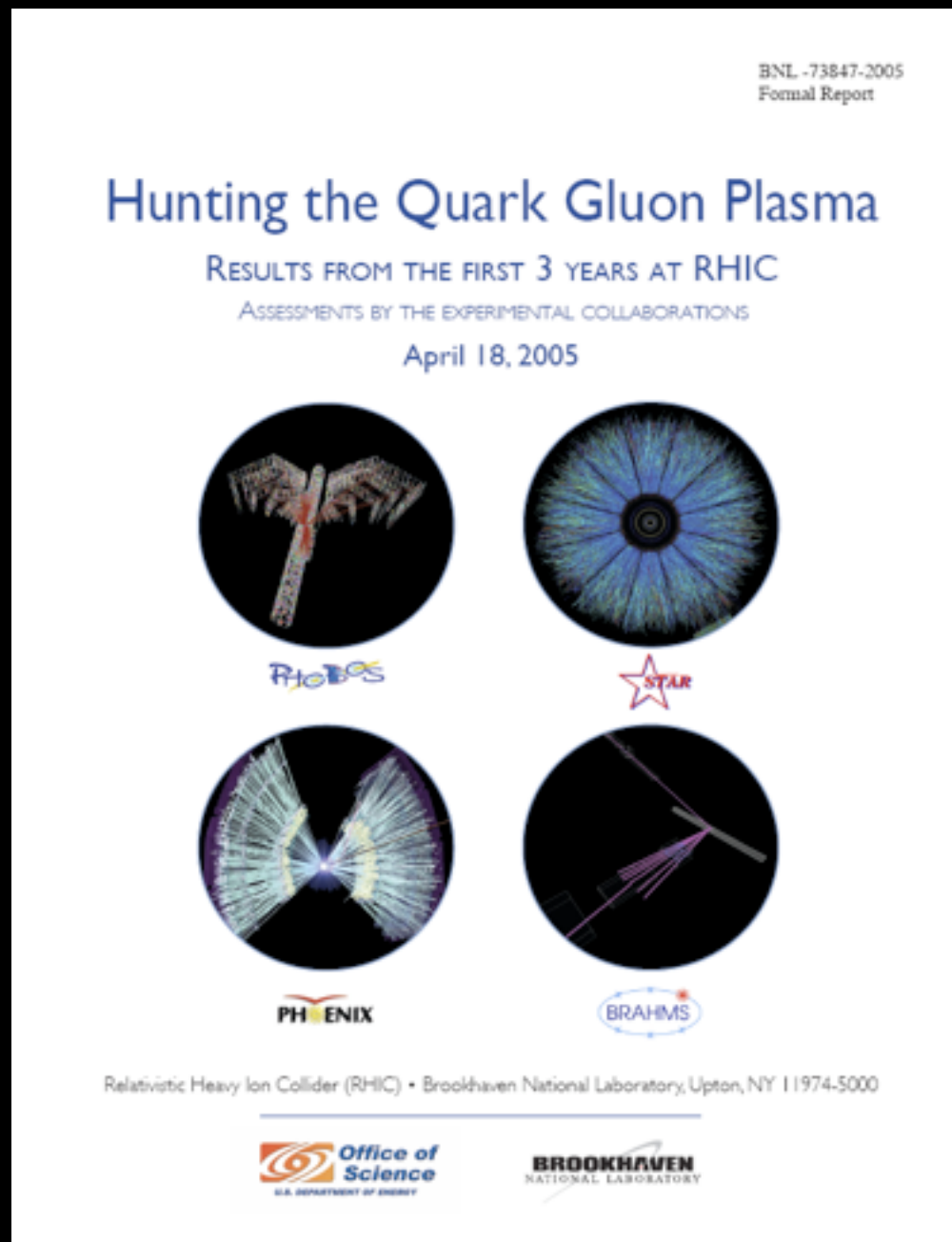


STAR Phys. Rev. Lett. 87, 182301 (2001)

The species dependence is sensitive to the EoS

RHIC Scientists Serve Up “Perfect” Liquid

New state of matter more remarkable than predicted --
raising many new questions
April 18, 2005

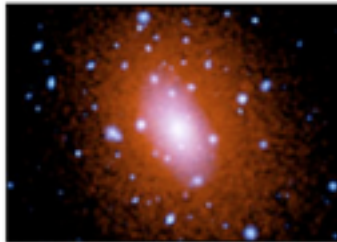


RHIC Scientists Serve Up "Perfect" Liquid

New state of matter more remarkable than predicted -- raising many new questions

April 18, 2005

Early Universe Went With the Flow



Posted April 18, 2005 5:57PM

Between 2000 and 2003 the lab's Relativistic Heavy Ion Collider repeatedly smashed the nuclei of gold atoms together with such force that their energy briefly generated trillion-degree temperatures. Physicists think of the collider as a time machine, because those extreme temperature conditions last prevailed in the universe less than 100 millionths of a second after the big bang.

Universe May Have Begun as Liquid, Not Gas

Associated Press
Tuesday, April 19, 2005; Page A05

The Washington Post

New results from a particle collider suggest that the universe behaved like a liquid in its earliest moments, not the fiery gas that was thought to have pervaded the first microseconds of existence.

Early Universe was a liquid

Quark-gluon blob surprises particle physicists.

by Mark Peplow
news@nature.com

nature

The Universe consisted of a perfect liquid in its first moments, according to results from an atom-smashing experiment.

Scientists at the Relativistic Heavy Ion Collider (RHIC) at Brookhaven National Laboratory on Long Island, New York, have spent five years searching for the quark-gluon plasma that is thought to have filled our Universe in the first microseconds of its existence. Most of them are now convinced they have found it. But, strangely, it seems to be a liquid rather than the expected hot gas.

New State of Matter Is 'Nearly Perfect' Liquid

Physicists working at Brookhaven National Laboratory announced today that they have created what appears to be a new state of matter out of the building blocks of atomic nuclei, quarks and gluons. The researchers unveiled their findings--which could provide new insight into the composition of the universe just moments after the big bang--today in Florida at a meeting of the American Physical Society.

**SCIENTIFIC
AMERICAN**

There are four collaborations, dubbed BRAHMS, PHENIX, PHOBOS and STAR, working at Brookhaven's Relativistic Heavy Ion Collider (RHIC). All of them study what happens when two interacting beams of gold ions smash into one another at great velocities, resulting in thousands of subatomic collisions every second. When the researchers analyzed the patterns of the atoms' trajectories after these collisions, they found that the particles produced in the collisions tended to move collectively, much like a school of fish does. Brookhaven's associate laboratory director for high energy and nuclear physics, Sam Aronson, remarks that "the degree of collective interaction, rapid thermalization and extremely low viscosity of the matter being formed at RHIC make this the most nearly perfect liquid ever observed."



Image: BNL

Early Universe was 'liquid-like'

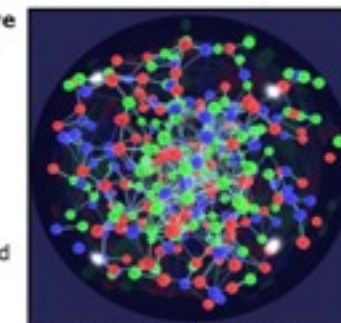
Physicists say they have created a new state of hot, dense matter by crashing together the nuclei of gold atoms.

BBC NEWS

The high-energy collisions prised open the nuclei to reveal their most basic particles, known as quarks and gluons.

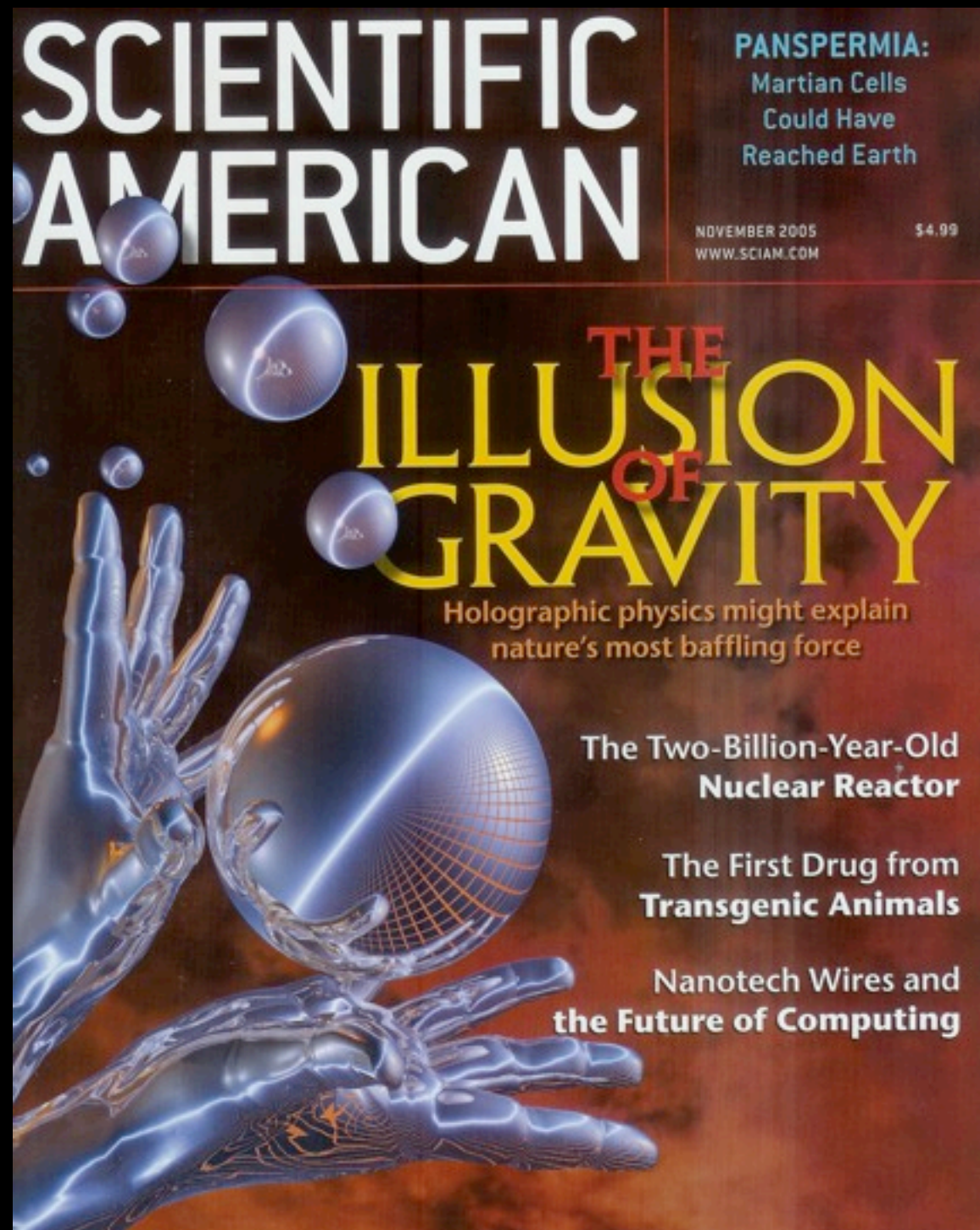
The researchers, at the US Brookhaven National Laboratory, say these particles were seen to behave as an almost perfect "liquid".

The work is expected to help scientists explain the conditions that existed just milliseconds after the Big Bang.



The impression is of matter that is more strongly interacting than predicted

AdS/CFT



the four force
I first con
graphic con
a specific the
dynamics in a
ary spacetime
ately excited
string theory
ture was made
Stephen S. Gu
of Princeton
Institute for A
ton, N.J. Son
have contribu
jectures and g
metastable an
theories, pro

So f

that it is com
ample has be
mathematics.

Mysteries
How does
ion of gravity
Black holes? I
enik Hawking
Stephen W. H
of Cambridge

sult. This radiation comes out of the black hole at a specific temperature. For all ordinary physical systems, a theory called statistical mechanics explains temperature in terms of the motion of the microscopic constituents. This theory explains the temperature of a glass of water or the temperature of the sun. What about the temperature of a black hole? To understand it, we would need to know what the microscopic constituents of the black hole are and how they behave. Only a theory of quantum gravity can tell us that.

Some aspects of the thermodynamics of black holes have raised doubts as to whether a quantum-mechanical theory of gravity could be developed at all. It seemed as if quantum mechanics itself might break down in the face of effects taking place in black holes. For a black

have an extremely low shear viscosity—smaller than any known fluid. Because of the holographic equivalence, strongly interacting quarks and gluons at high temperatures should also have very low viscosity.

A test of this prediction comes from the Relativistic Heavy Ion Collider (RHIC) at Brookhaven National Laboratory, which has been colliding gold nuclei at very high energies. A prelimi-

nary holographic theory for our universe; there is no convenient place to put the hologram.

An important lesson that one can draw from the holographic conjecture, however, is that quantum gravity, which has perplexed some of the best minds on the planet for decades, can be very simple when viewed in terms of the right variables. Let's hope we will soon find a simple description for the big bang! ■

MORE TO EXPLORE

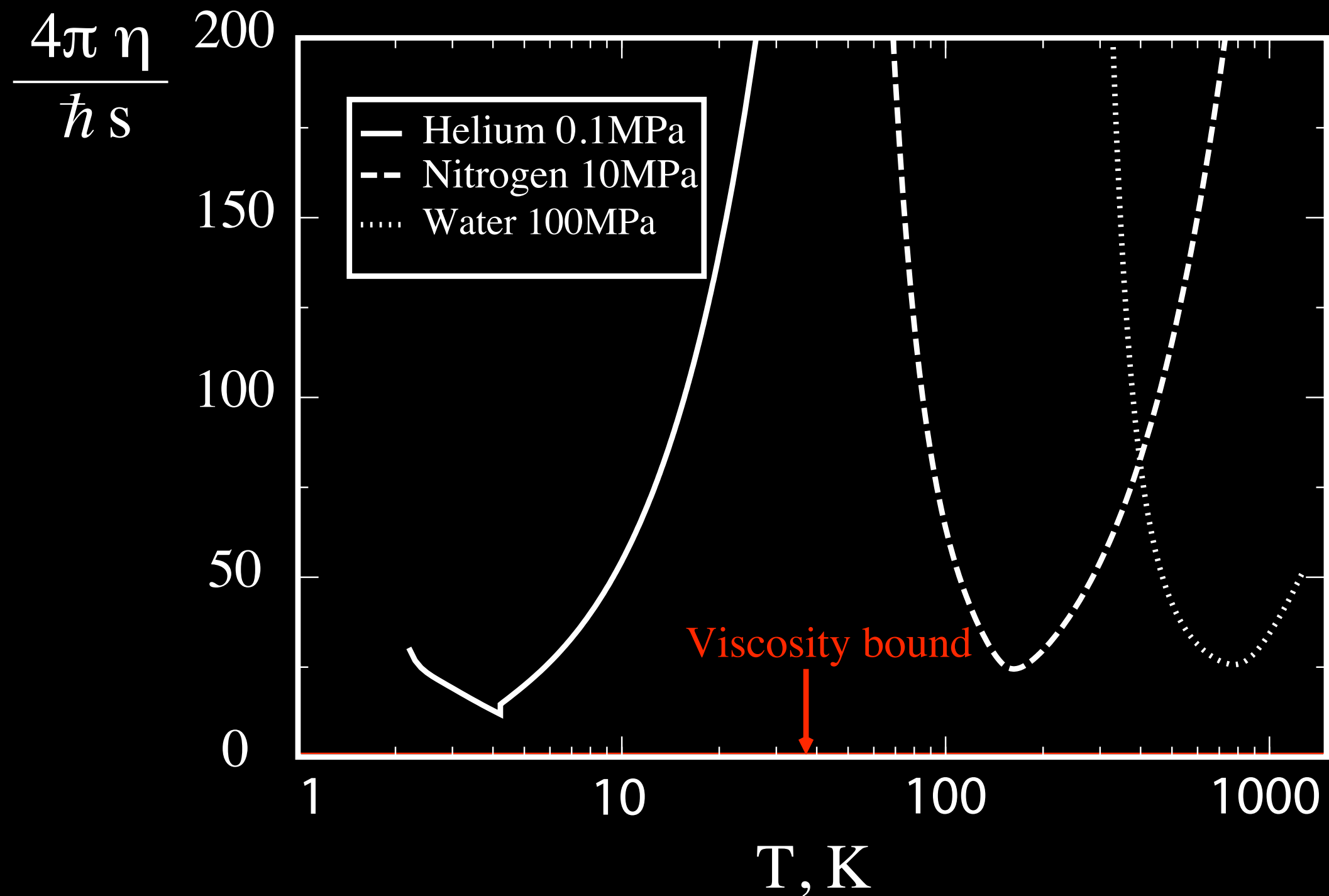
Anti-de Sitter Space and Holography. Edward Witten in *Advances in Theoretical and Mathematical Physics*, Vol. 2, pages 253–281, 1998. Available online at <http://arxiv.org/abs/hep-th/9802150>
Gauge Theory Correspondence from Non-Compact String Theory. S. Gubser, I. R. Klebanov and A. N. Polyakov in *Applied Physics Letters*, Vol. 428, pages 105–114, 1996. <http://arxiv.org/abs/hep-th/9602109>
The Theory Formerly Known as Strings. Michael J. Duff in *Scientific American*, Vol. 276, No. 2, pages 114–119, February 1997.
The Elegant Universe. Brian Greene. Published by W. W. Norton and Company, 2003.
A string theory Web site is at superstringtheory.com

www.sciam.com

SCIENTIFIC AMERICAN 63

November, 2005 Scientific American “The Illusion of Gravity” J. Maldacena

AdS/CFT



Kovtun, Son, Starinets, PRL 94 (2005) 111601

quantifying the viscosity

NA49: C. Alt et al., Phys. Rev. C68, 034903 (2003)

for ideal hydrodynamics:

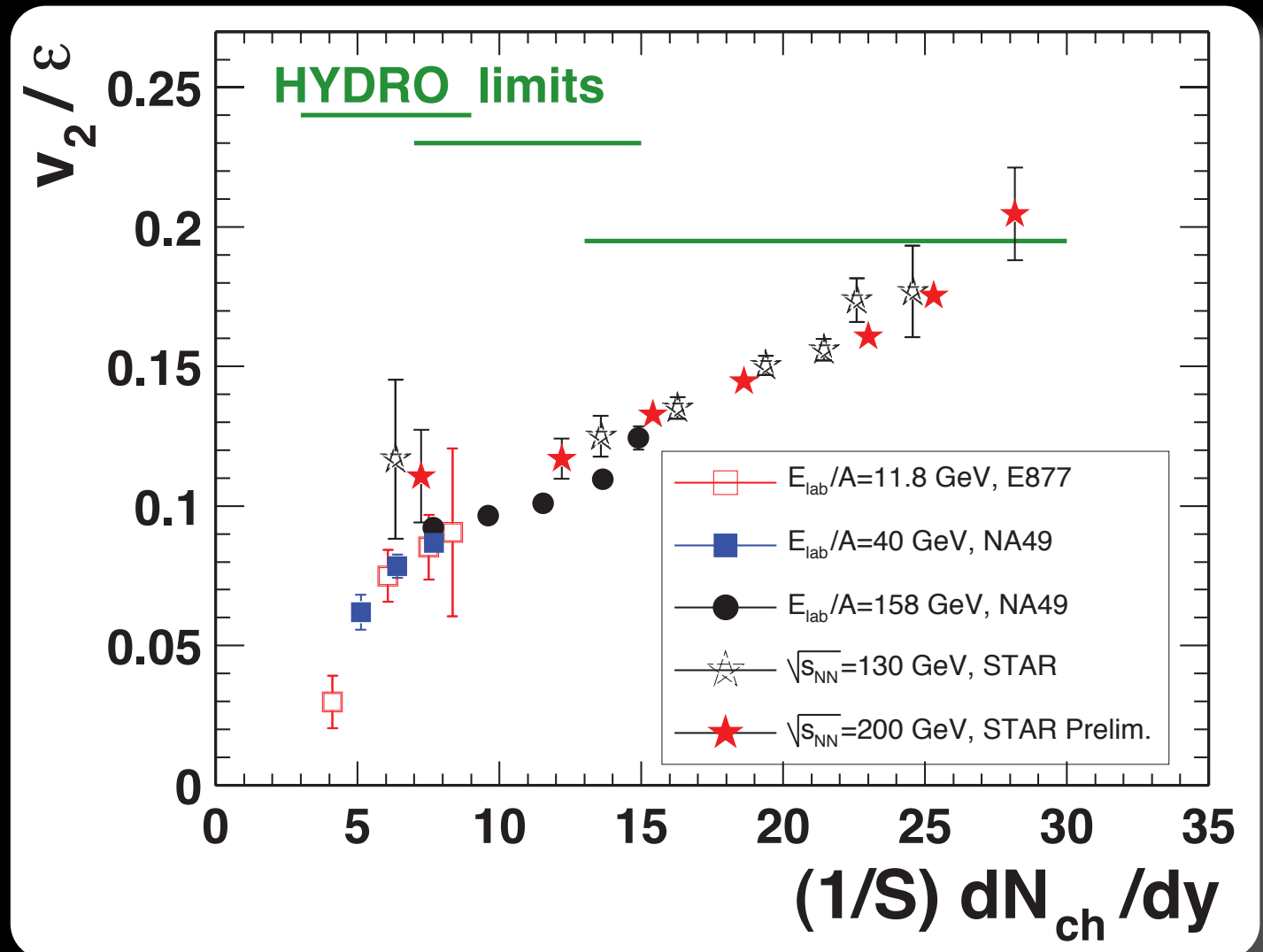
$$\frac{v_2}{\epsilon} = h$$

in the Low Density Limit (LDL):

$$\frac{v_2}{\epsilon} \propto \sigma \frac{1}{S} \frac{dN}{dy}$$

H. Heiselberg and A. M. Levy,
Phys. Rev. C 59, 2716 (1999)

S. A. Voloshin and A. M. Poskanzer,
Phys. Lett. B 474, 27 (2000)



Hydro limits from P.F. Kolb, J. Sollfrank, U.W. Heinz; Phys.Rev.C62:054909,2000.

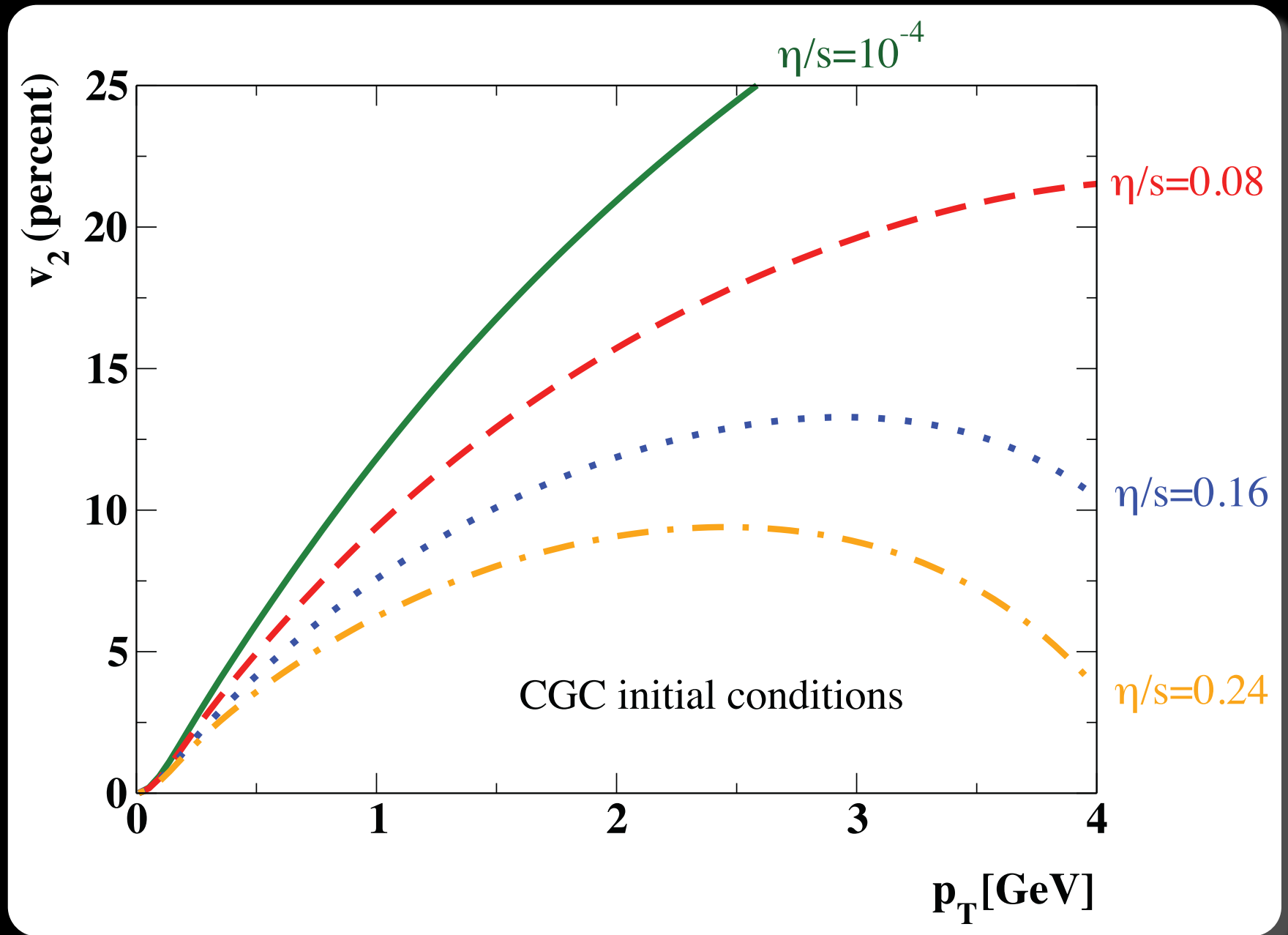
data not understood in ideal hydrodynamics alone!
corrections needed: parton cascade, viscous hydro, hadron
cascade

hydro Calculations of $v_2(p_t)$

small shear viscosity
(close to the
conjectured lower
bound) leads to a
significant reduction
of the predicted
elliptic flow

D. Teaney PRC68 034913 (2003)

in recent years
strong theoretical
effort to incorporate
viscous corrections
in the hydro
calculations



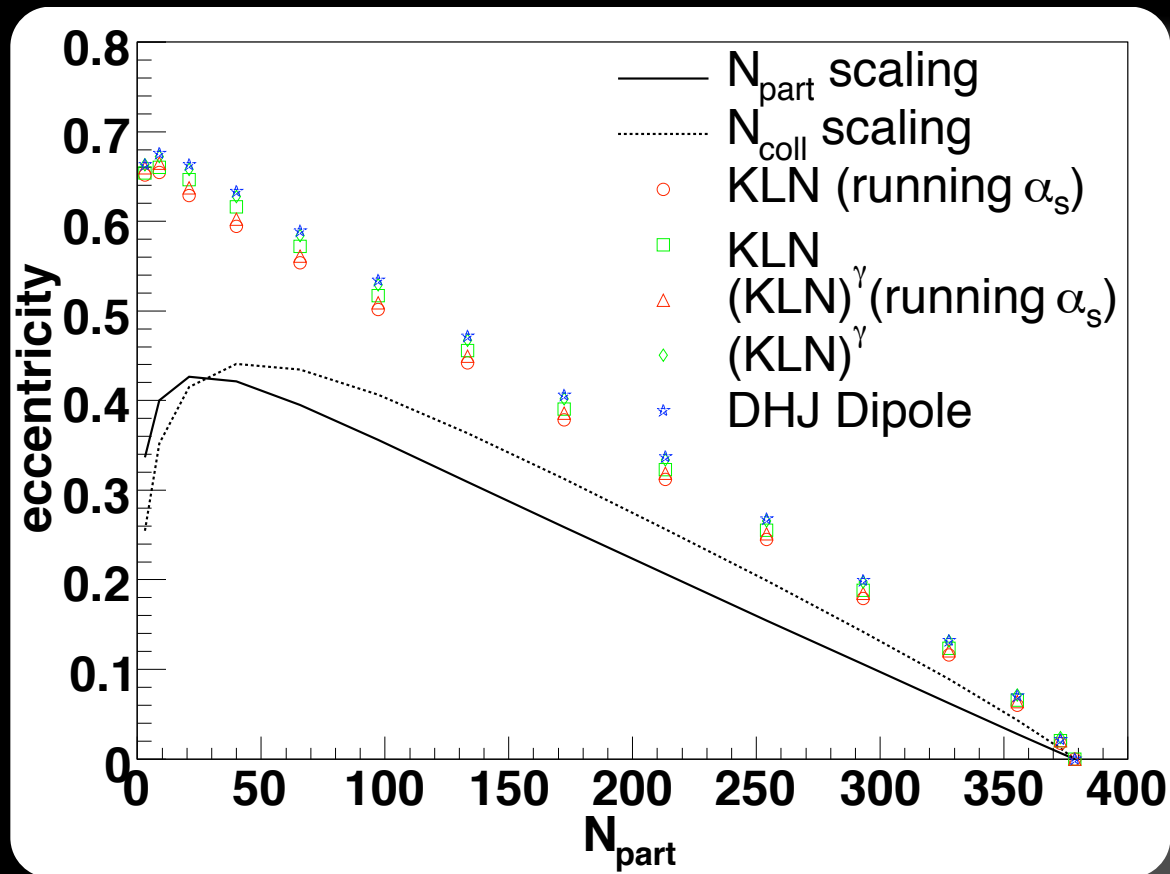
Matthew Luzum, Paul Romatschke arXiv:0804.4015

viscous corrections go as $p_t^2 \rightarrow$ shift of maximum v_2 vs p_t

towards estimates of η/s

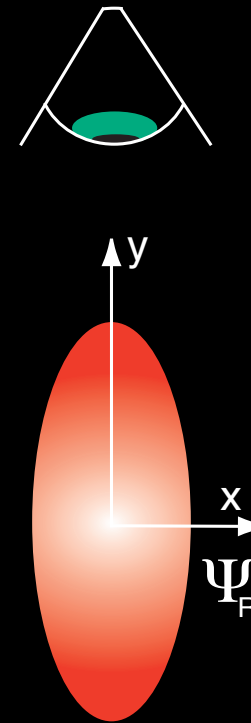
- strong experimental/theoretical effort to understand better the uncertainties for the determination of v_2 for the most central and peripheral collisions, and at higher transverse momenta
- strong theoretical/experimental effort to quantify the uncertainties in the initial conditions (main contribution is the eccentricity, mean magnitude and its fluctuation)
- uncertainties are now better understood (next slides)

initial conditions: ϵ

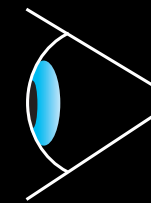


H-J. Drescher et al., Phys.Rev.C74:044905,2006

- estimates of the eccentricity vary significantly $\sim 30\%$!
- difference between CGC and Glauber N_{part} scaling largest
- v_2 follows ϵ



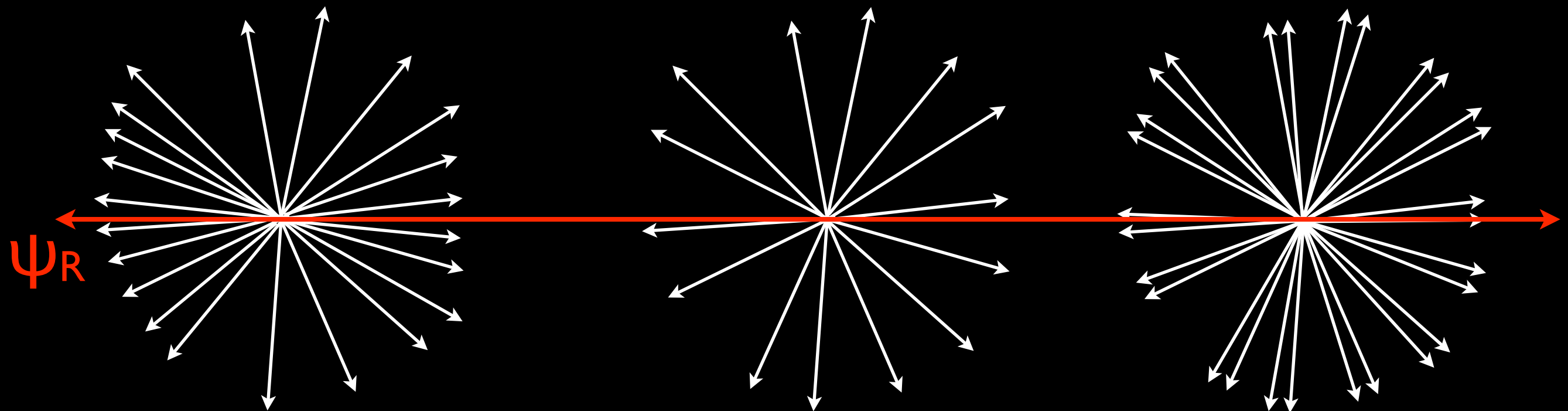
$$\epsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$



measure elliptic flow

$$v_n = \langle \cos(n(\phi - \Psi_R)) \rangle = \langle e^{in(\phi - \Psi_R)} \rangle$$

$$\langle e^{in(\phi_i - \phi_j)} \rangle = \langle e^{in(\phi_i - \Psi_R)} e^{in(\Psi_R - \phi_j)} \rangle \approx \langle e^{in(\phi_i - \Psi_R)} \rangle \langle e^{in(\Psi_R - \phi_j)} \rangle = (v_n \{2\})^2$$



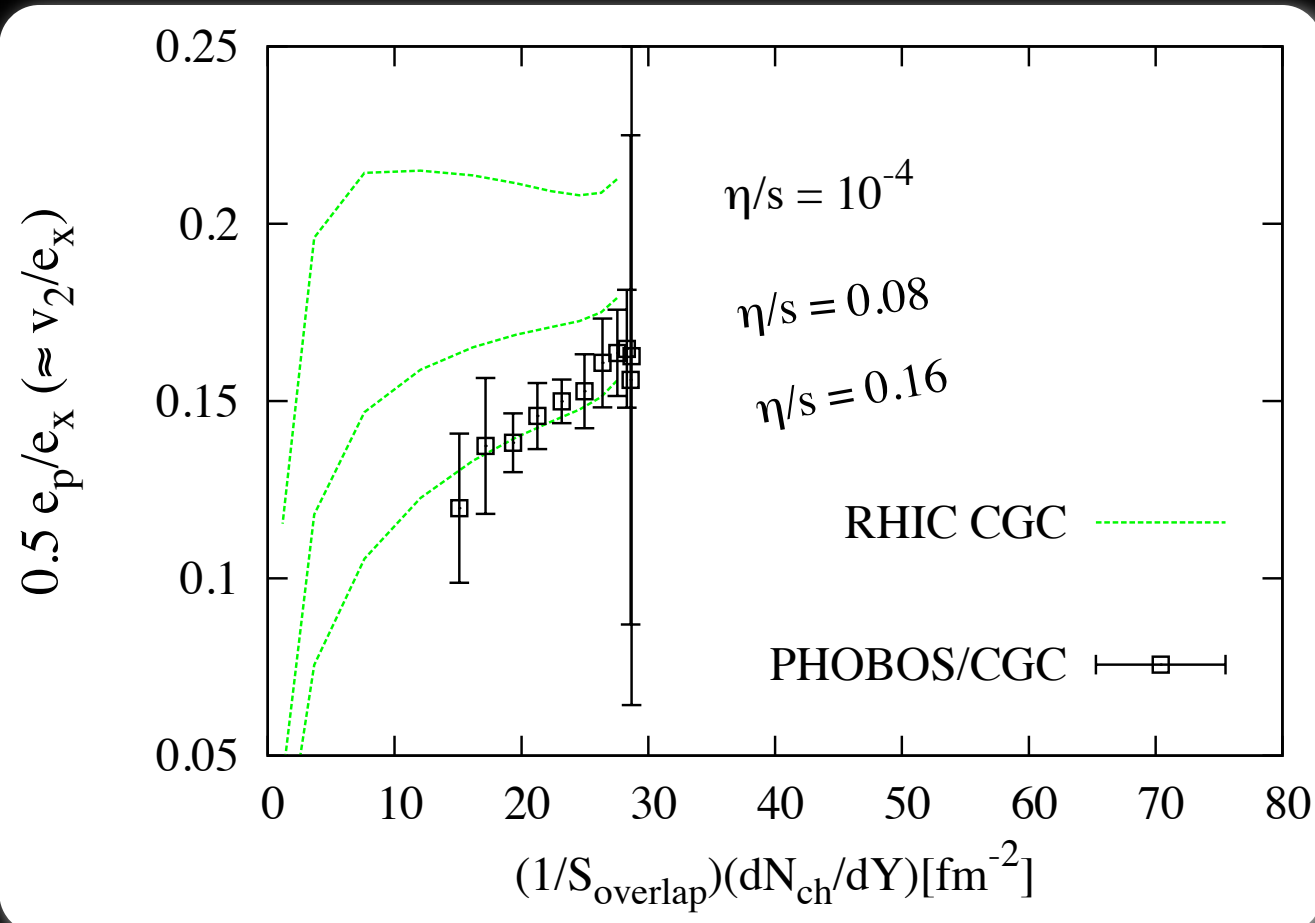
$$v_2 > 0, v_2 \{2\} > 0$$

$$v_2 = 0, v_2 \{2\} = 0$$

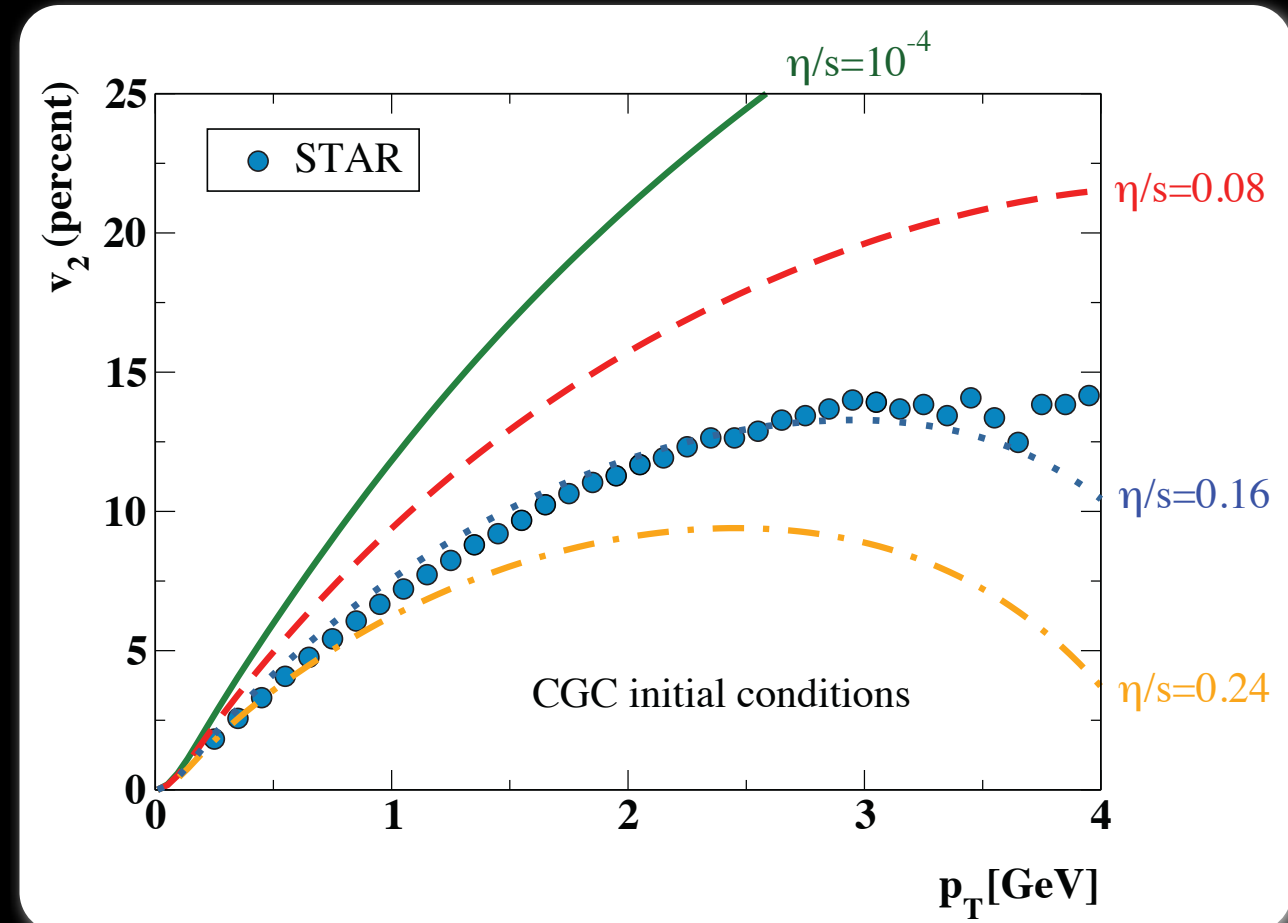
$$v_2 = 0, v_2 \{2\} > 0$$

use more sophisticated multiparticle correlation methods

viscous hydro CGC



Matthew Luzum and Paul Romatschke; arXiv:0901.4588 [nucl-th]



Matthew Luzum, Paul Romatschke arXiv:0804.4015

viscous hydro calculations using \sim lattice EoS and CGC ε describe the centrality dependence and p_t dependence using $\eta/s = 2/4\pi$

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April 18, 2005

- now improved quantitative understanding of the created system at RHIC
- viscous correction are important!
- initial conditions (e.g. ε , initial flow fields) not sufficiently constrained!
- \sim lattice EoS ☺ + \sim CGC ε ☺ $\rightarrow \sim \eta/s = 2/4\pi$ ☺ describes centrality and transverse momentum dependence of v_2 very well!
 - current estimates of η/s are $< 4x$ the conjectured lower bound from AdS/CFT

highlights at RHIC

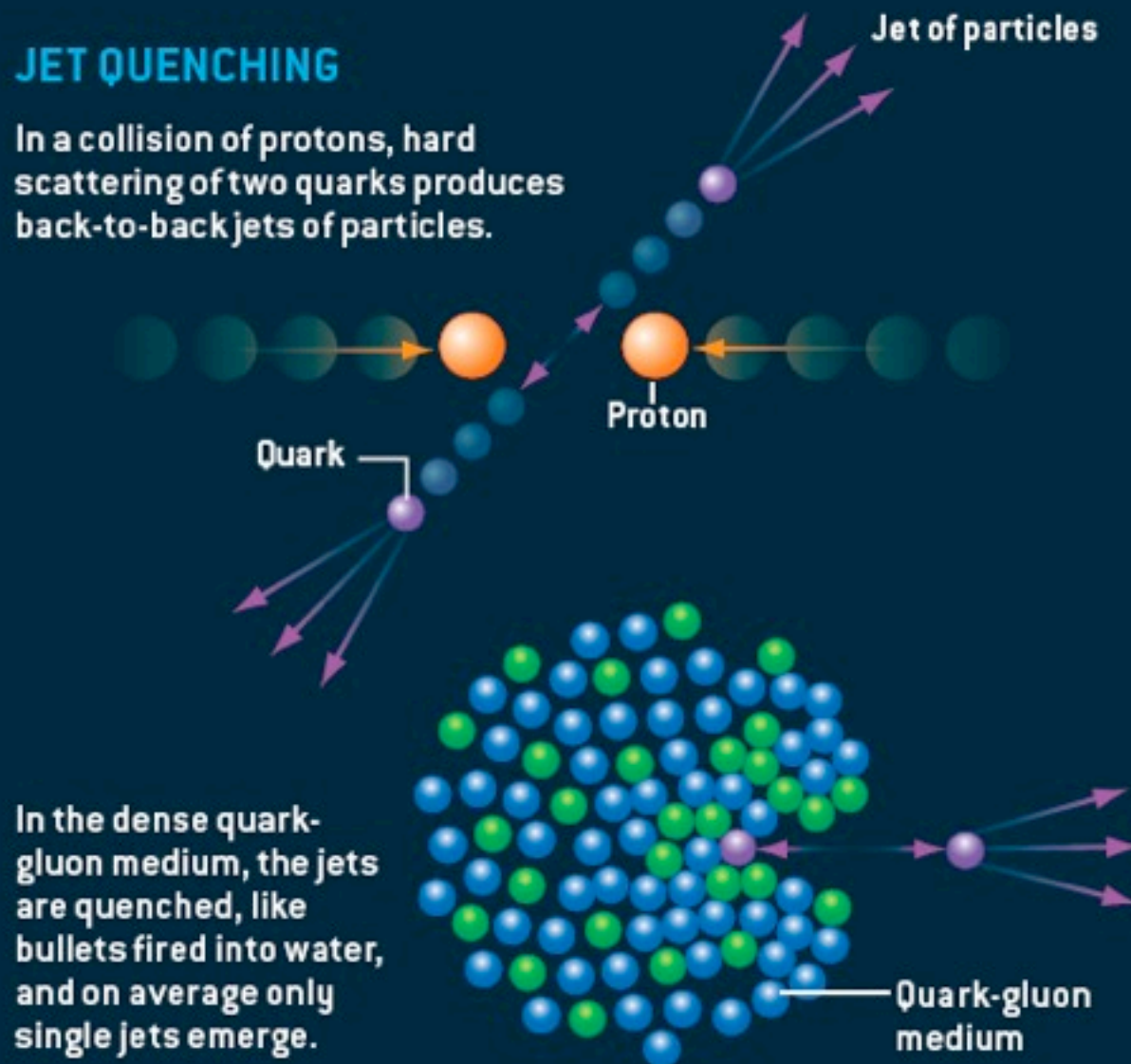
EVIDENCE FOR A DENSE LIQUID

M. Roirdan and W. Zajc, Scientific American 34A May (2006)

Two phenomena in particular point to the quark-gluon medium being a dense liquid state of matter: jet quenching and elliptic flow. Jet quenching implies the quarks and gluons are closely packed, and elliptic flow would not occur if the medium were a gas.

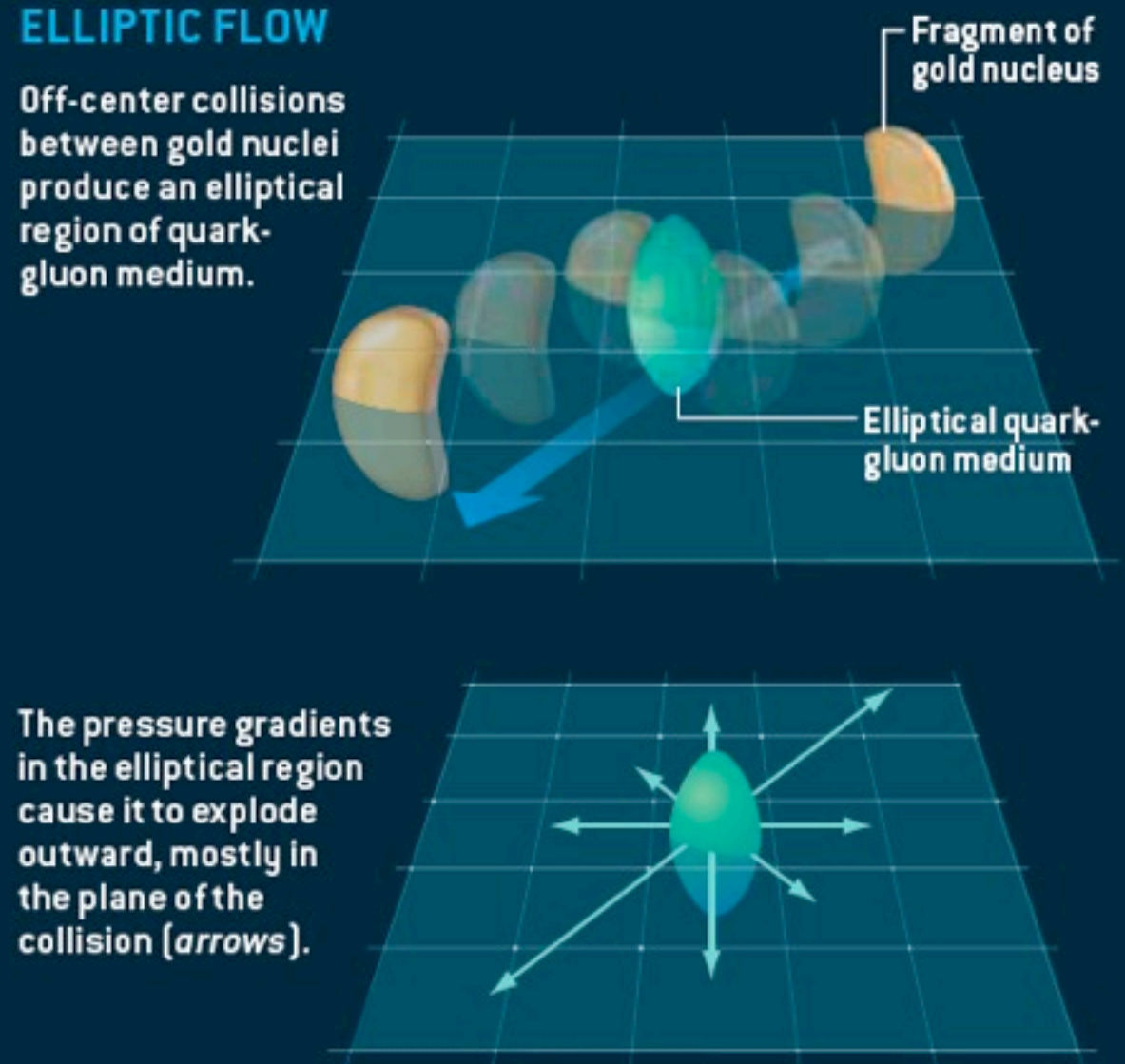
JET QUENCHING

In a collision of protons, hard scattering of two quarks produces back-to-back jets of particles.

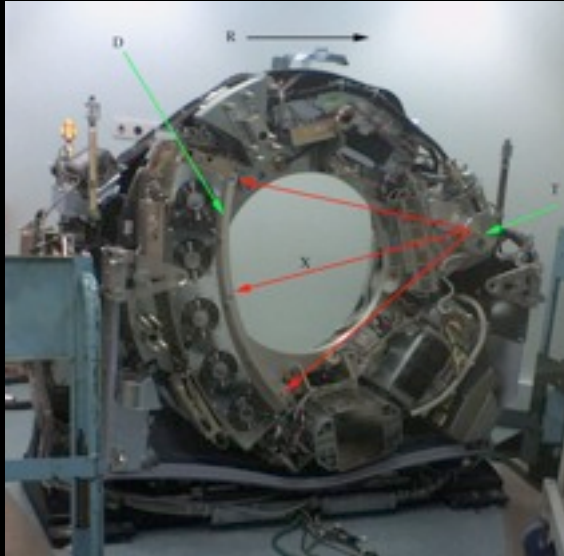


ELLIPTIC FLOW

Off-center collisions between gold nuclei produce an elliptical region of quark-gluon medium.



parton energy loss



probe the density of the created system

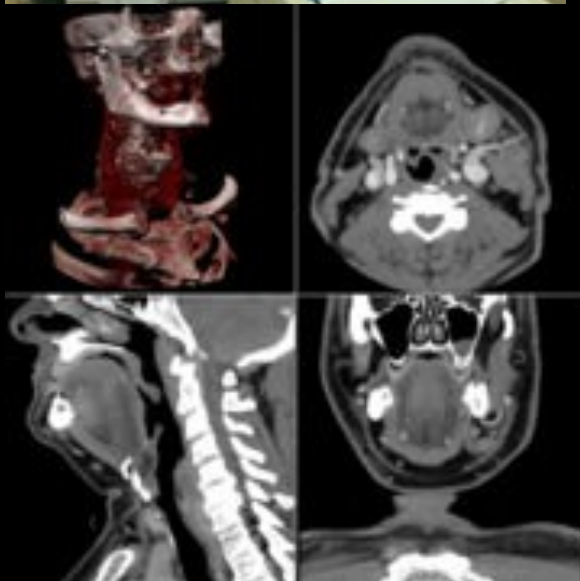
probe needs to be calibrated

✓ high momentum partons

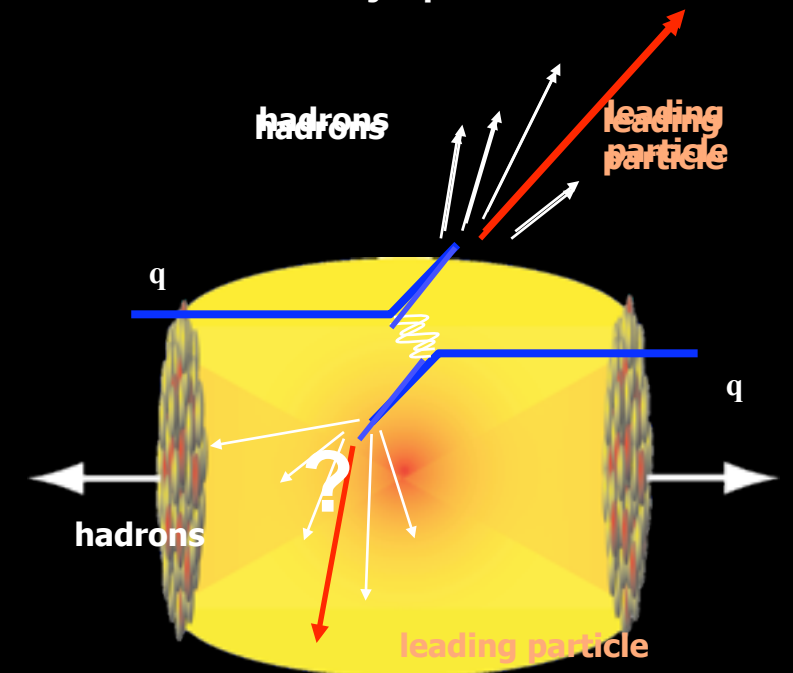
probe needs to interact with the matter

✓ parton energy loss

interactions should be understood and calculable

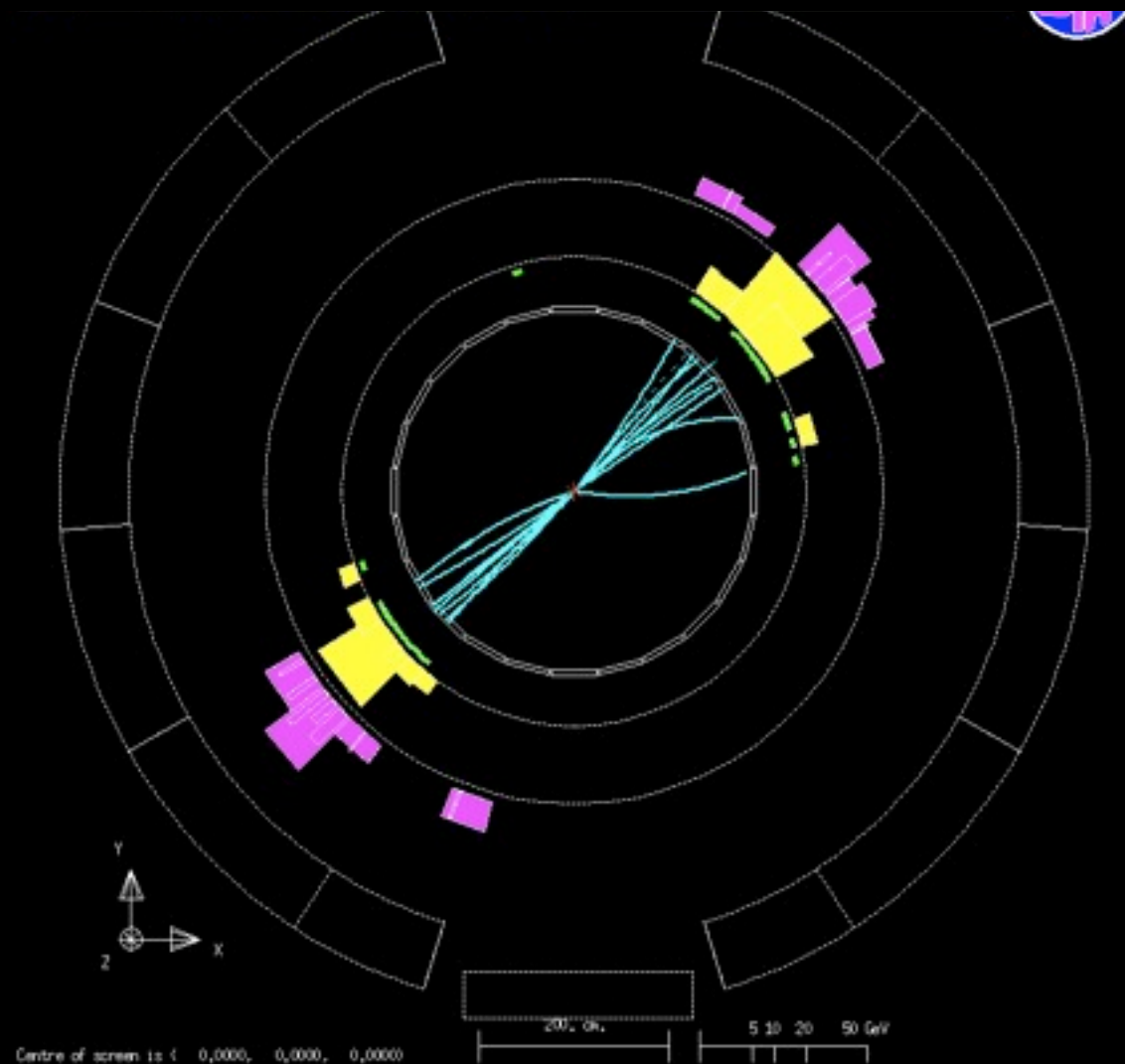


schematic view of jet production

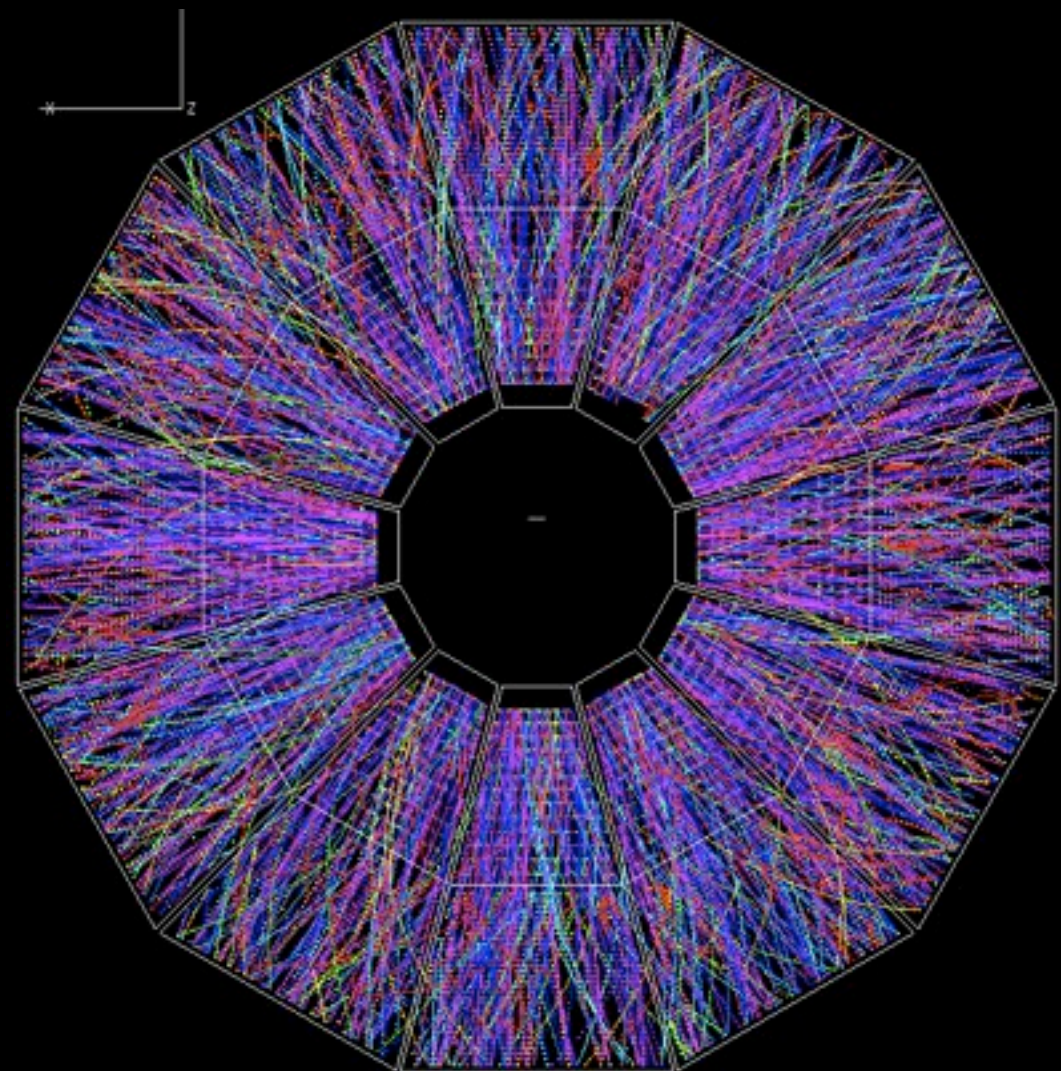


jets in a heavy-ion environment

Jets in e^+e^-



Jets in Au + Au at 200 GeV



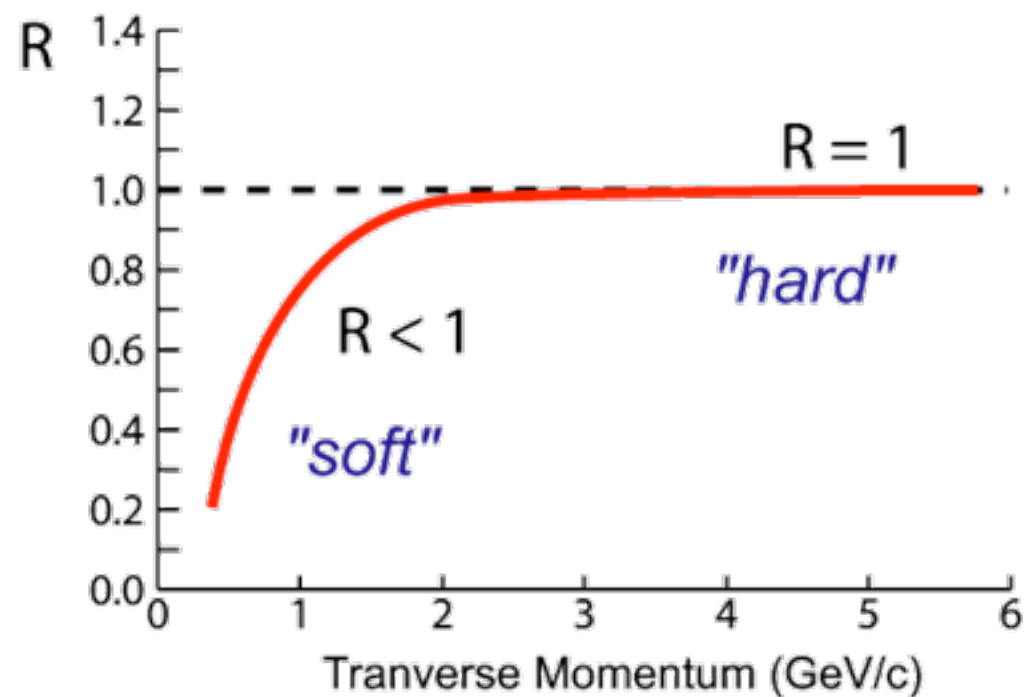
heavy-ion collisions are a complicated environment
to do full jet reconstruction!

parton energy loss

we measure the particle yield as
function of transverse momentum in
nucleus-nucleus and in nucleon+nucleon
collisions

we construct the ratio:

$$R = \frac{\text{Yield}_{A+A} / \langle N_{\text{binary}} \rangle}{\text{Yield}_{p+p}}$$



no "nuclear effects":
 $R < 1$ in regime of soft physics
 $R = 1$ at high- p_t where hard
scattering dominates

suppression:
 $R < 1$ at high- p_t

parton energy loss

$$R = \frac{\text{Yield}_{A+A} / \langle N_{\text{binary}} \rangle}{\text{Yield}_{p+p}}$$

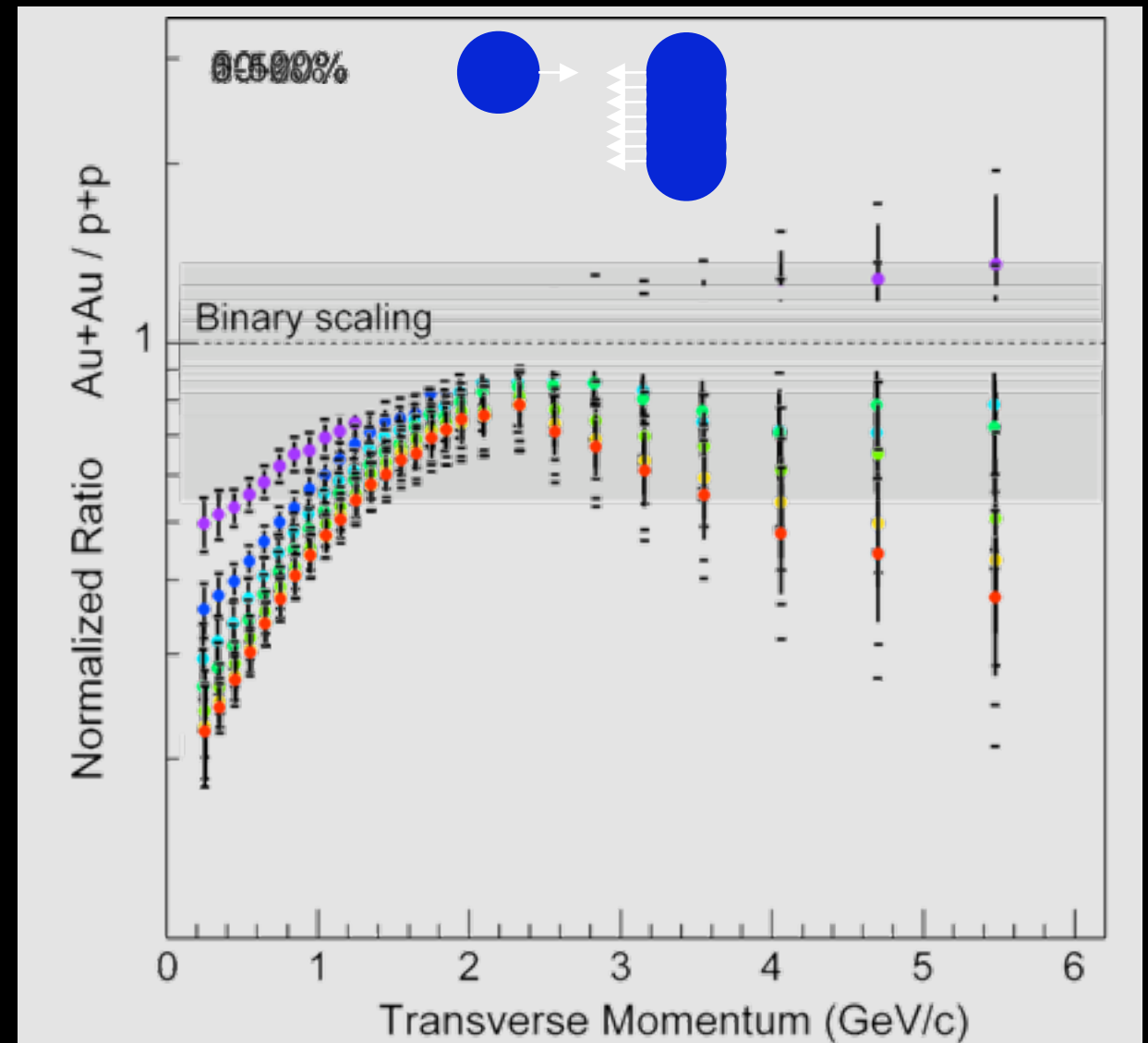
no “nuclear effects”:

$R < 1$ in regime of soft physics

$R = 1$ at high- p_t where hard scattering dominates

suppression:

$R < 1$ at high- p_t



parton energy loss

$$R = \frac{\text{Yield}_{A+A} / \langle N_{\text{binary}} \rangle}{\text{Yield}_{p+p}}$$

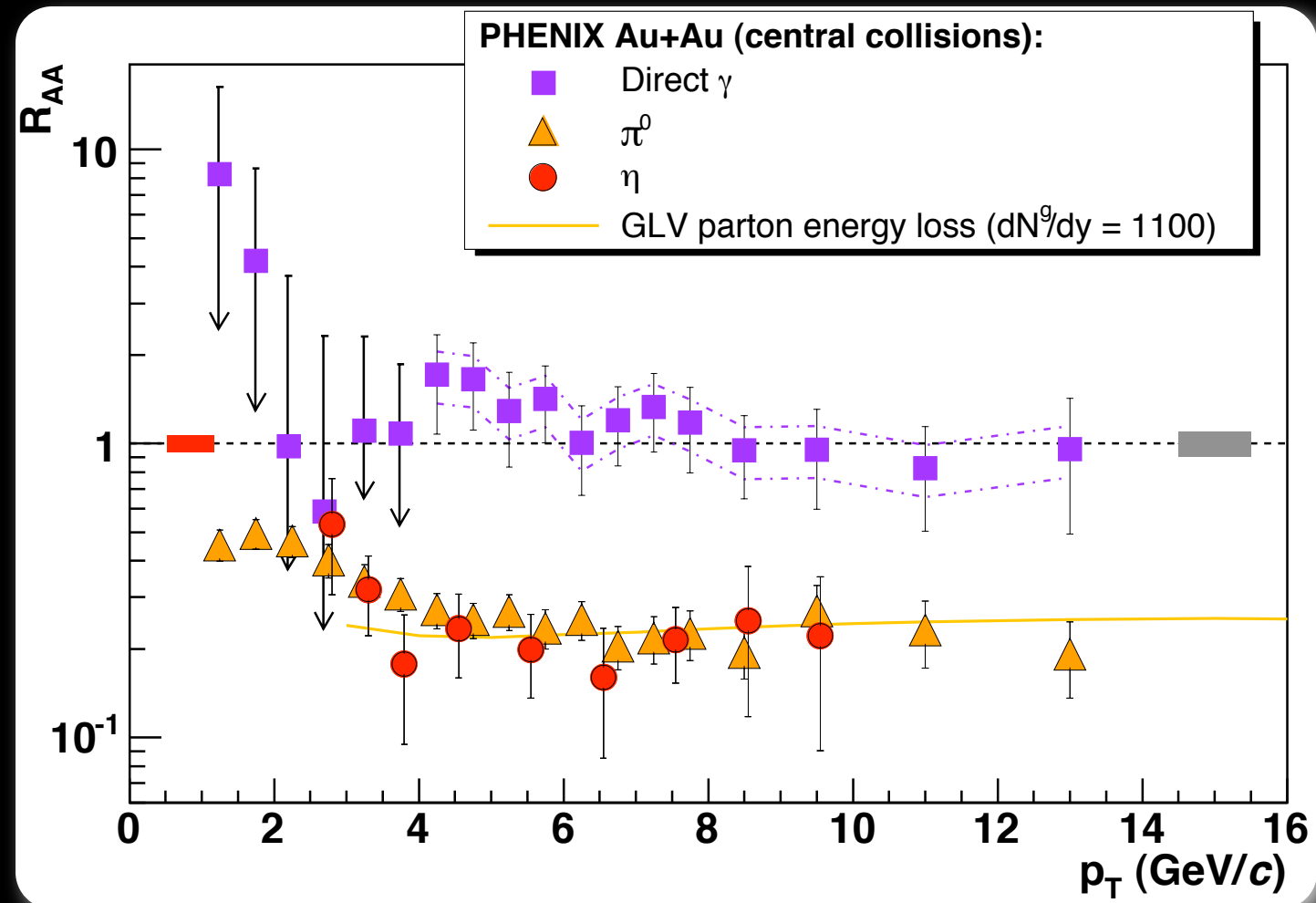
no “nuclear effects”:

$R < 1$ in regime of soft physics

$R = 1$ at high- p_t where hard scattering dominates

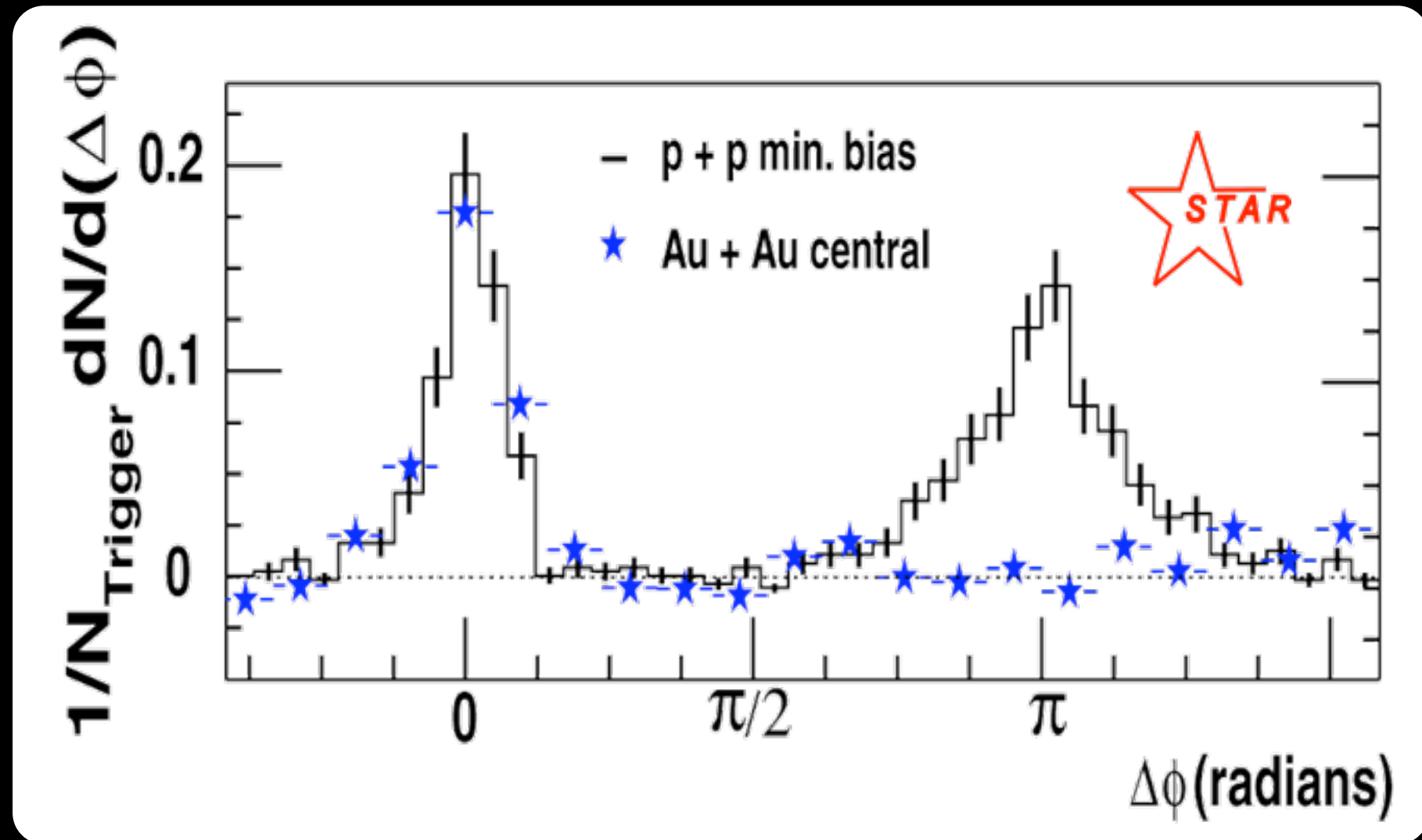
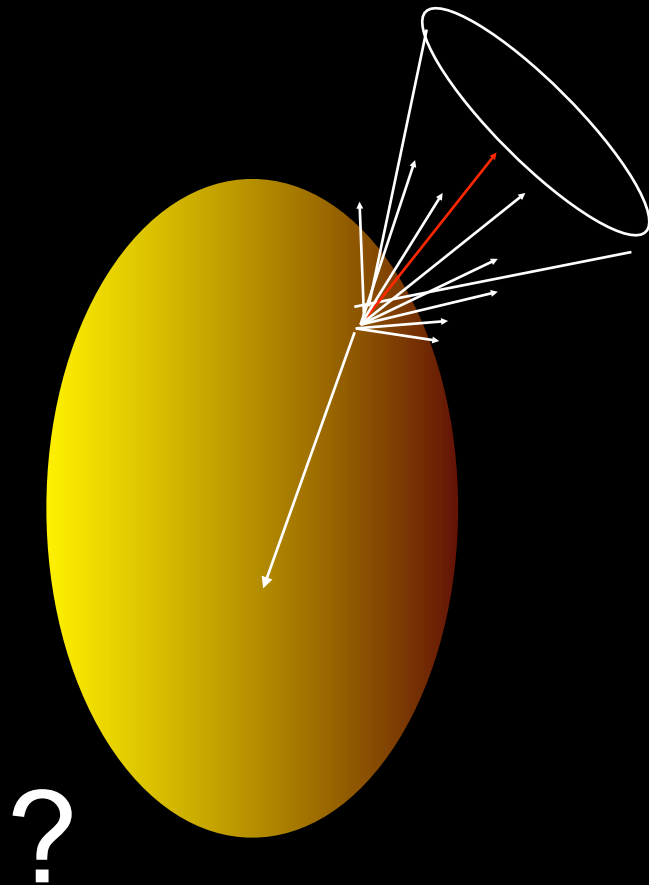
suppression:

$R < 1$ at high- p_t



very strong suppression!
medium density extracted 30-50 times normal
nuclear matter density

parton energy loss



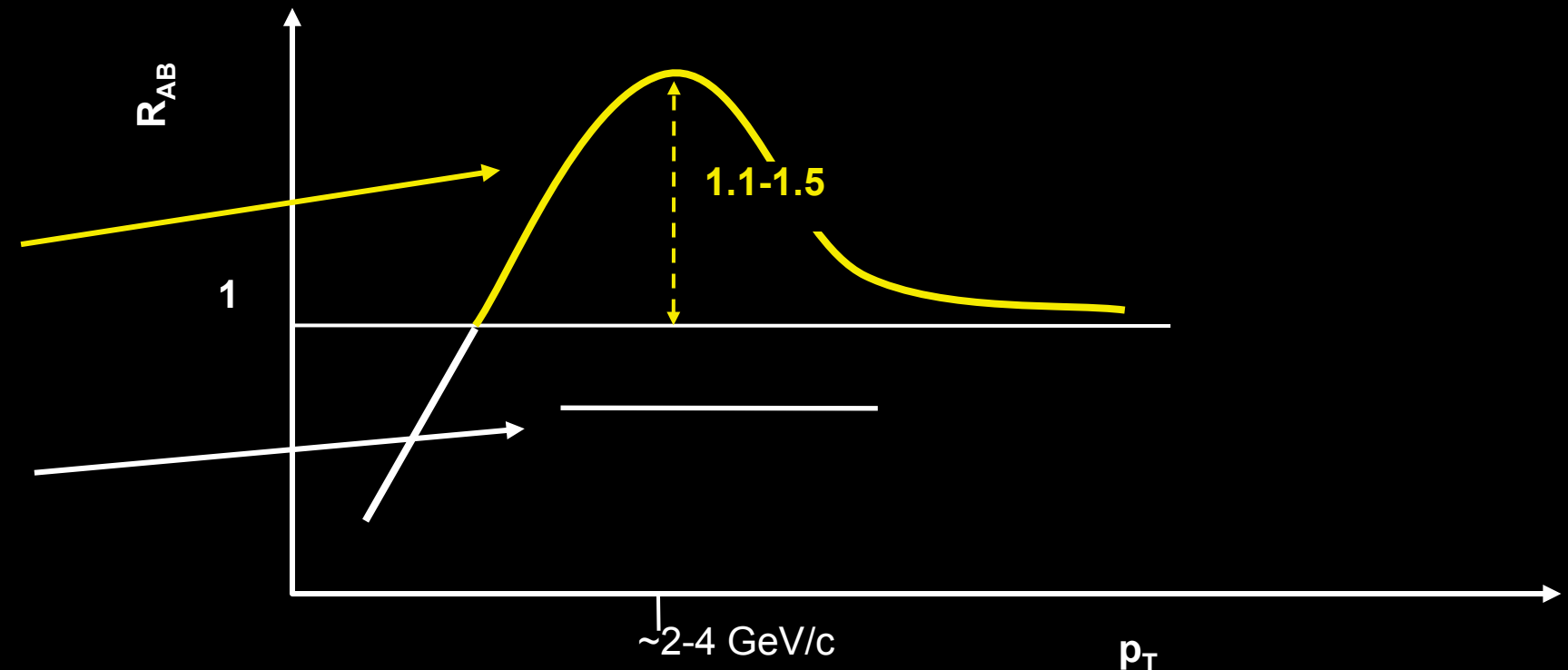
the away side disappears completely!
medium density extracted 30-50 times normal
nuclear matter density

what to expect in d+A?

Inclusive spectra

If Au+Au suppression is final state

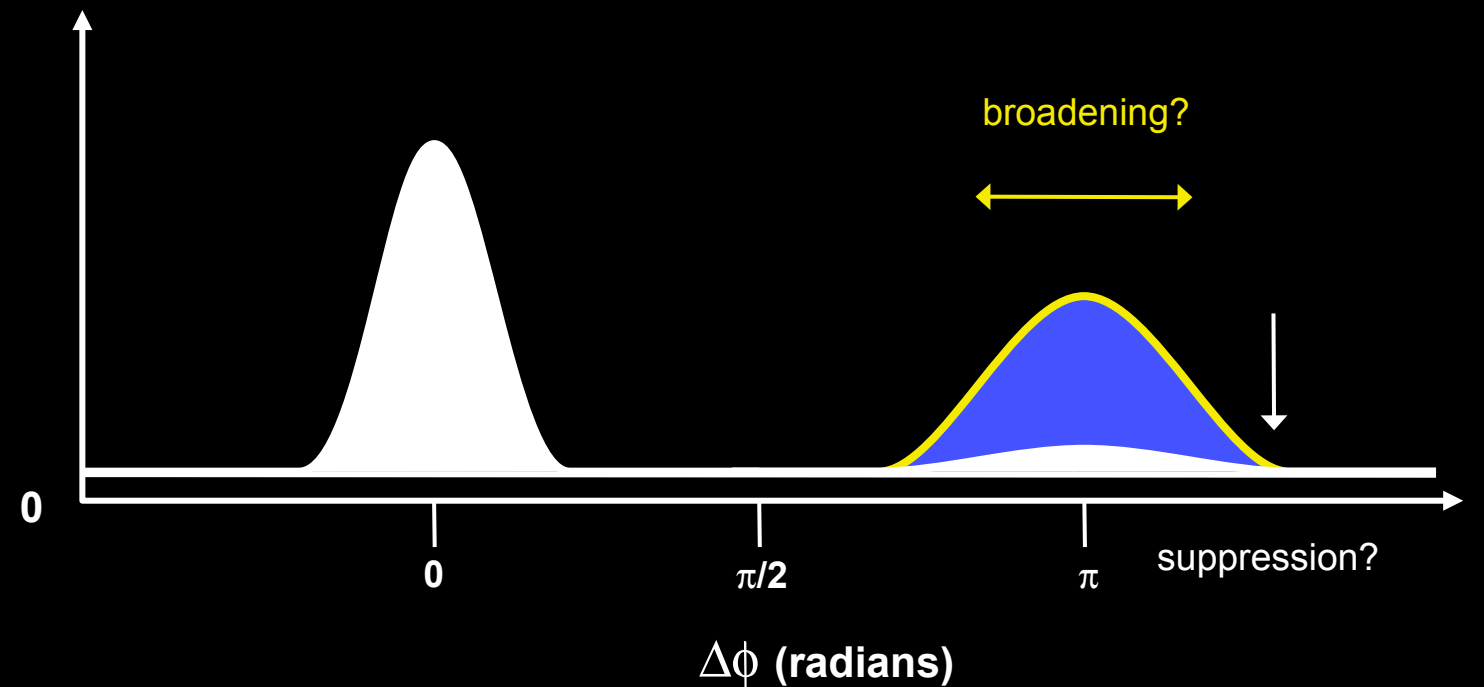
If Au+Au suppression is initial state (KLM saturation: 0.75)



High p_T hadron pairs

pQCD: no suppression, small broadening due to Cronin effect

saturation models:
suppression due to mono-jet contribution?

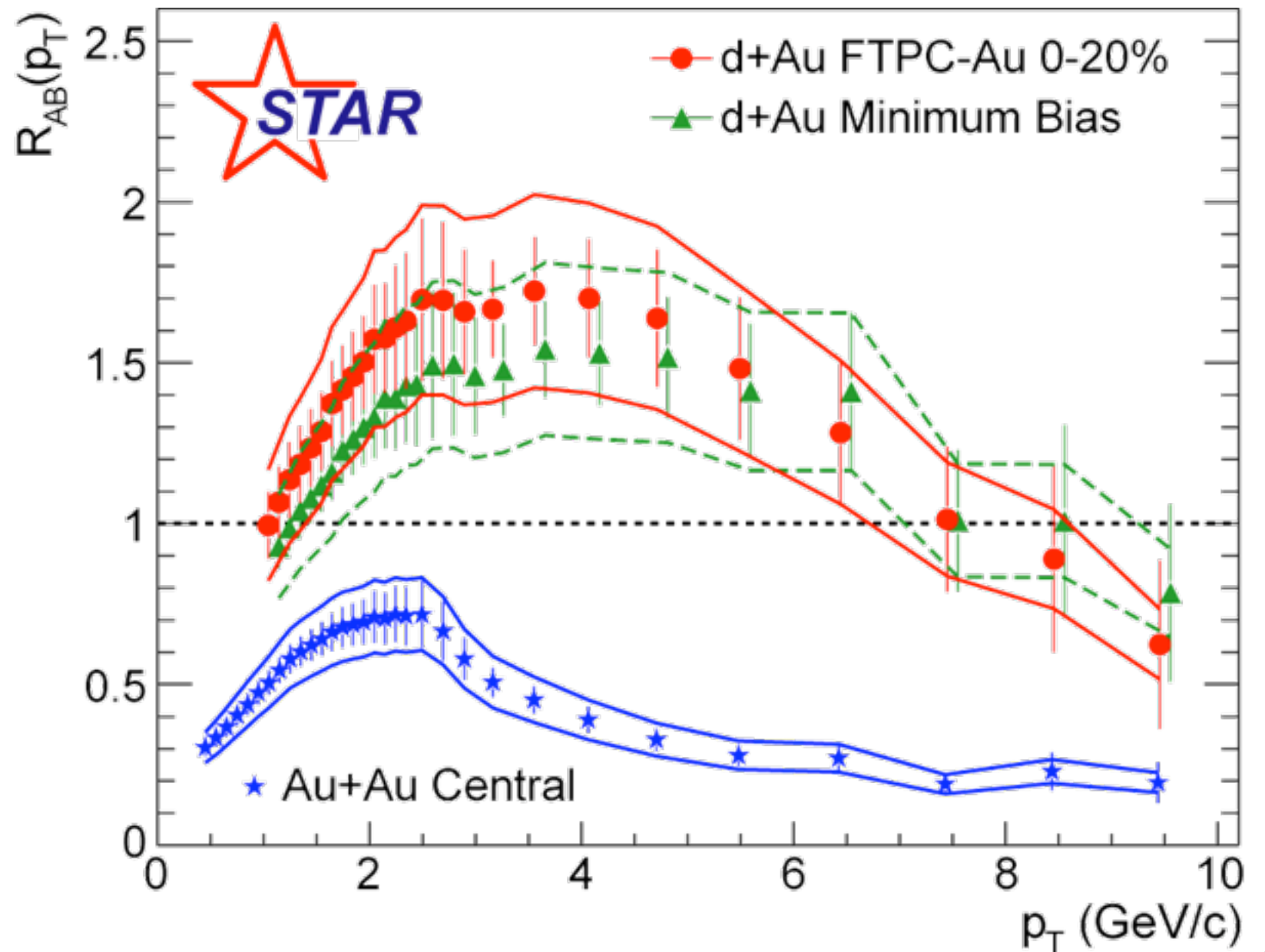


All effects strongest in central d+Au collisions

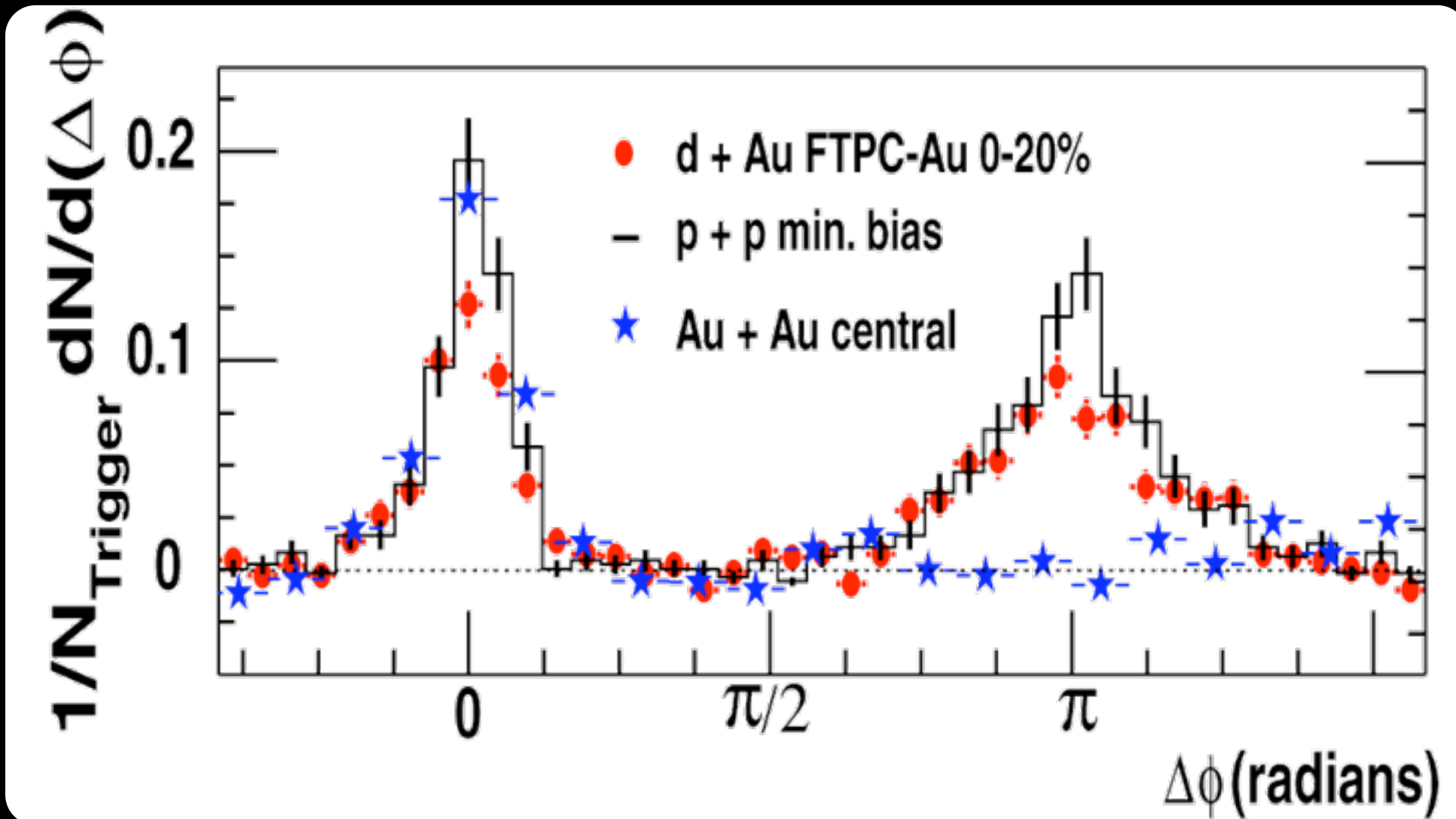
d+Au

ratio is larger than
unity in d+Au
opposite to
Au+Au!

suppression is final
state medium
effect!

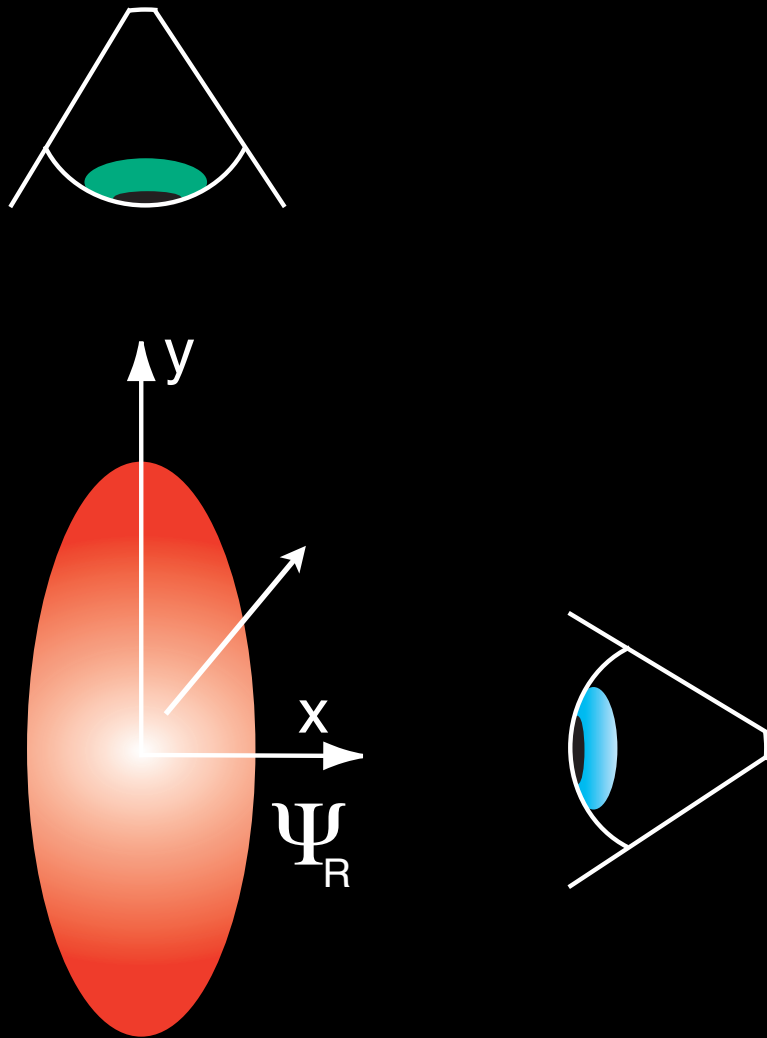


d+Au

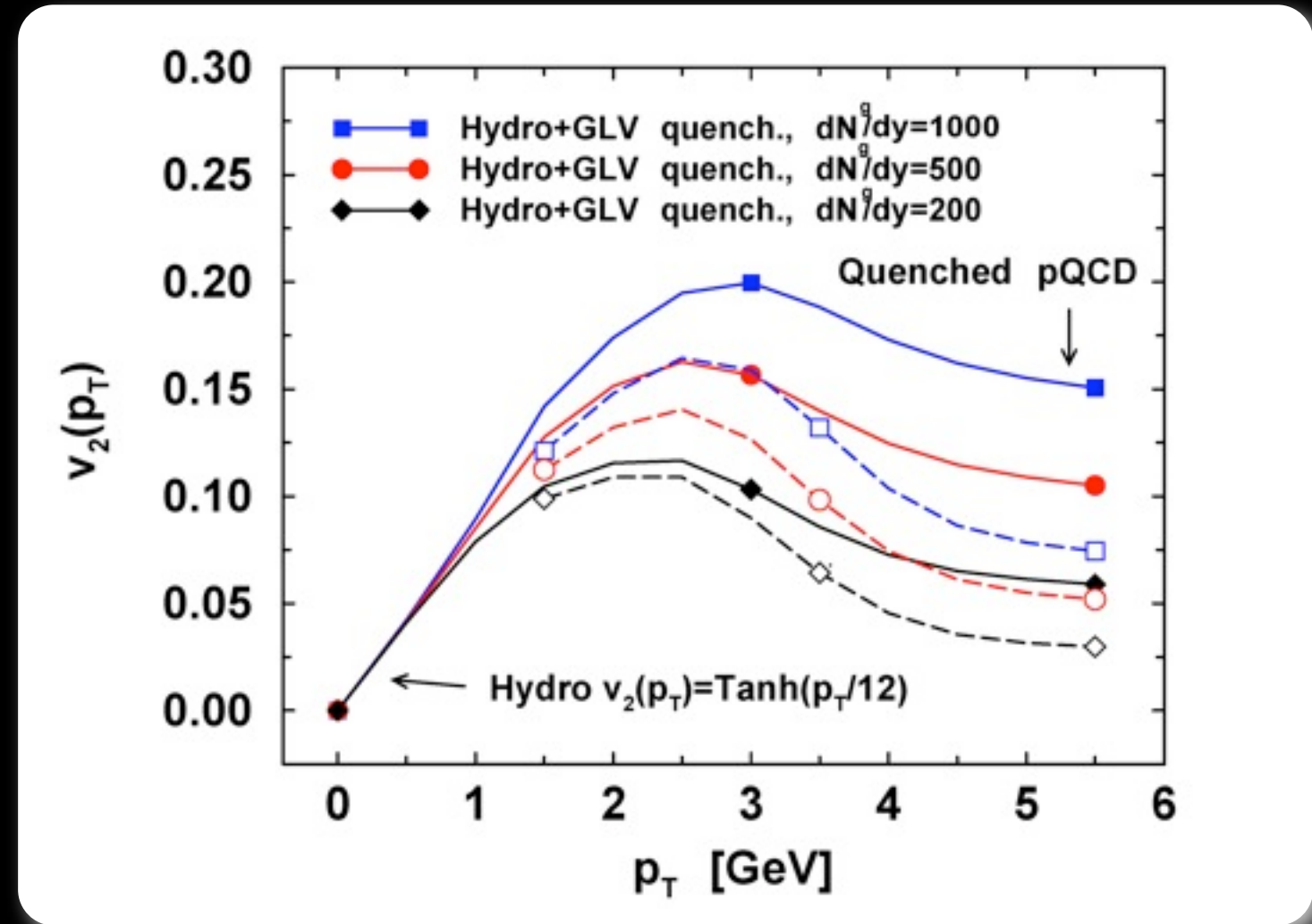


correlation in d+Au resembles pp and very different from Au+Au!
suppression is final state medium effect!

parton energy loss



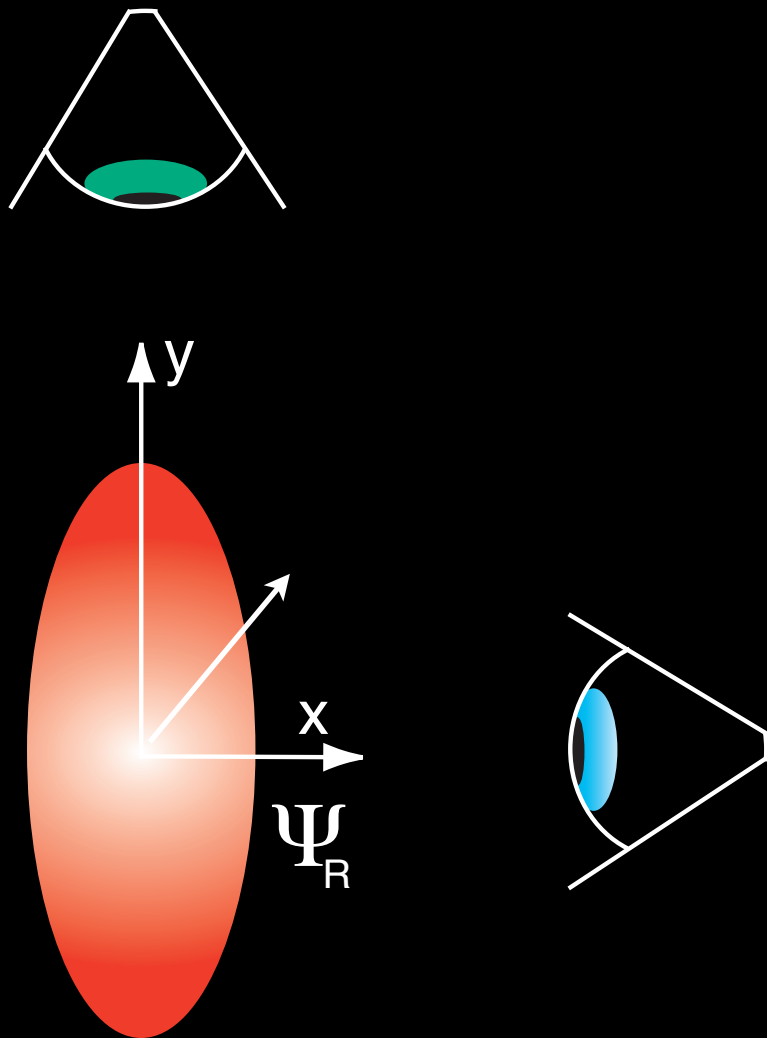
$$v_2 = \langle \cos 2(\phi - \Psi_R) \rangle$$



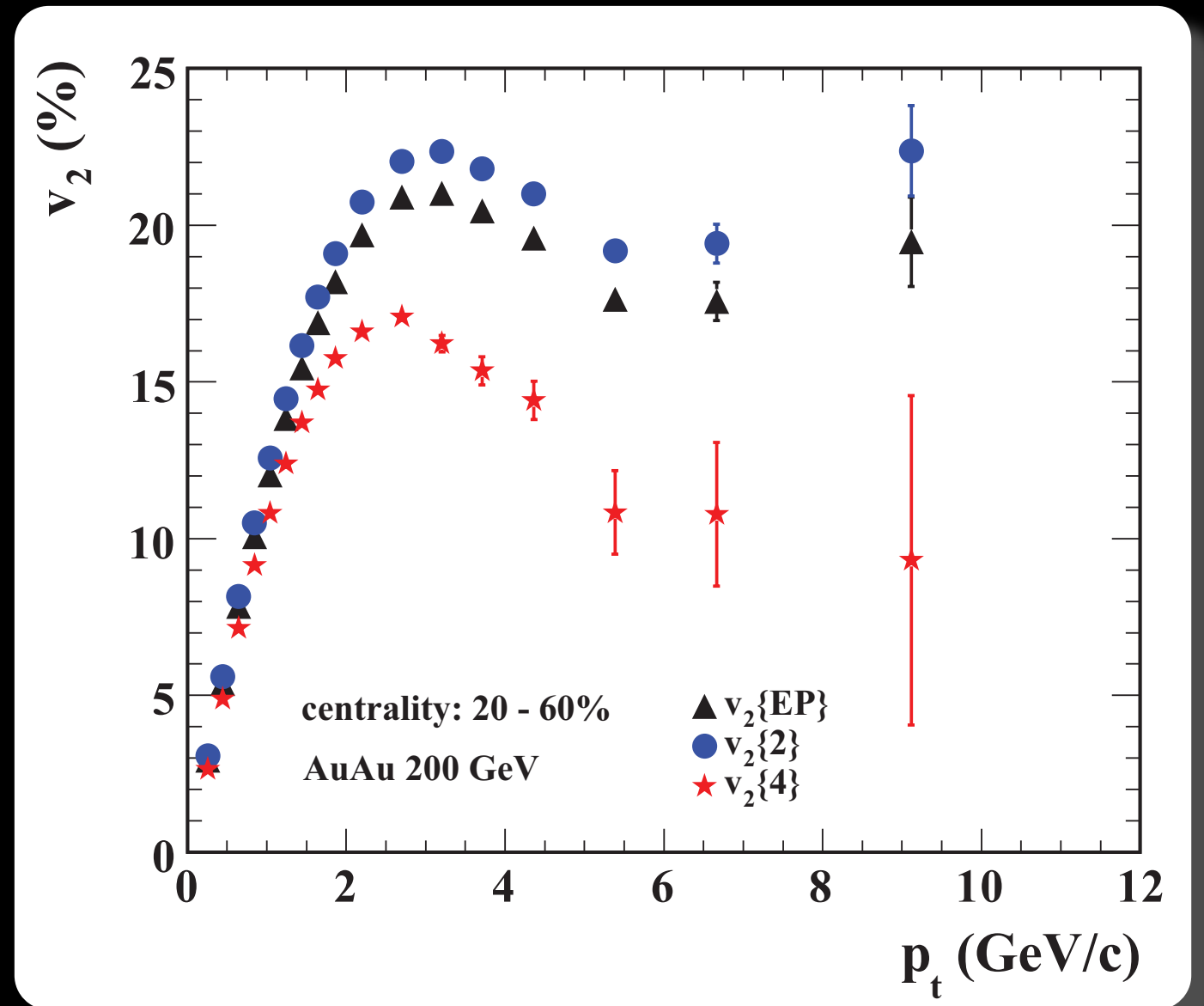
M. Gyulassy, I. Vitev and X.N. Wang
PRL 86 (2001) 2537

R.S, A.M. Poskanzer, S.A. Voloshin,
nucl-ex/9904003

parton energy loss



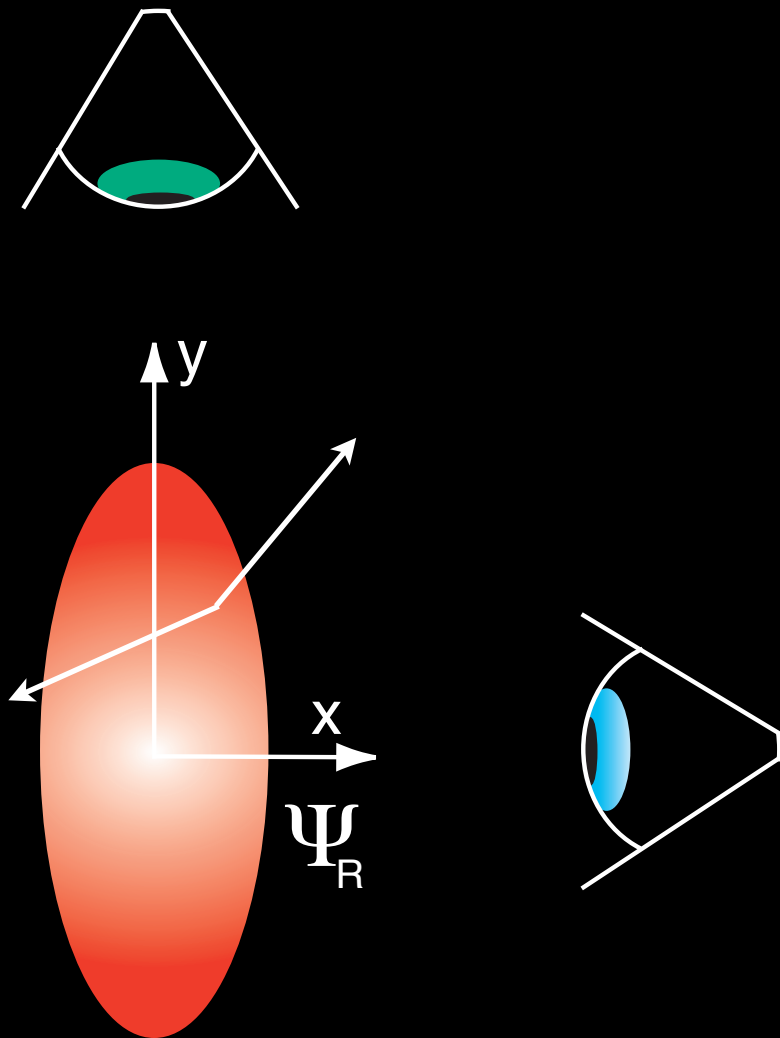
$$v_2 = \langle \cos 2(\phi - \Psi_R) \rangle$$



Yuting Bai, Nikhef PhD thesis

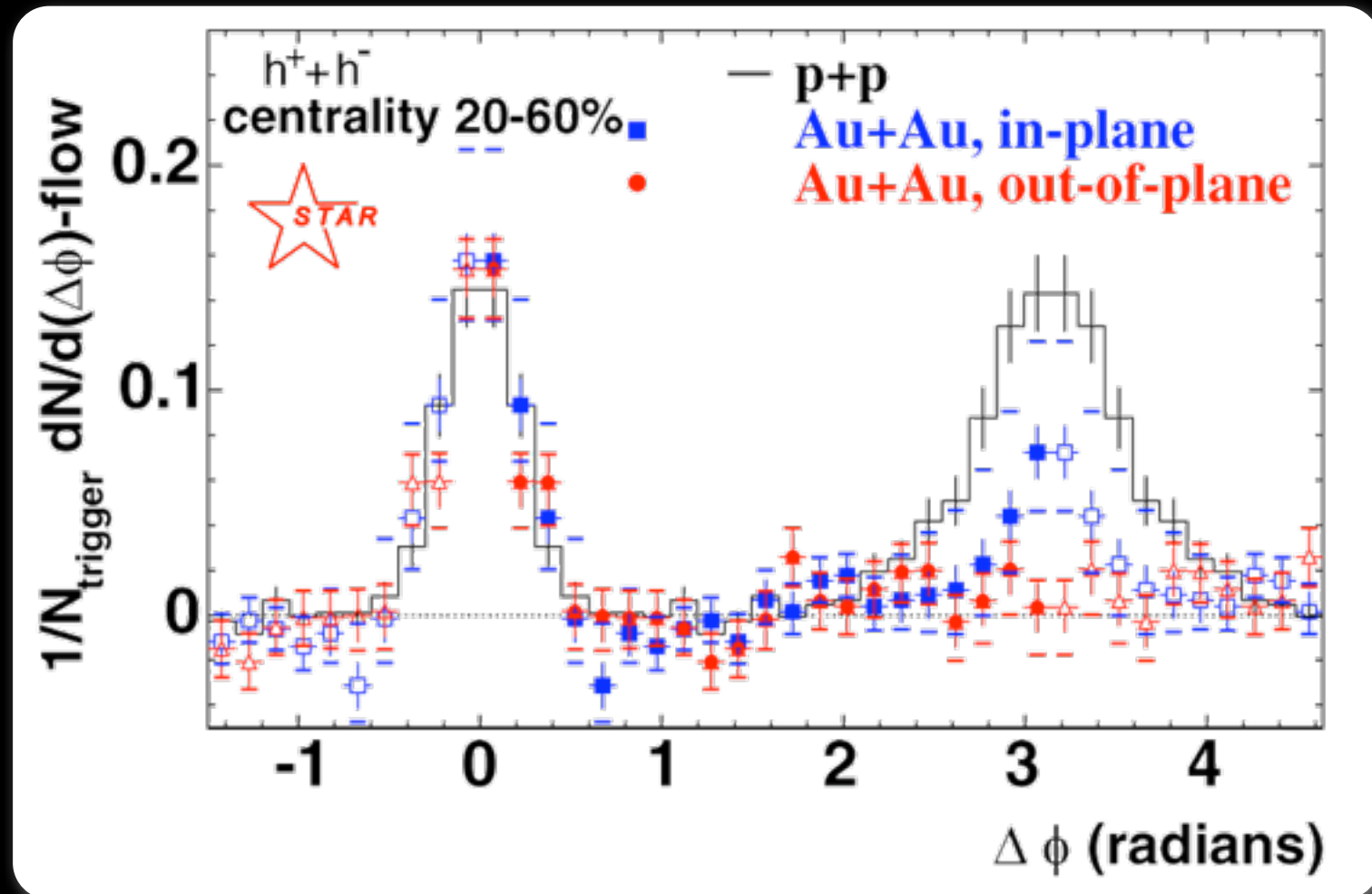
strong path length dependence observed!

parton energy loss



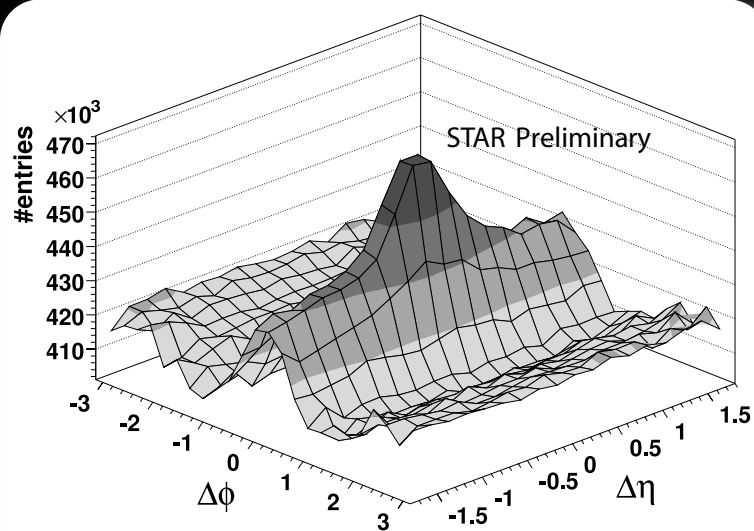
$$v_2 = \langle \cos 2(\phi - \Psi_R) \rangle$$

path length dependence observed!

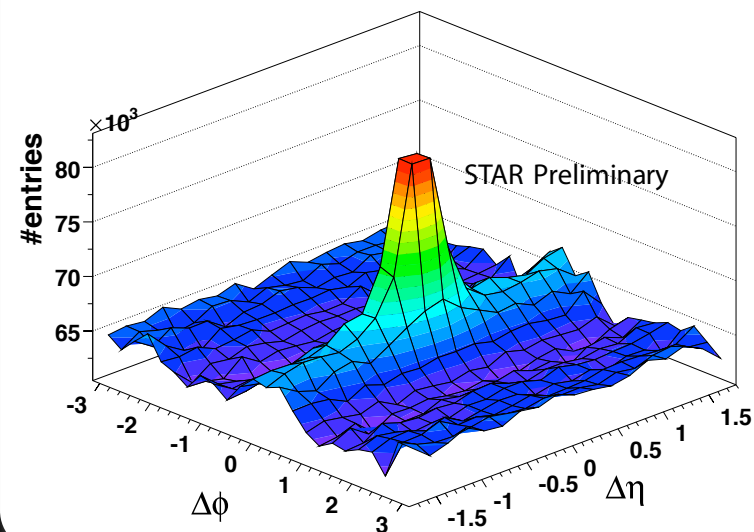


simple? the ridge

$$2 < p_t^{\text{asso}} < p_t^{\text{trigg}}$$

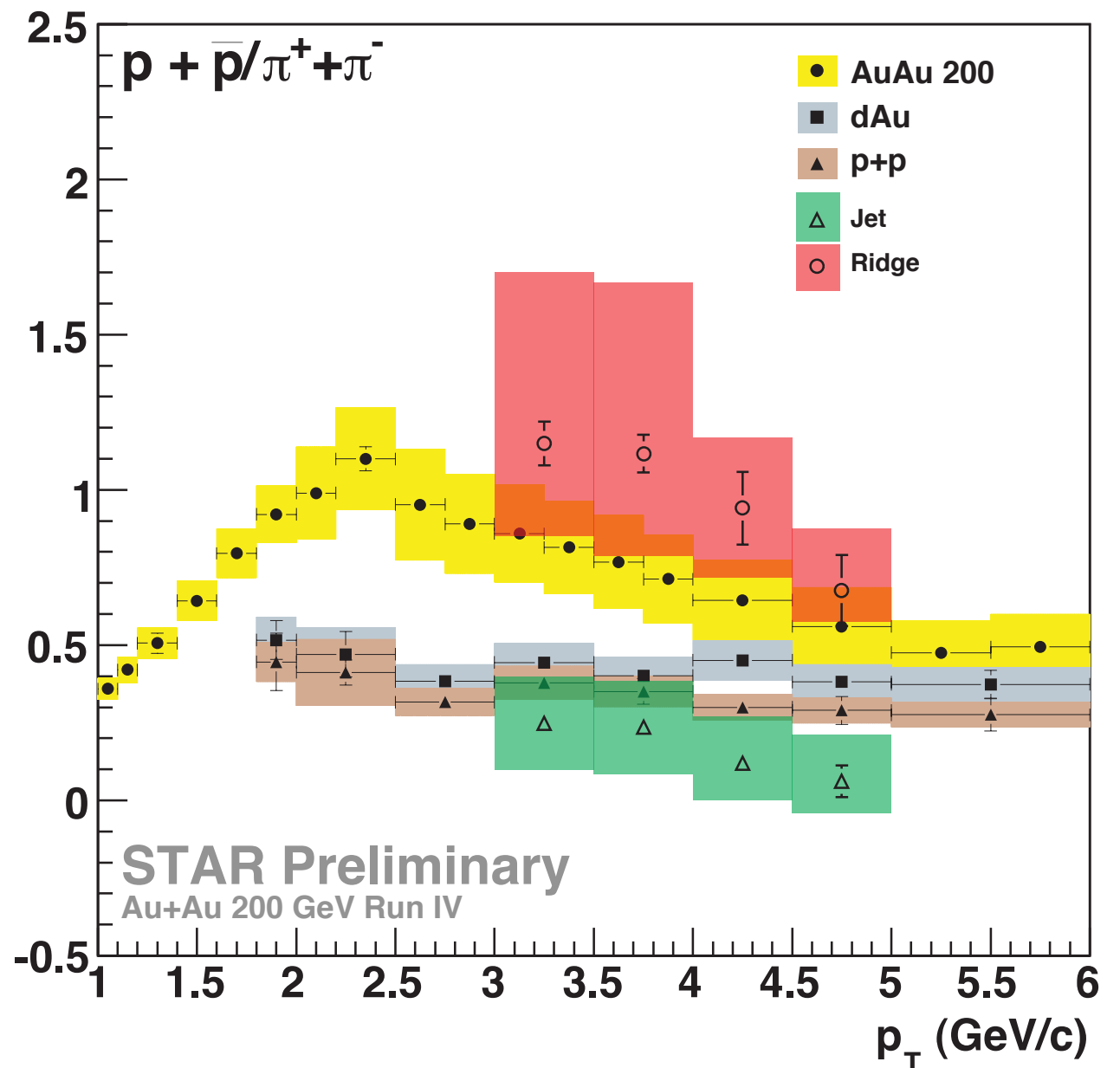


$$3 < p_t^{\text{trig}} < 4 \text{ GeV/c}$$



$$4 < p_t^{\text{trig}} < 6 \text{ GeV/c}$$

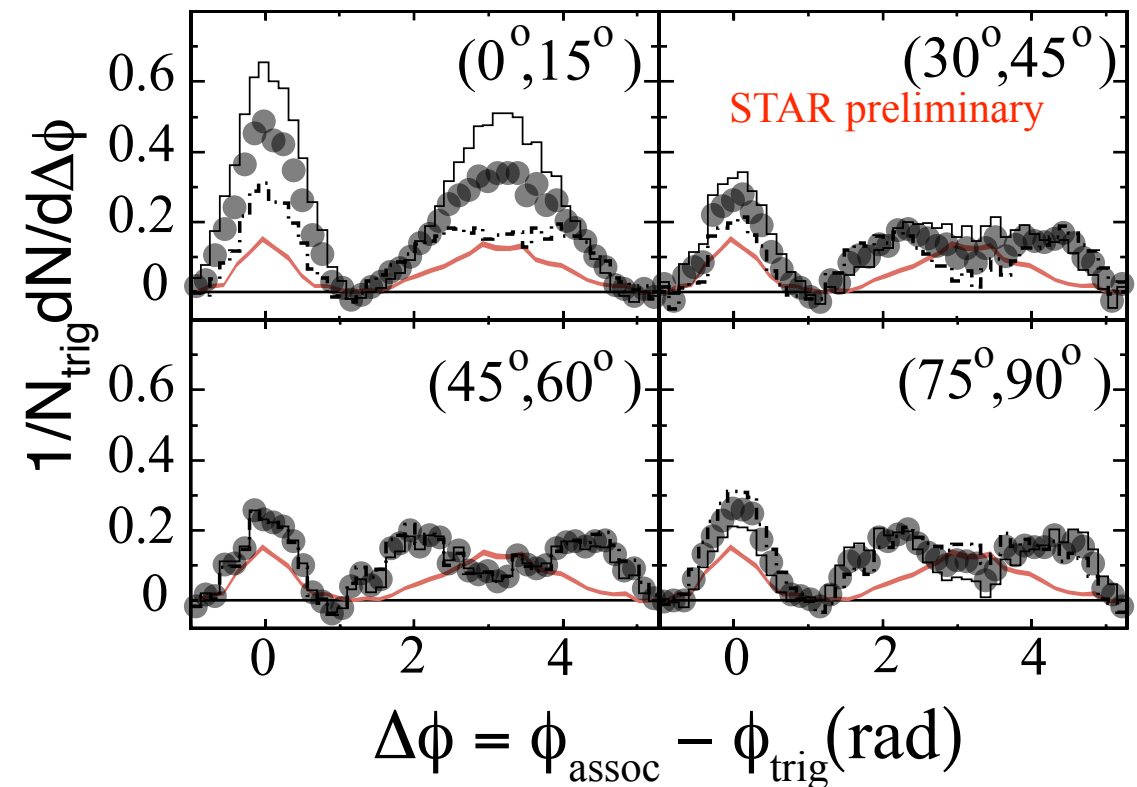
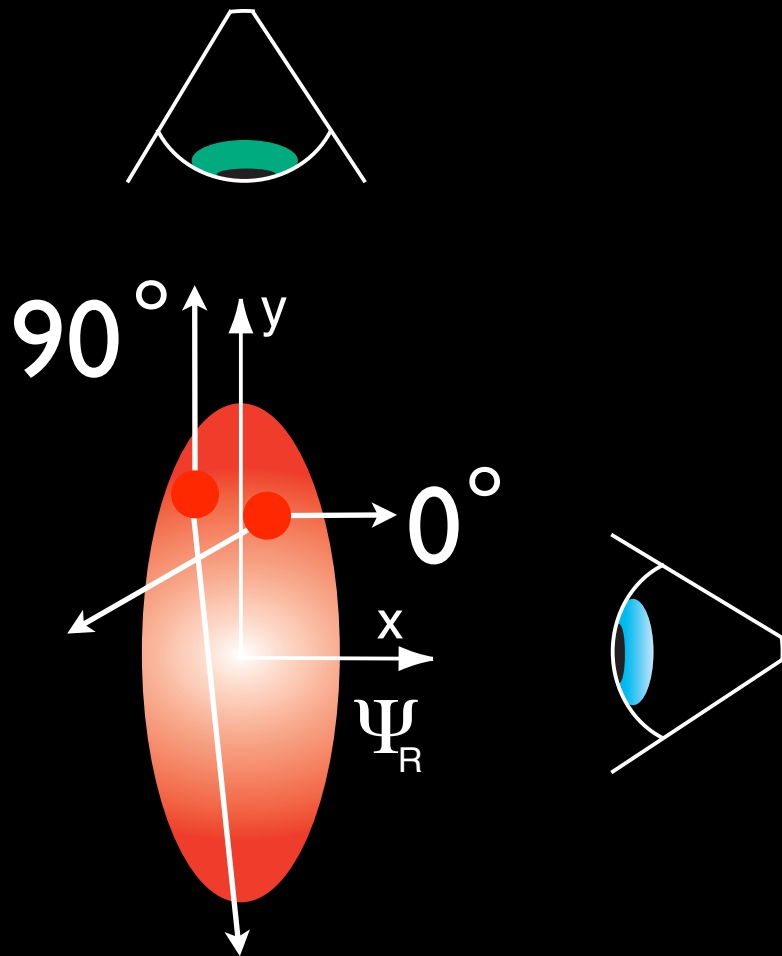
in Au+Au a “ridge” of correlated yield with the trigger particle



the ridge has the same properties as the bulk

simple? the away side

$$3 < p_t^{\text{trig}} < 4 \text{ GeV}/c \quad 1.0 < p_t^{\text{asso}} < 1.5 \text{ GeV}/c$$



Aoki Feng

The away side changes from single peak (in plane) to double peak (out of plane)

strongly debated measurement and interpretation
could provide additional acces to the sound speed

high p_t probes in $A+A$

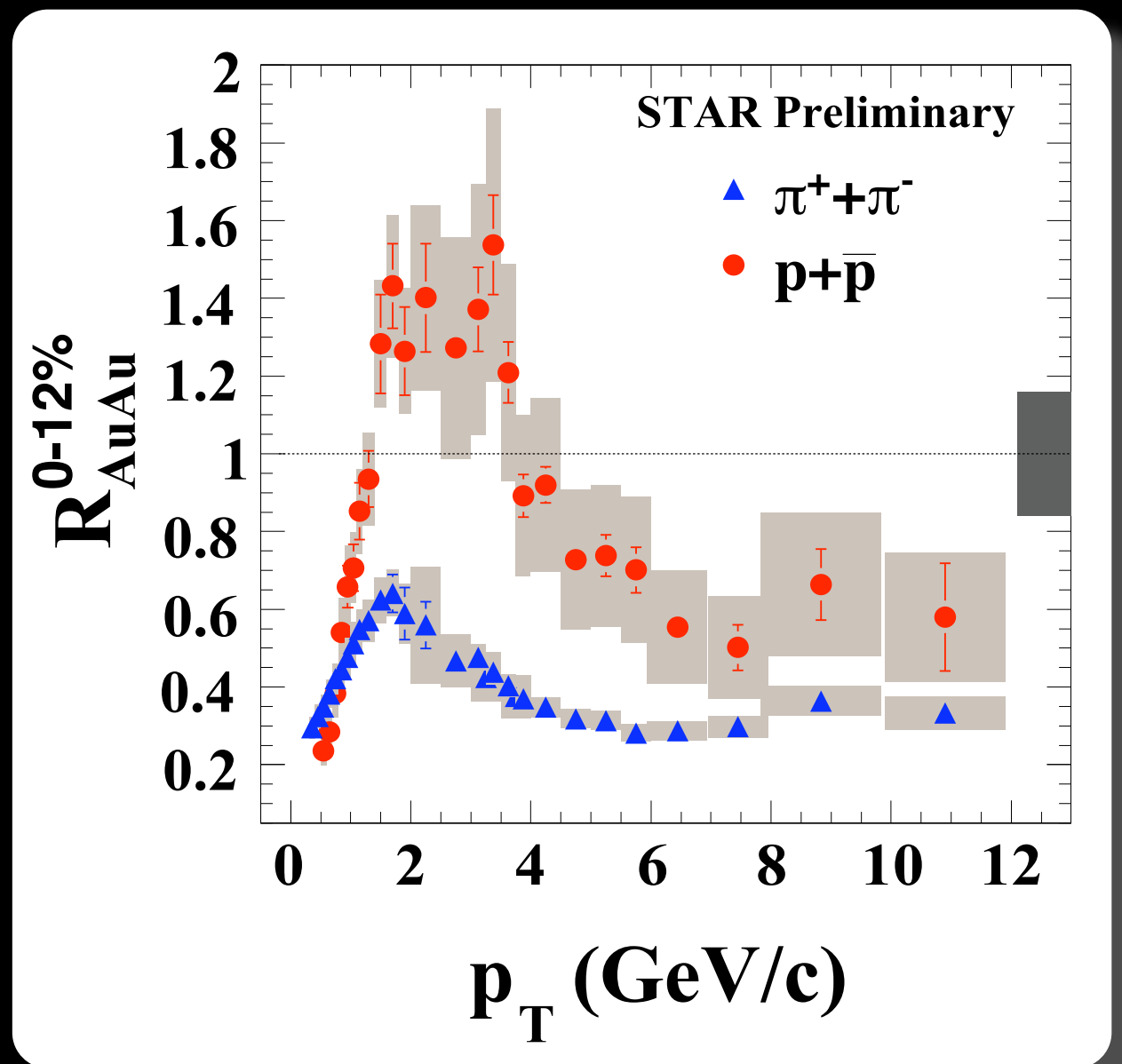
- source rather well understood (with $p+p$ and $p+A$ data available at the same energy)
- the probes clearly strongly interact with the medium
- to calibrate the interaction of the probes with the medium is the main part of the ongoing research
 - color factors
 - heavy quarks

Color Factors?

$$\langle \Delta E \rangle \propto \alpha_s C \langle \hat{q} \rangle L^2$$

$$\frac{\Delta E_g}{\Delta E_q} \approx \frac{9}{4}$$

protons produced more
from gluons, pions
produced more from
quarks: **expect stronger
suppression protons**



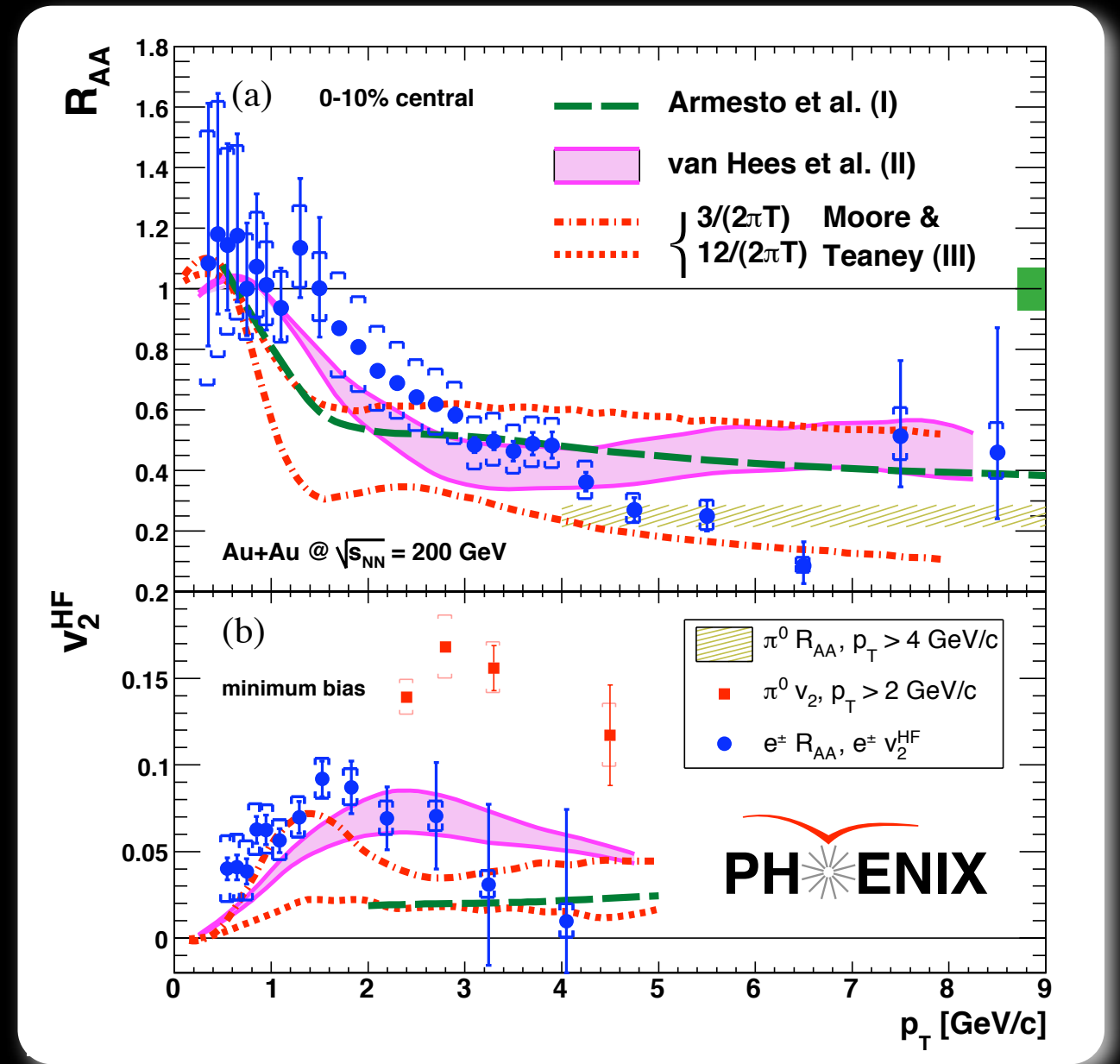
color factor effects of the magnitude 9/4 are
not observed!

Dead Cone?

probability

$$\frac{1}{[\Theta^2 + (m_Q/E_Q)^2]^2} \propto$$

in medium dead cone
implies less energy loss for
heavy quarks



PHENIX [Phys. Rev. Lett. 98, 172301 \(2007\)](#)

heavy quarks show similar suppression as light quarks!

RHIC: current summary

- A tremendous amount of striking observables at RHIC
- Leading to surprising conclusions about the matter created at RHIC
 - sQGP, perfect liquid
- Measurements and theory becoming more quantitative
 - flow: viscous corrections are addressed
 - high- p_t : testing energy loss mechanism, starting full jet reconstruction

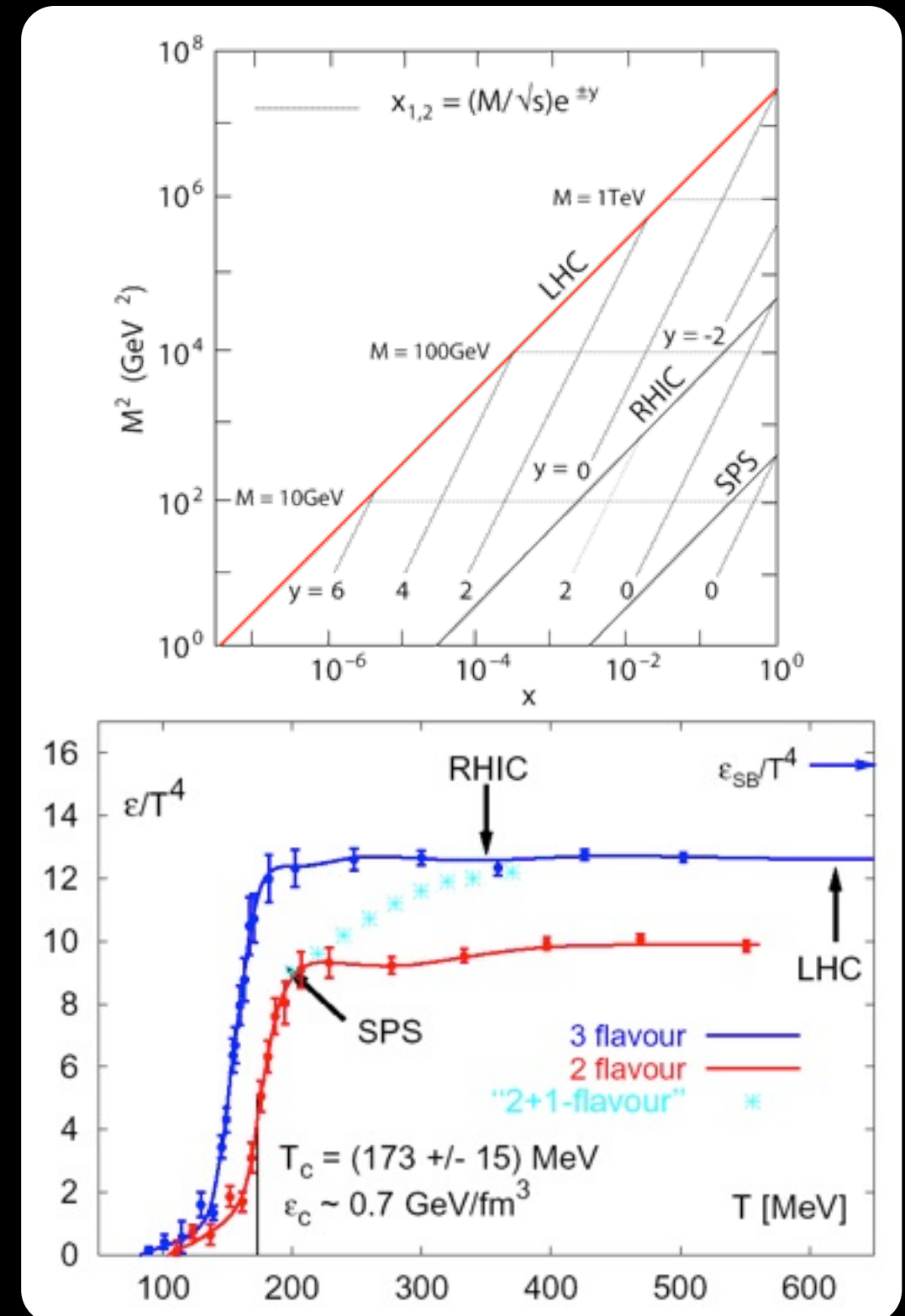
From SPS, RHIC to the LHC

	SPS	RHIC	LHC	
$\sqrt{s_{NN}}$ (GeV)	17	200	5500	
dN/dy	500	850	1500-4000	
τ_{QGP}^0 (fm/c)	1	0.2	0.1	
T/T_c	1.1	1.9	3-4	Hotter
ε (GeV/fm ³)	3	5	15-60	Denser
τ_{QGP} (fm/c)	≤ 2	2-4	≥ 10	Longer
τ_f (fm/c)	~ 10	20-30	30-40	
V_f (fm ³)	few 10^3	few 10^4	Few 10^5	Bigger



From SPS, RHIC to the LHC

- not just super sized, a new regime!
 - high density pdf's (saturised) determine particle production
 - parton dynamics dominate the fireball expansion
- new tools!
 - hard processes contribute significantly to the cross section
 - weakly interacting hard probes become available
- detailed understanding of QGP
- possible surprises!



Jets at the LHC

- At LHC >90% of the particle production from hard collisions, jet rates are high at energies at which jets can be reconstructed over the large background from the underlying event
- More than 1 jet > 20 GeV per central collision (more than 100 > 2 GeV!)
- Reach to about 200 GeV
- Provides lever arm to measure the energy dependence of the medium induced energy loss

1 month of running	
$E_T >$	N_{jets}
50 GeV	2.0×10^7
100 GeV	1.1×10^6
150 GeV	1.6×10^5
200 GeV	4.0×10^4

penetrating probes: heavy quarks

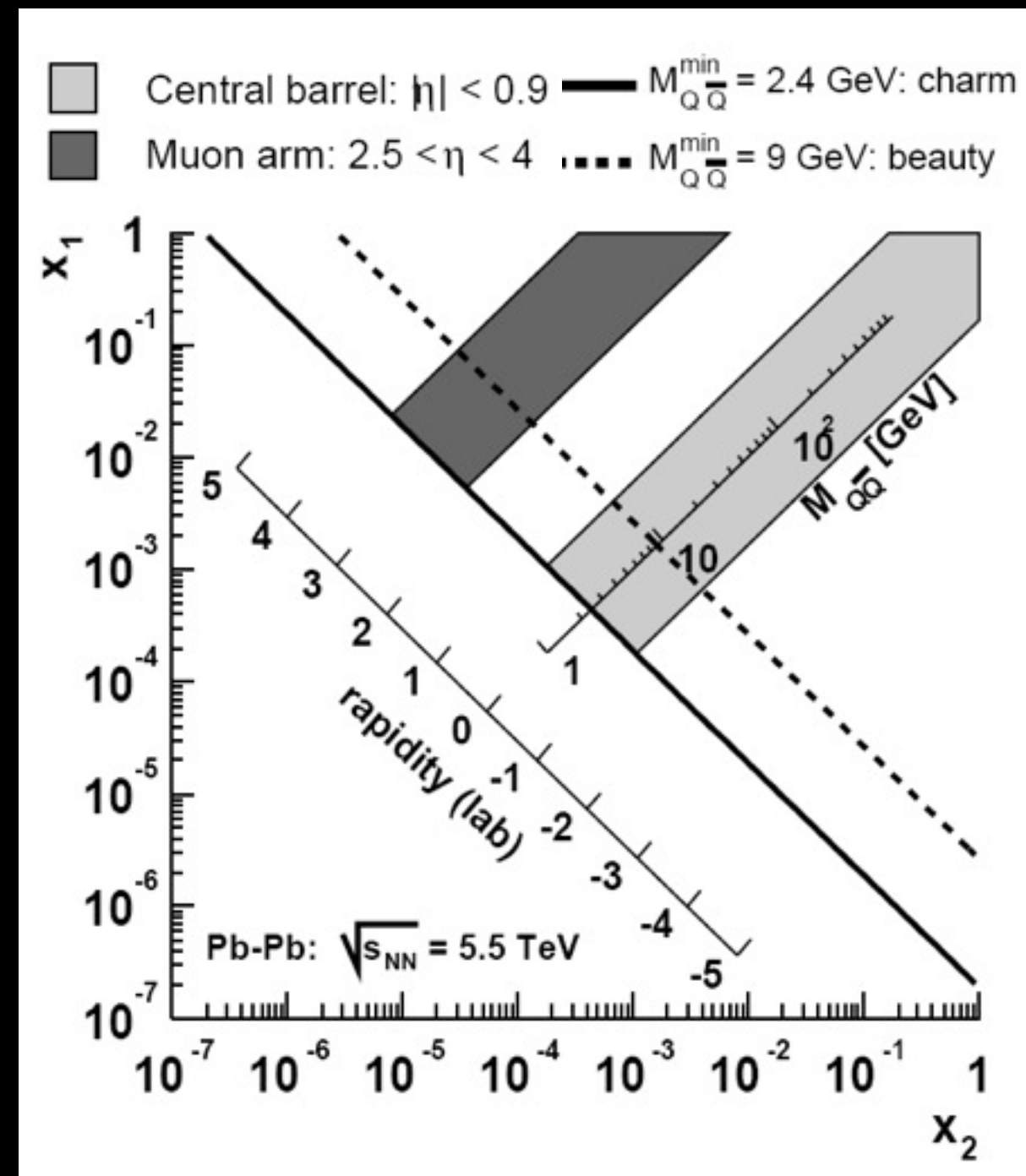
	SPS PbPb Cent	RHIC AuAu Cent	LHC pp	LHC pPb	LHC PbPb Cent
N_{cc}/evt	0.2	10	0.2	1	115
N_{bb}/evt	-	0.05	0.007	0.03	5

- produced early and calculable: $\tau \propto 1/m_Q$
- Relatively long lifetime: $\tau_{\text{decay}} \gg \tau_{QGP}$
- detailed test of parton energy loss
 - dead cone effect

Probability:

$$\propto \frac{1}{[\theta^2 + (m_Q / E_Q)^2]^2}$$

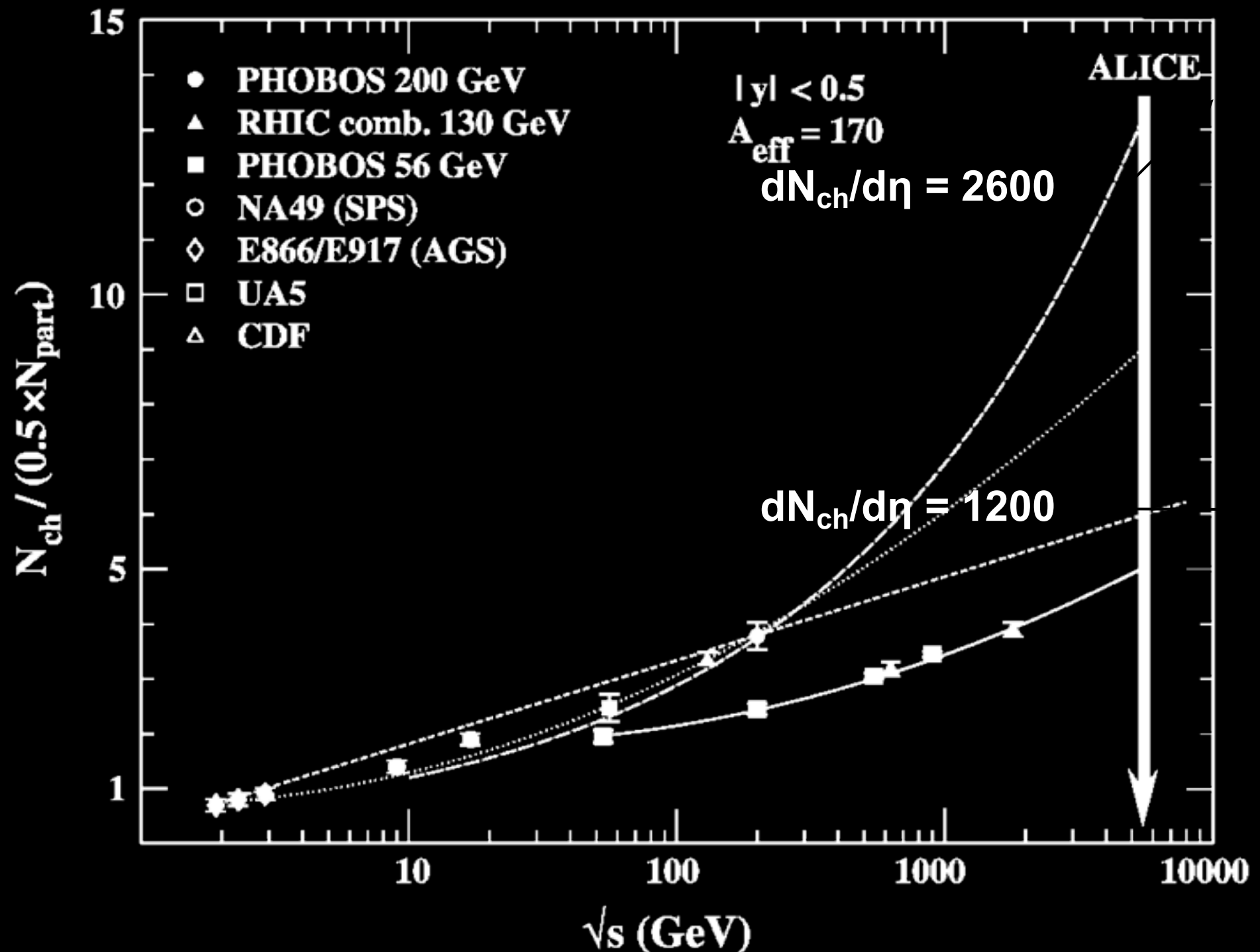
- In medium dead cone implies less energy loss
- probes small x ($10^{-3} - 10^{-5}$)



particle multiplicities

large
uncertainty in
predicted
multiplicities
at the LHC

complicates
designing a
detector!



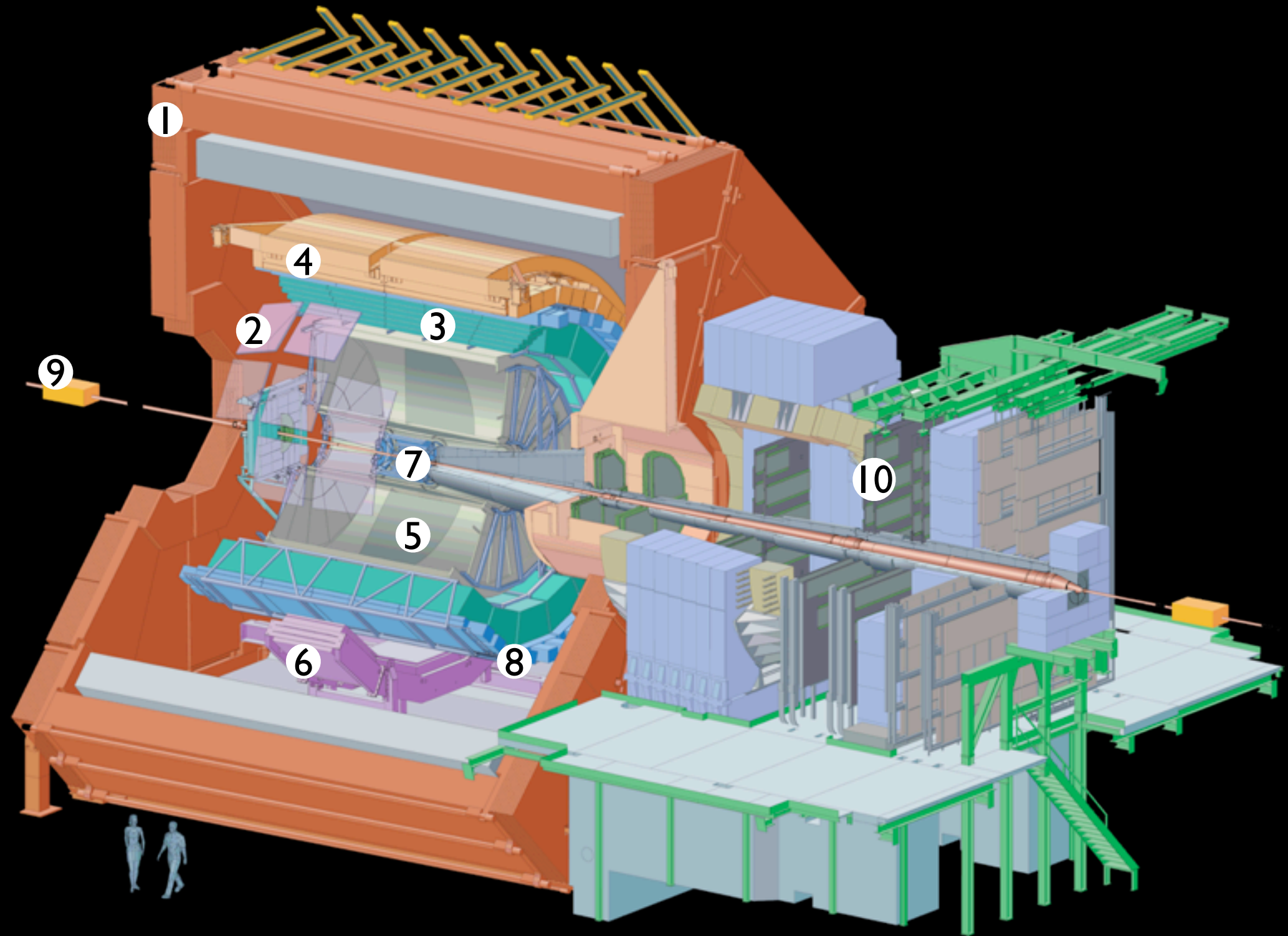


Alice Detector

- measure low- p_t particles (< 100 MeV/c)
- measure high- p_t tracks (100 GeV/c)
- particle identification over this range
- able to measure large multiplicities (4000 per unit rapidity)
- measure rare probes (direct γ , charm, bottom)

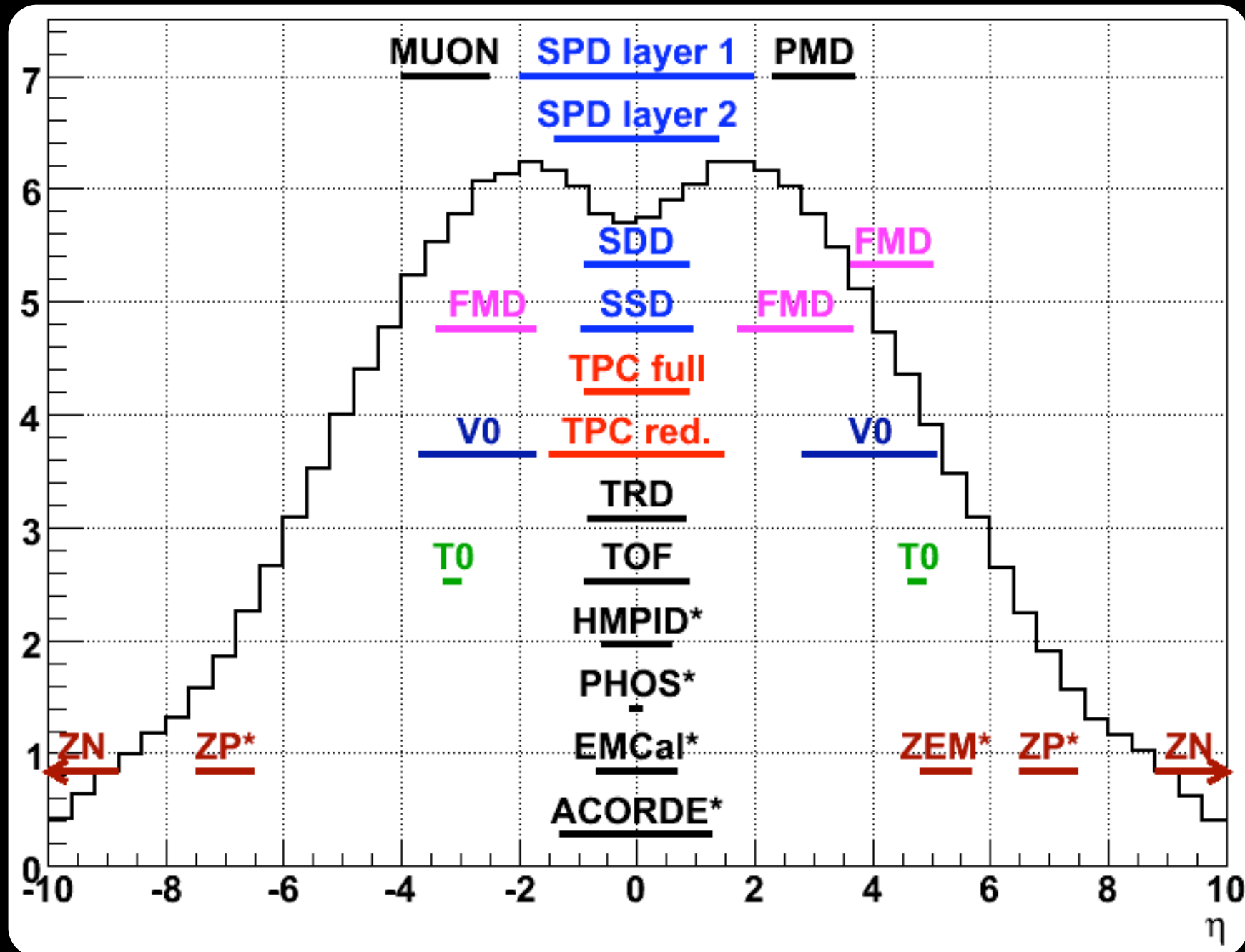
ALICE at the LHC

- 1. L3 magnet
- 2. HMPID
- 3. TRD
- 4. EMCAL
- 5. TPC
- 6. PHOS
- 7. ITS
- 8. TOF
- 9. ZDC
- 10. Muon system



~1000 collaborators from 109 institutes in 31 countries

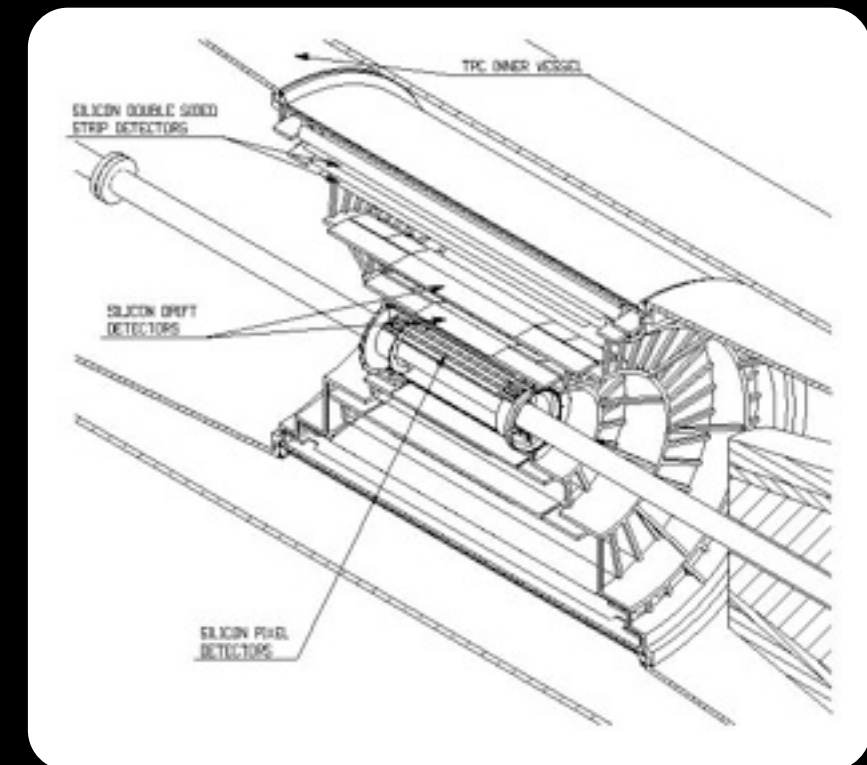
Alice Detector



Inner Tracker System

low mass: 8 % X_0

6 layers		R	$\sigma_{r\phi}$	σ_Z
Layer 1	pixels	4 cm	12 μm	100 μm
Layer 2	pixels	8 cm	12 μm	100 μm
Layer 3	drift	15 cm	38 μm	28 μm
Layer 4	drift	24 cm	38 μm	28 μm
Layer 5	double sided strip	38 cm	17 μm	800 μm
Layer 6	double sided strip	43 cm	17 μm	800 μm



- needed to reconstruct secondary vertices
- needed to track very low momenta particles

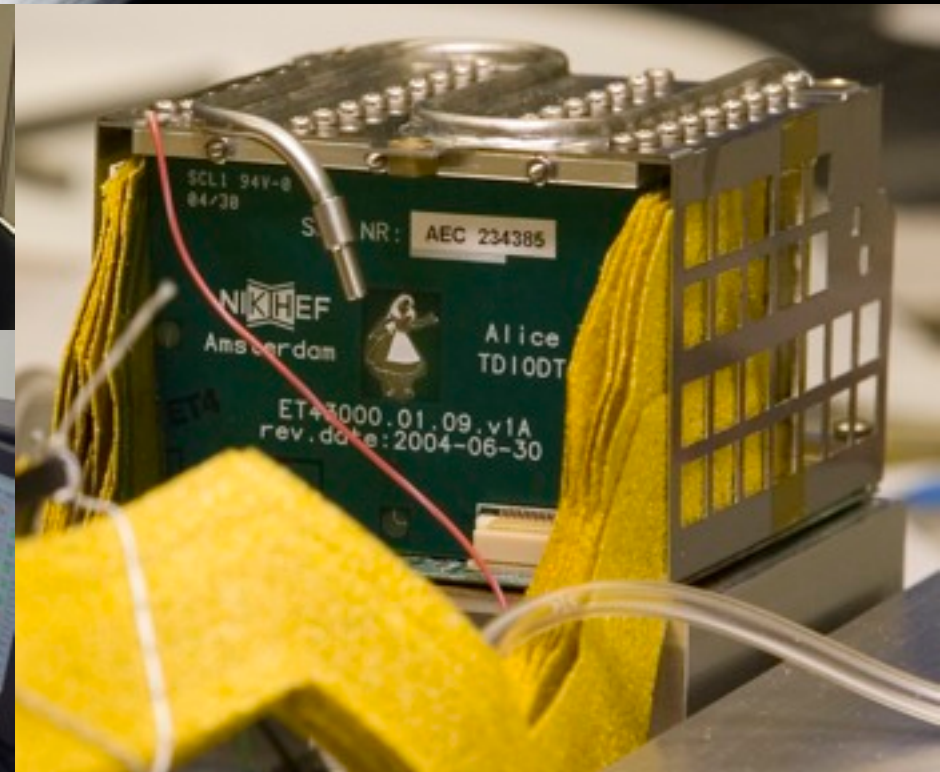
Completed SSD Ladder



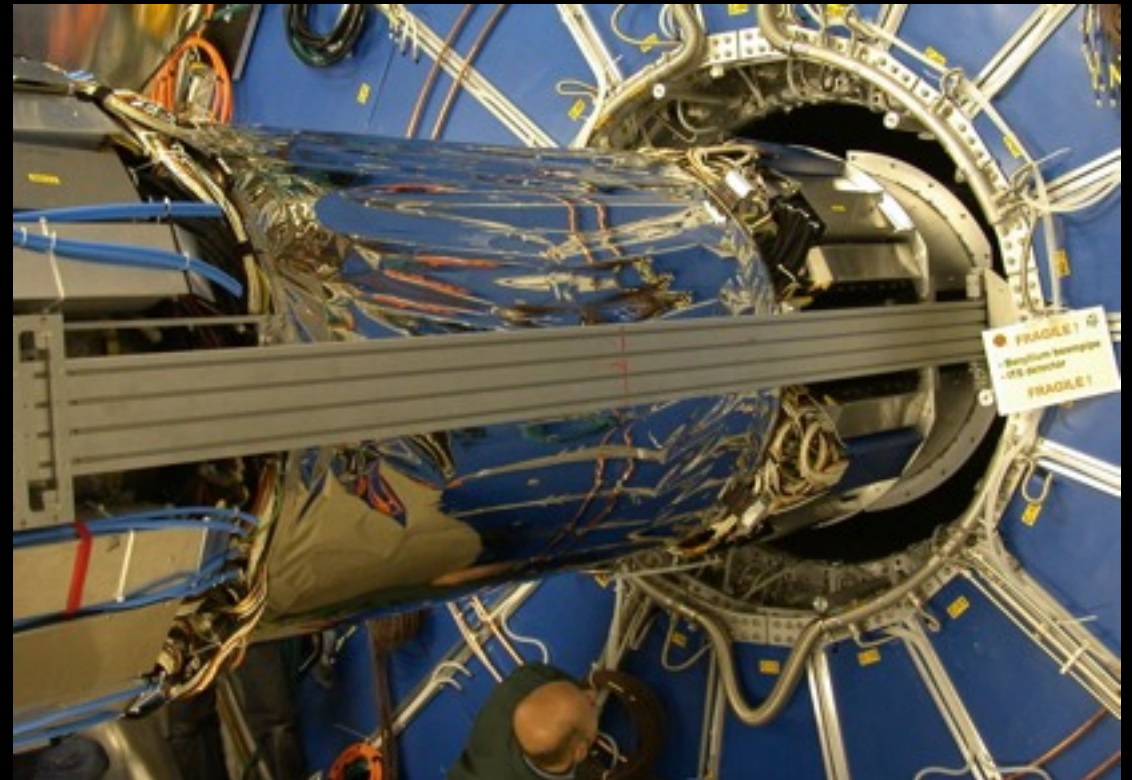
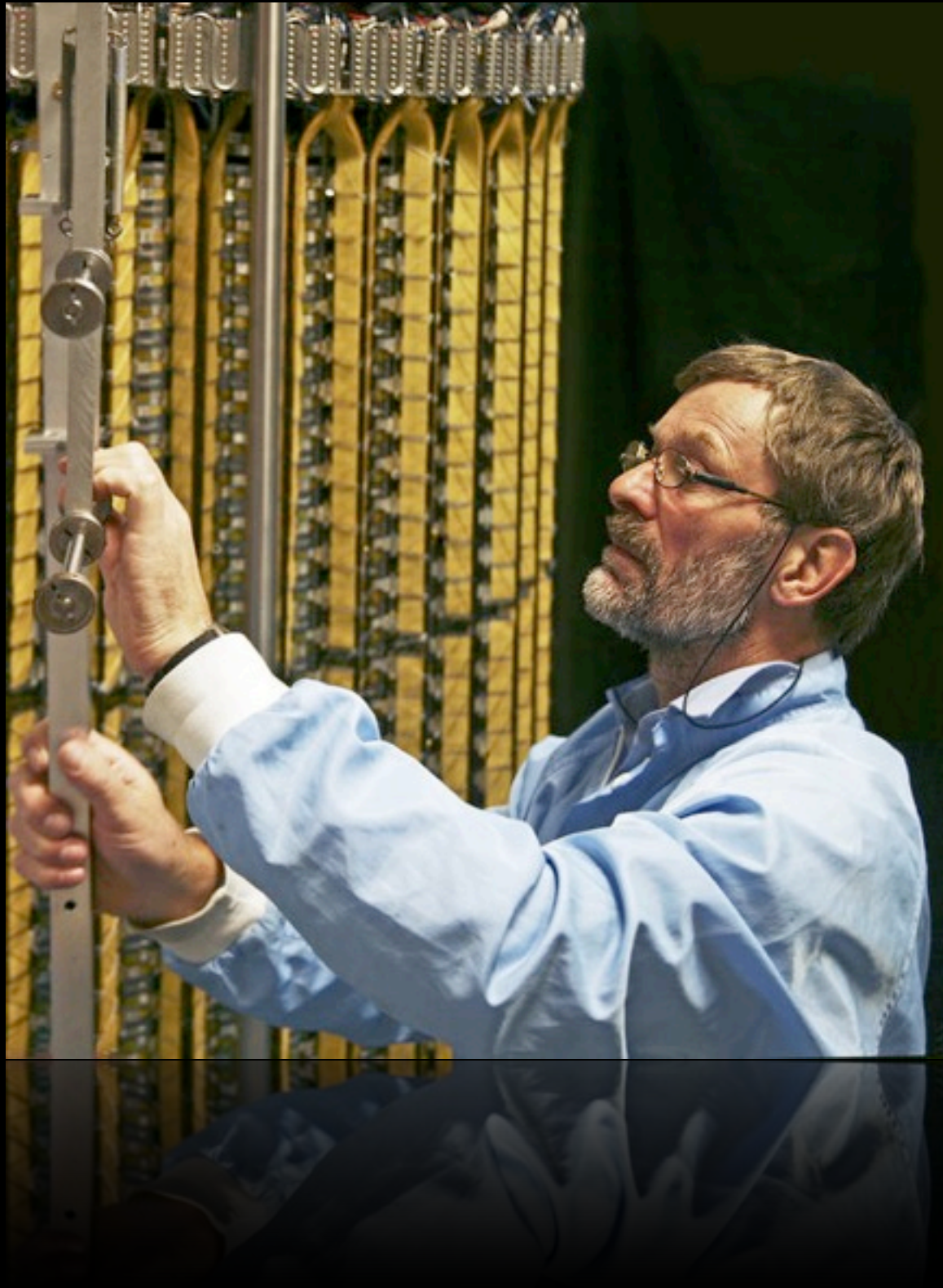
The main image shows a large, complex electronic assembly, the Completed SSD Ladder, featuring a dense grid of yellow and silver components mounted on a black frame. The assembly is shown in a perspective view, extending from the foreground into the background.

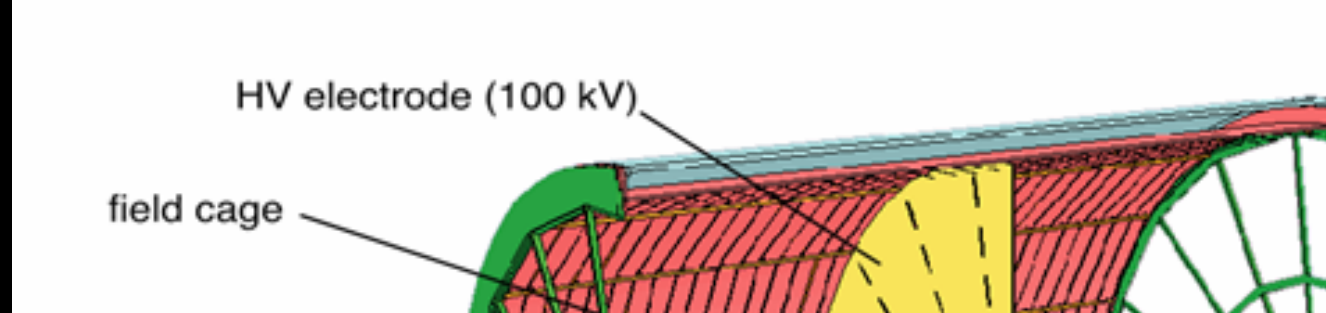
Three smaller inset images provide additional context:

- The top-right inset shows a person in a blue lab coat and hairnet working on a similar assembly in a cleanroom environment.
- The bottom-left inset shows a person in a blue lab coat and hairnet sitting at a desk with a computer monitor, likely monitoring the assembly process.
- The bottom-right inset shows a close-up of a component labeled "SCL1 94V-0 04/30" and "NR: AEC 234385" with "NIKE" visible, mounted on a green circuit board.

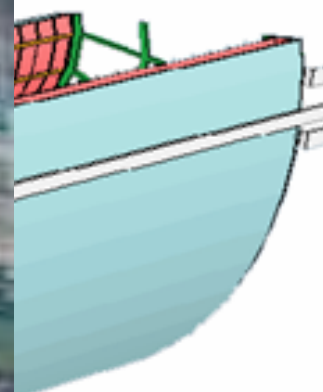


Full SSD





The
largest

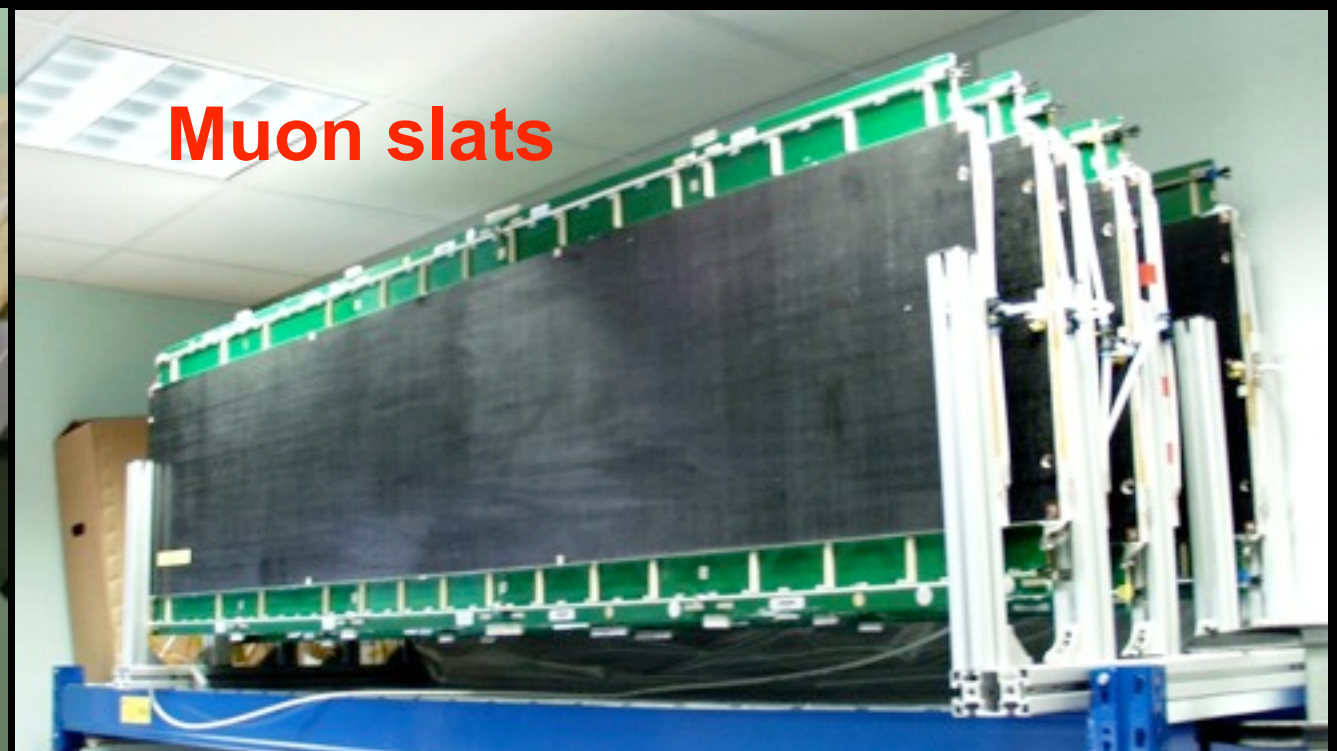
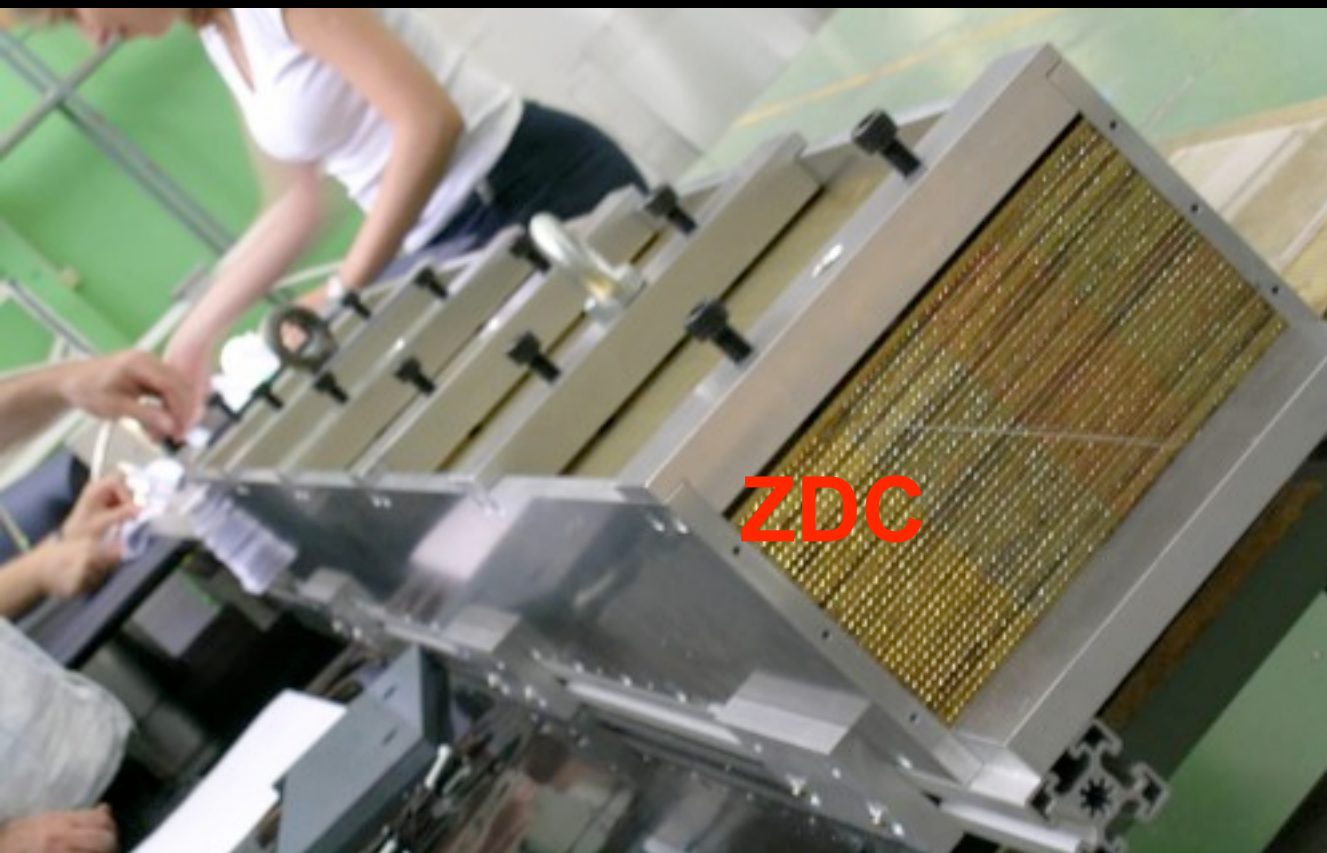


Field Cage



HV membrane (25 μm)

A few other detectors

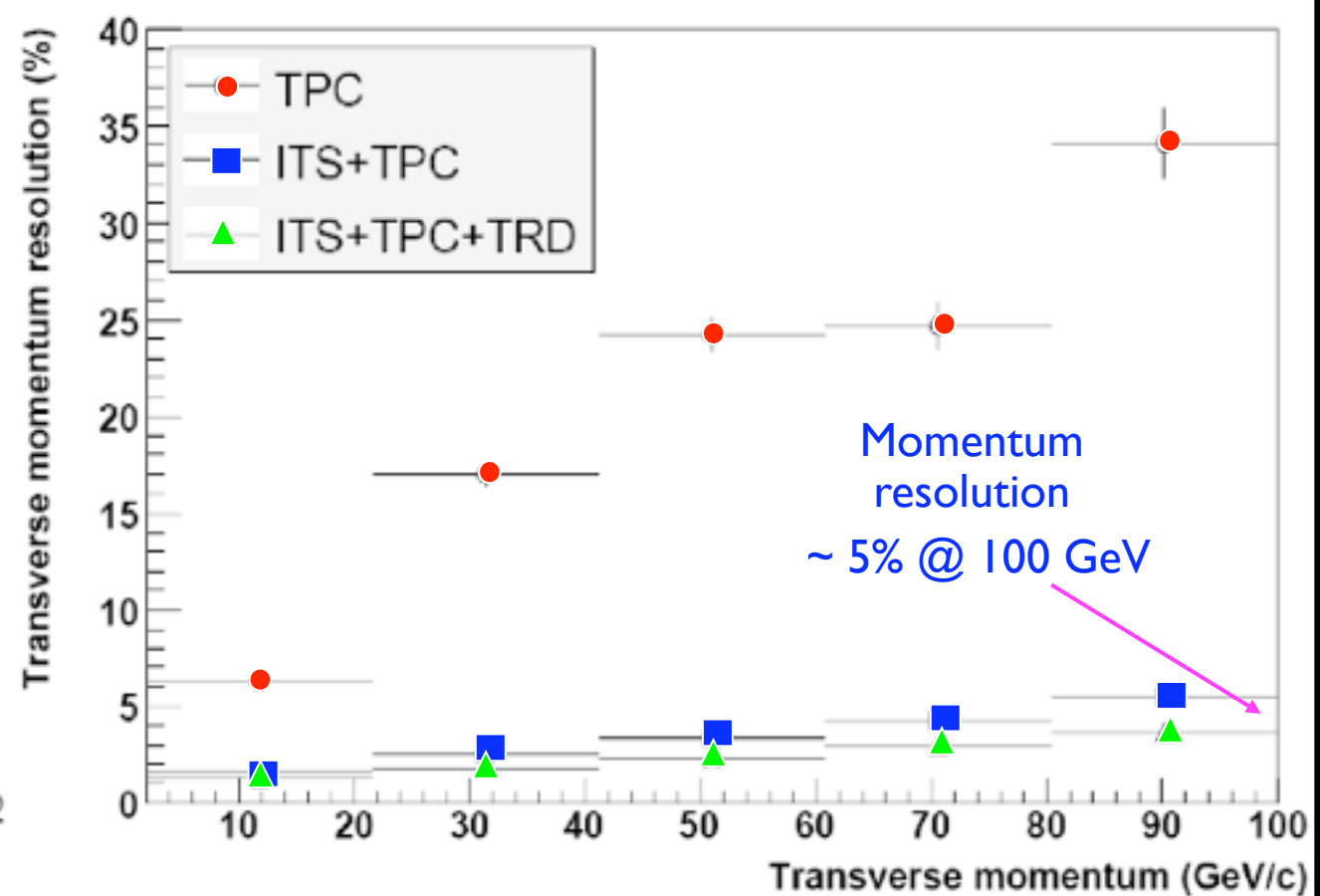
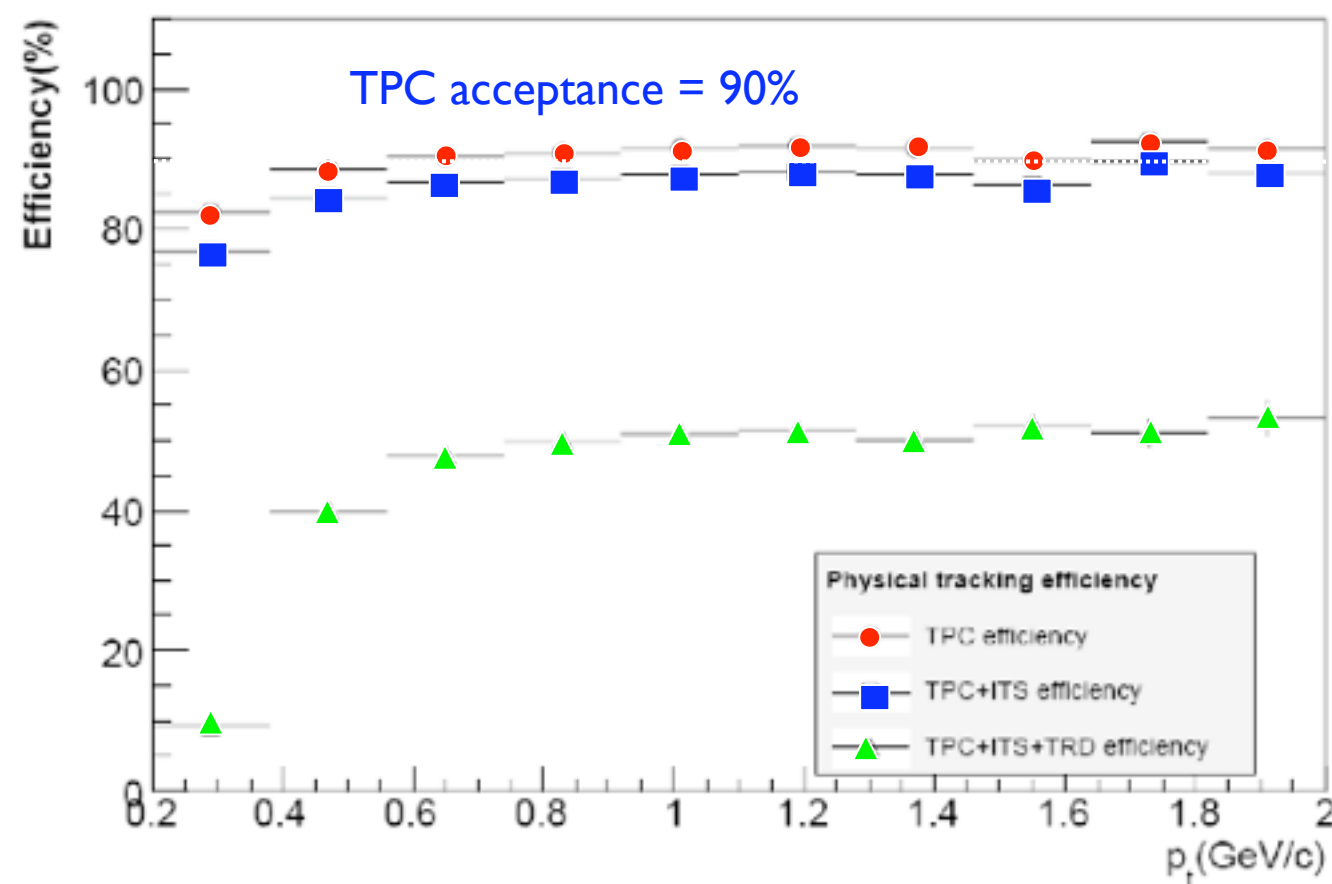


Tracking

- robust, redundant tracking from 0.1 to 100 GeV/c
 - modest solenoidal field (0.5 T) \Rightarrow easy pattern recognition
 - long lever arm \Rightarrow good momentum resolution
 - BL^2 : ALICE \sim CMS $>$ ATLAS
- small material budget \sim 10% X_0 vertex to end TPC
- Silicon Vertex Detector: 4-44 cm, 6 layers
- Time Projection chamber: 85-245 cm, 159 pad rows
- Transition Radiation Detector: 290-370 cm, 6x3 cm tracks

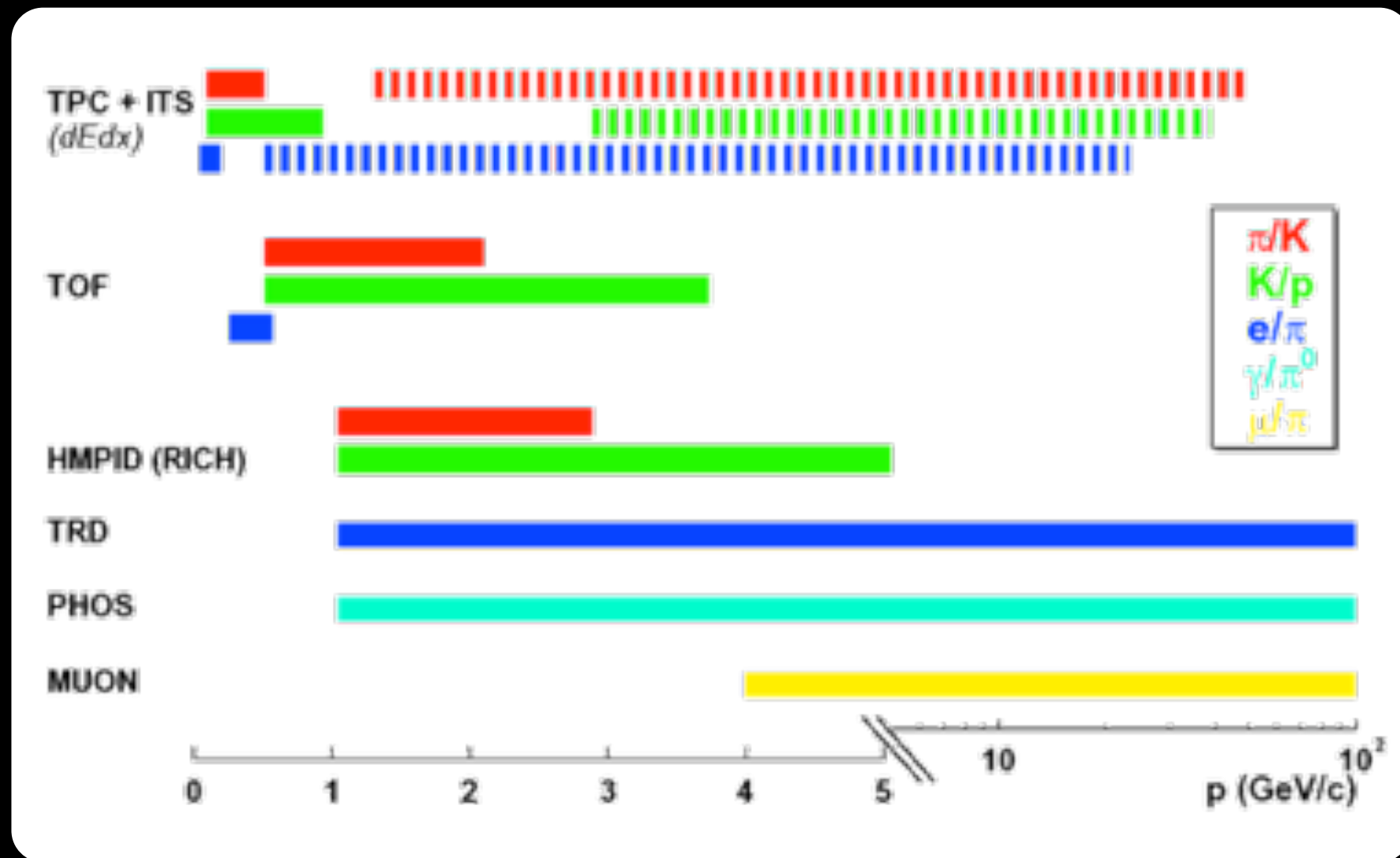
Tracking

- full GEANT simulation: central Pb-Pb, $dN_{ch}/dy = 6000$
- very little dependence on dN_{ch}/dy up to 8000 (important for systematics !)



particle identification

- techniques
 - dE/dx
 - time-of-flight
 - Cerenkov
 - transition radiation
 - calorimeter
 - spectrometer

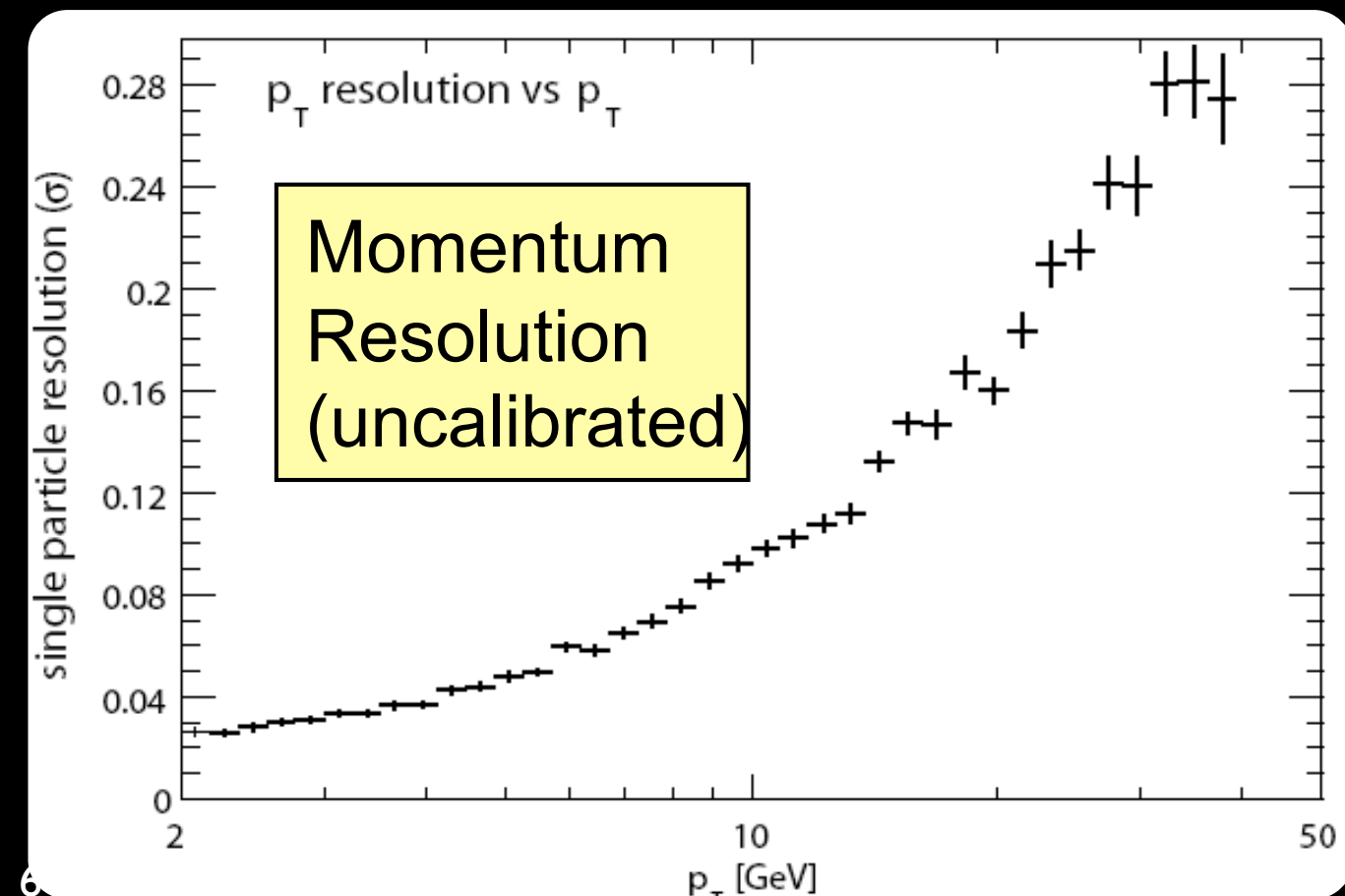
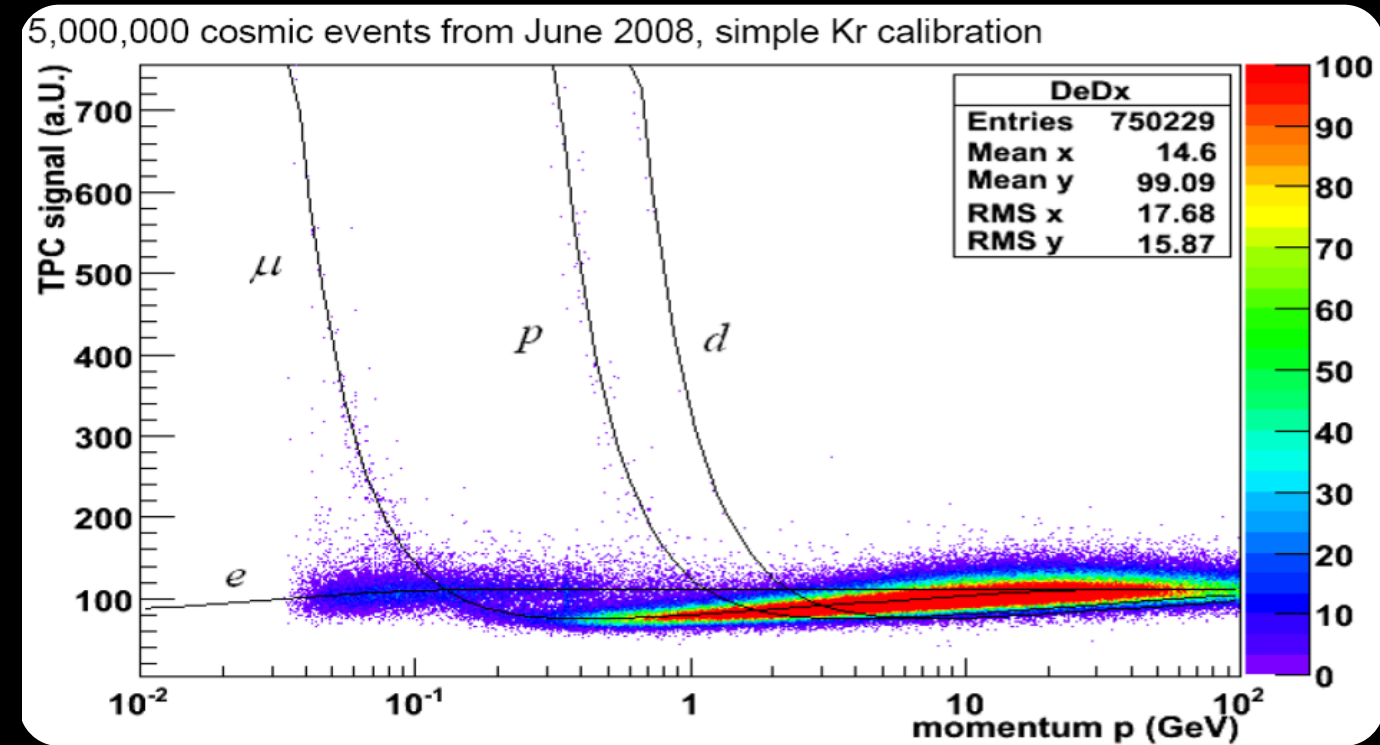
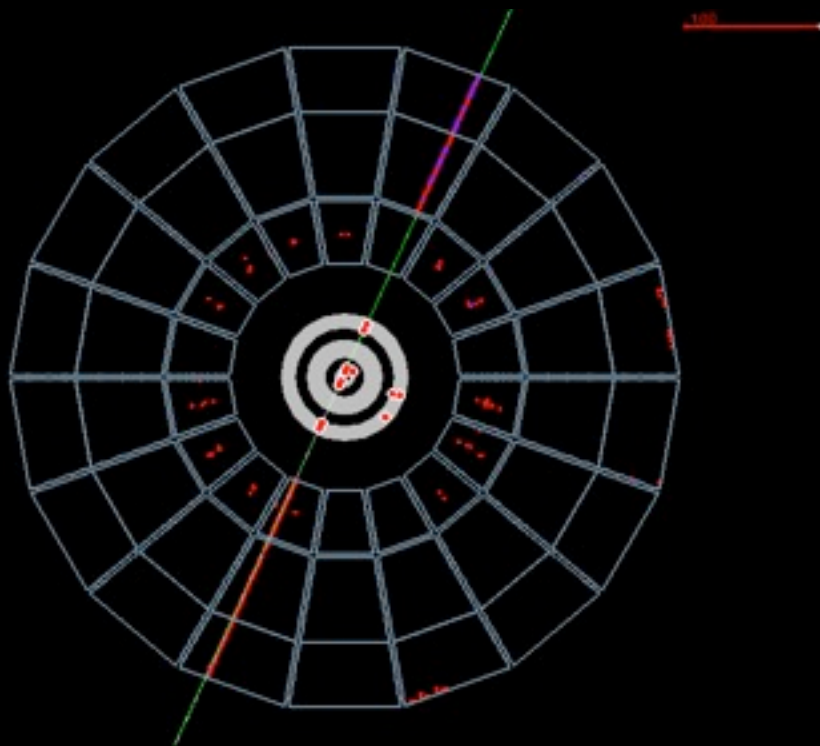


End of ALICE Installation 2008



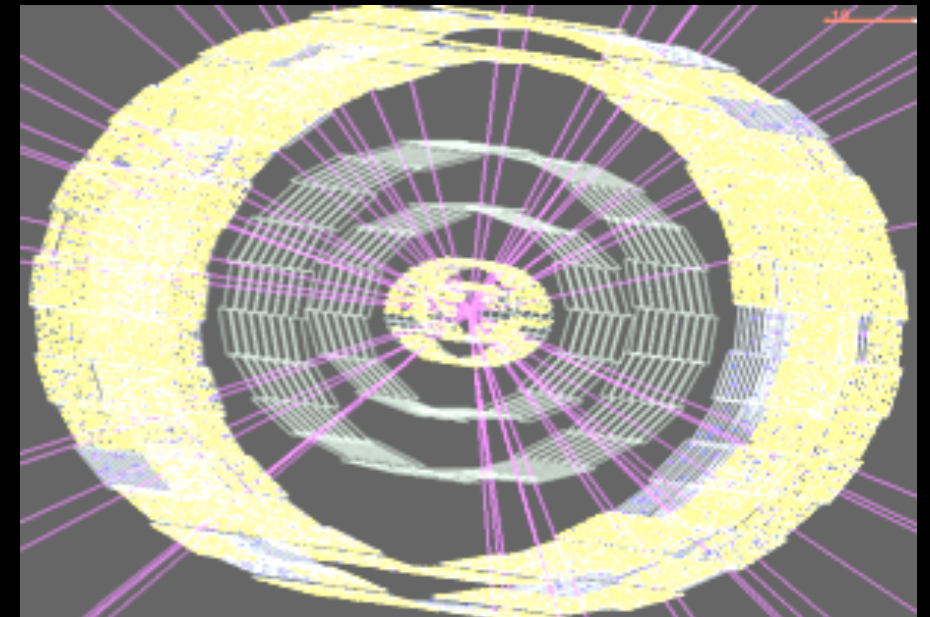
TPC Performance

- from cosmic data
 - dE/dx resolution $< 6\%$ (PPR goal 5.5%)
 - transverse momentum resolution 10% at 10 GeV/c, no calibration (PPR 5%)

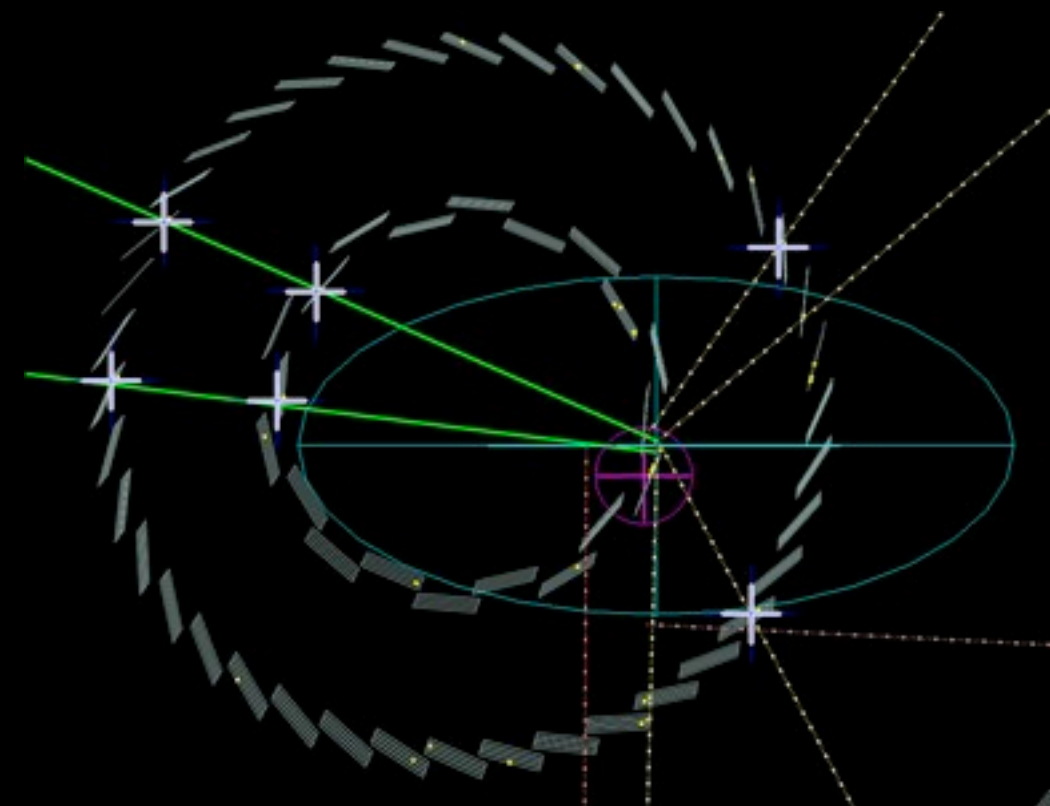
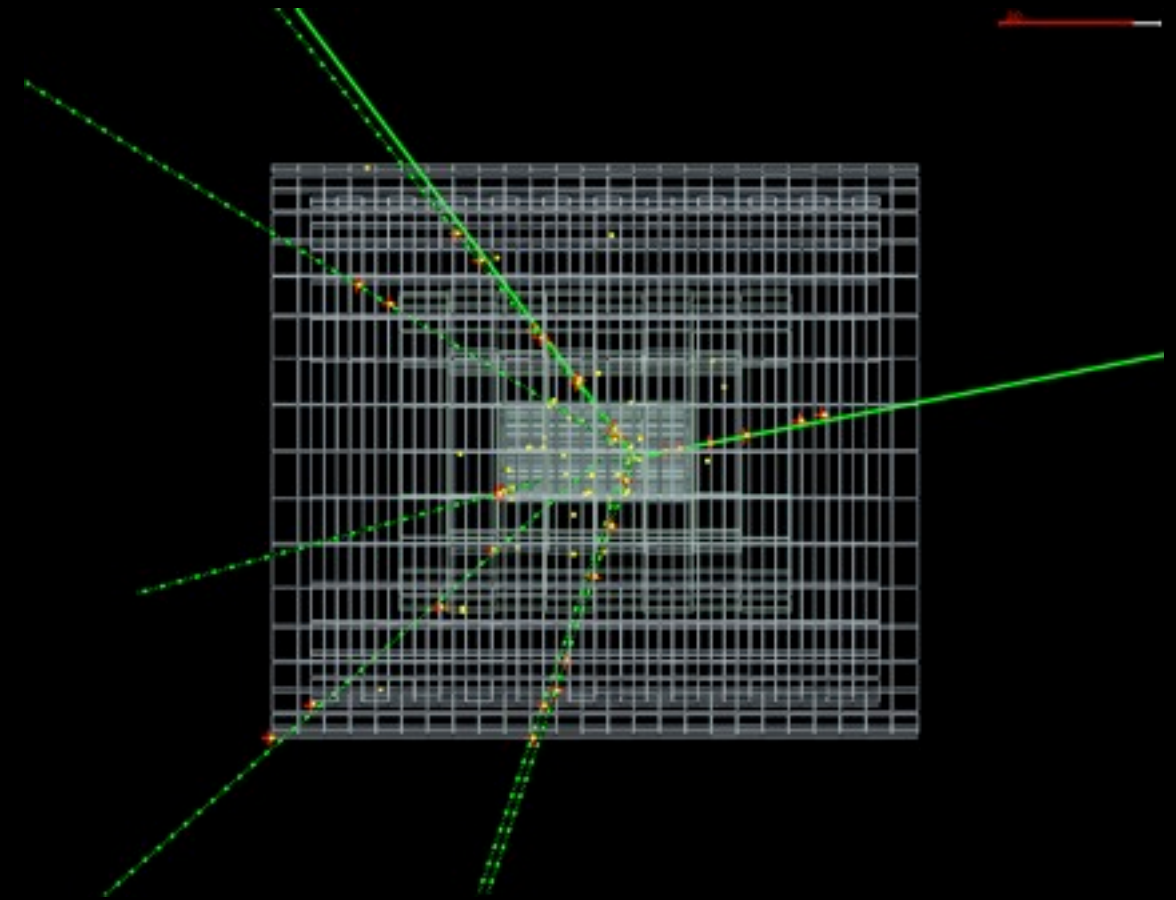
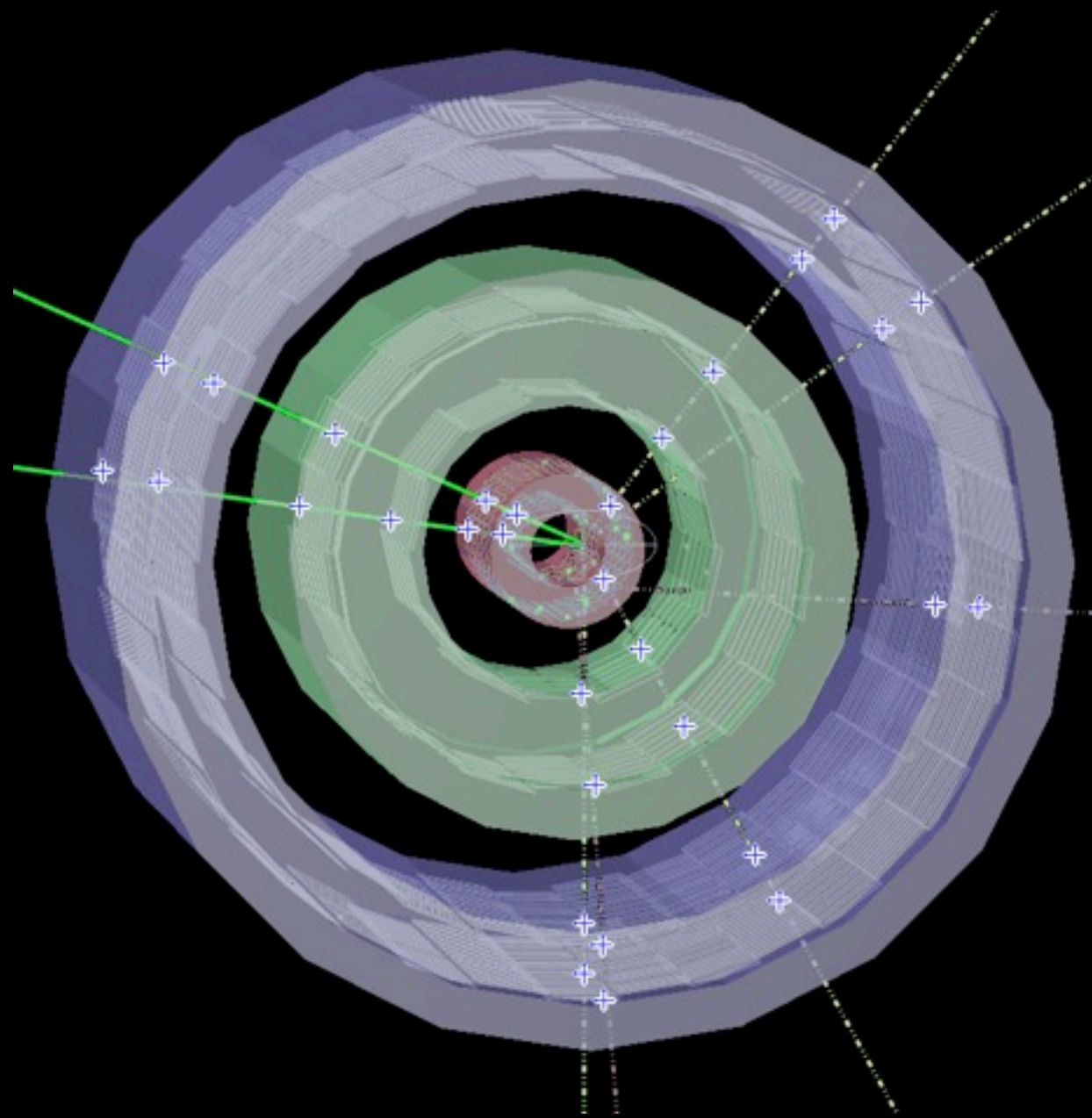


Injection Test

- Events in ALICE
 - Tl2 dump June 15
 - Injection test August
 - Circulating beam September
- Observed very high particle densities
 - 10s - 1000s of particles per cm²
 - sensitive detectors turned off
 - SPD and V0 always on
 - SDD, SSD, T0 and FMD sometimes



First Interaction I2-9



Summary

- heavy-ion experiments provide us access to the properties of matter at extreme temperatures and densities
- the LHC is the next chapter and a step above and beyond existing facilities
- ALICE is ready

Backup

- ☒ Event 73
 - ☒ Primary Vertex
 - ☐ Primary Vertex SPD
 - ☒ V0 offline vertex locations
 - ☒ V0 on-the-fly vertex locations
 - ☒ ESD v0
 - ☒ ITS Clusters
 - ☒ TPC Clusters
 - ☒ ESD Tracks
 - ☒ TPC 2D
 - ☒ TPC 3D
 - ☒ 3D sector 0
 - ☒ 3D sector 1
 - ☒ 3D sector 2
 - ☒ 3D sector 3

Style

Name Event 73::AliEveEventManager

TEveElement

Show: ☒ Self ☒ Children

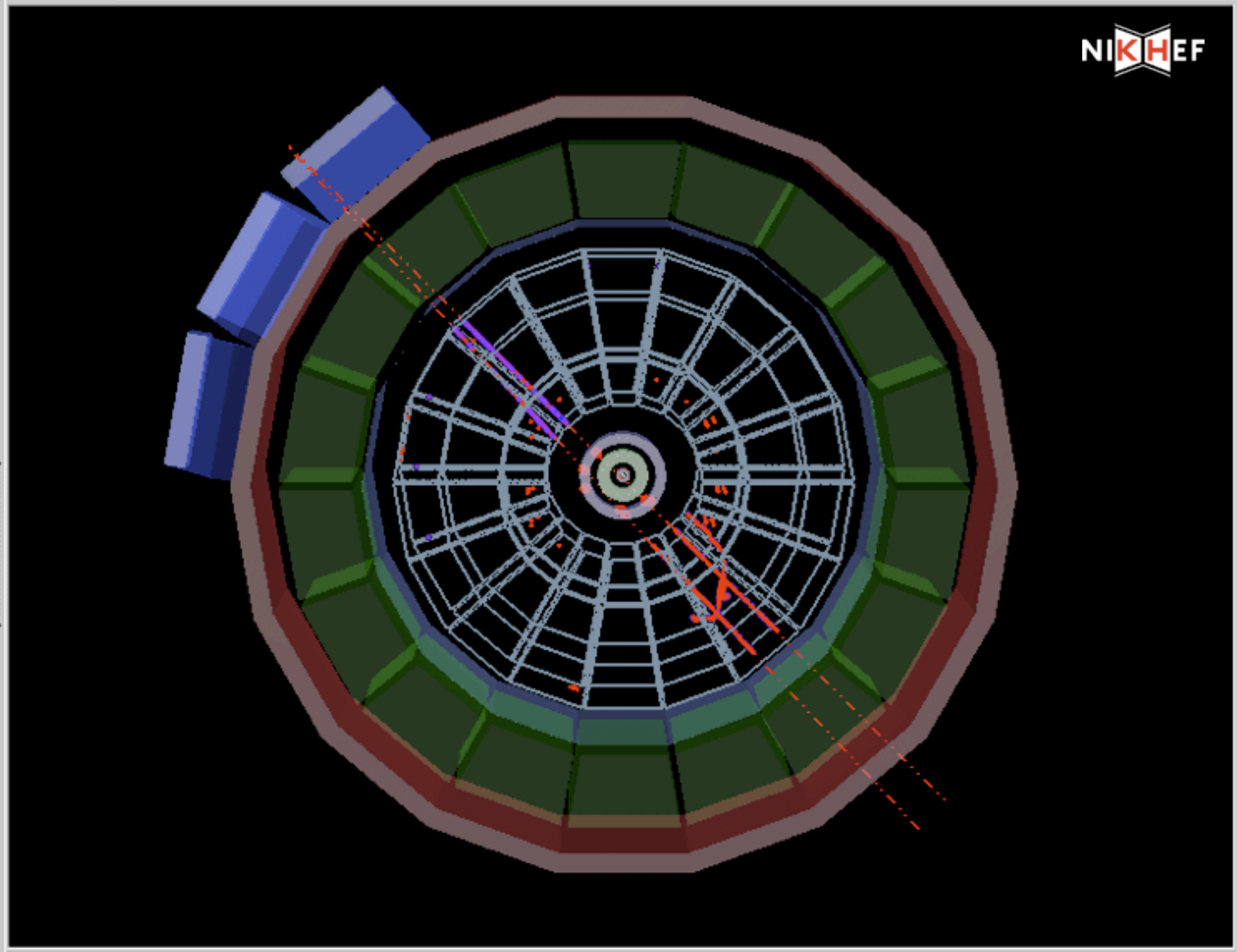
AliEveEventManager

Next Event

Event Information:

Raw-data event info:
 Run#: 60705
 Event type: 7 (PHYSICS_EVENT)
 Period: 1
 Orbit: 4c635e BC: 718
 Trigger: 1
 Detectors: 8002000d (ITSSPD ITS)
 Attributes:bf-0-30
 Timestamp: 48dabbac

ESD event info:
 Run#: 60705
 Active trigger classes: D0SC0
 Event type: 7 (PHYSICS_EVENT)
 Period: 1
 Orbit: 4c635e BC: 718
 Trigger: 1 (D0SC0)

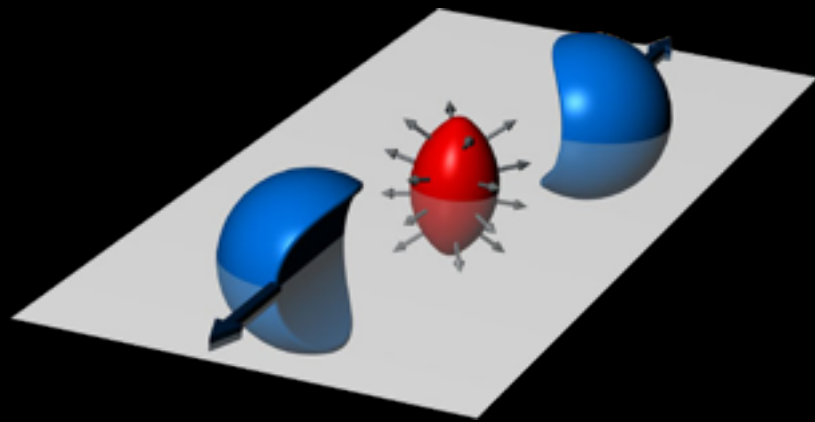


Command EventCtrl

First Prev 73 / 911 Next Last || Refresh || ☐ Autoload Time: 5 || TRG select:

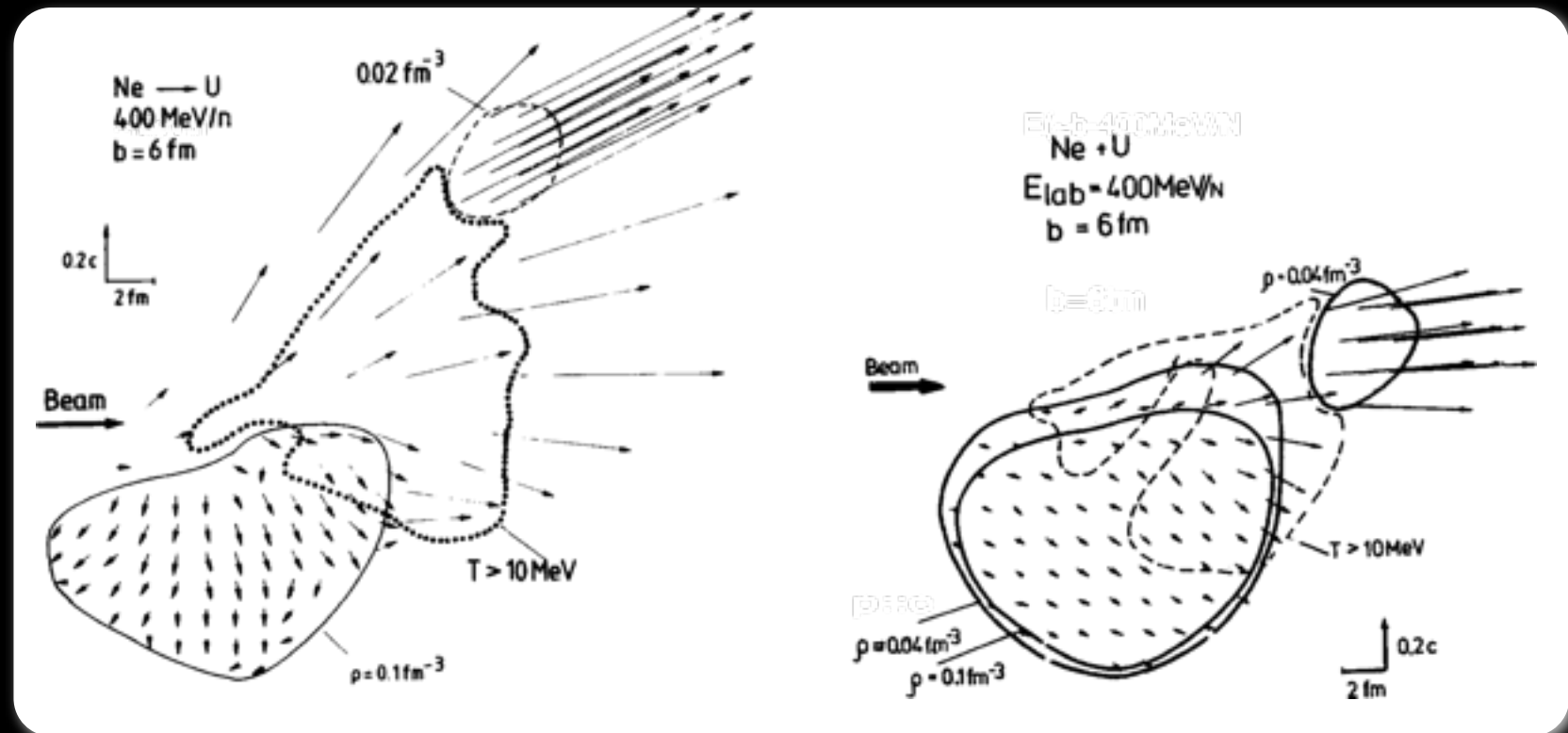
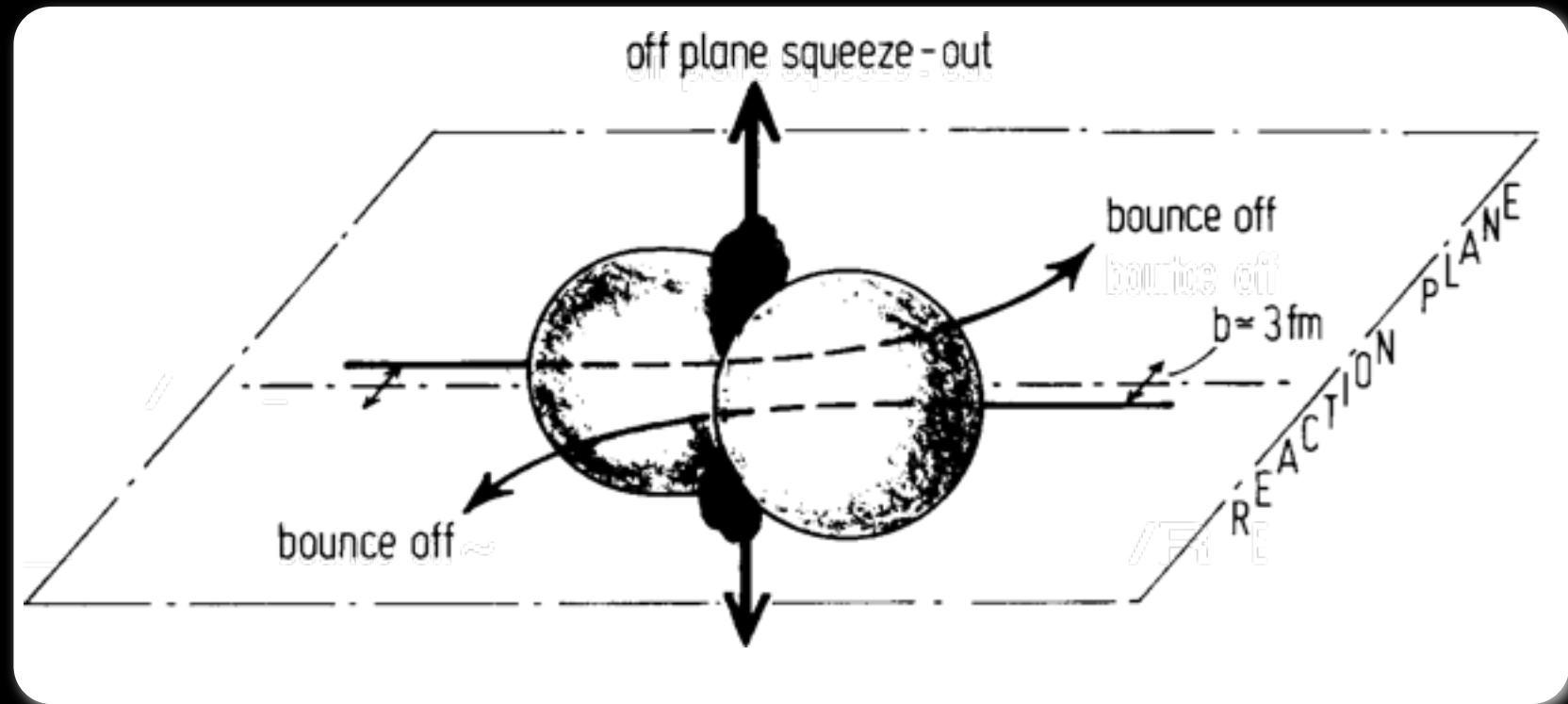
RAW event info: Run#: 60705 Event type: 7 (PHYSICS_EVENT) Period: 1 Orbit: 4c635e BC: 718
 Trigger: 1
 Detectors: 8002000d (ITSSPD ITSSSD TPC TRG)

Collective Motion



H. Stöcker and W. Greiner,
Phys. Rep. 137 (1986) 277

quickly
recognized as
important probe

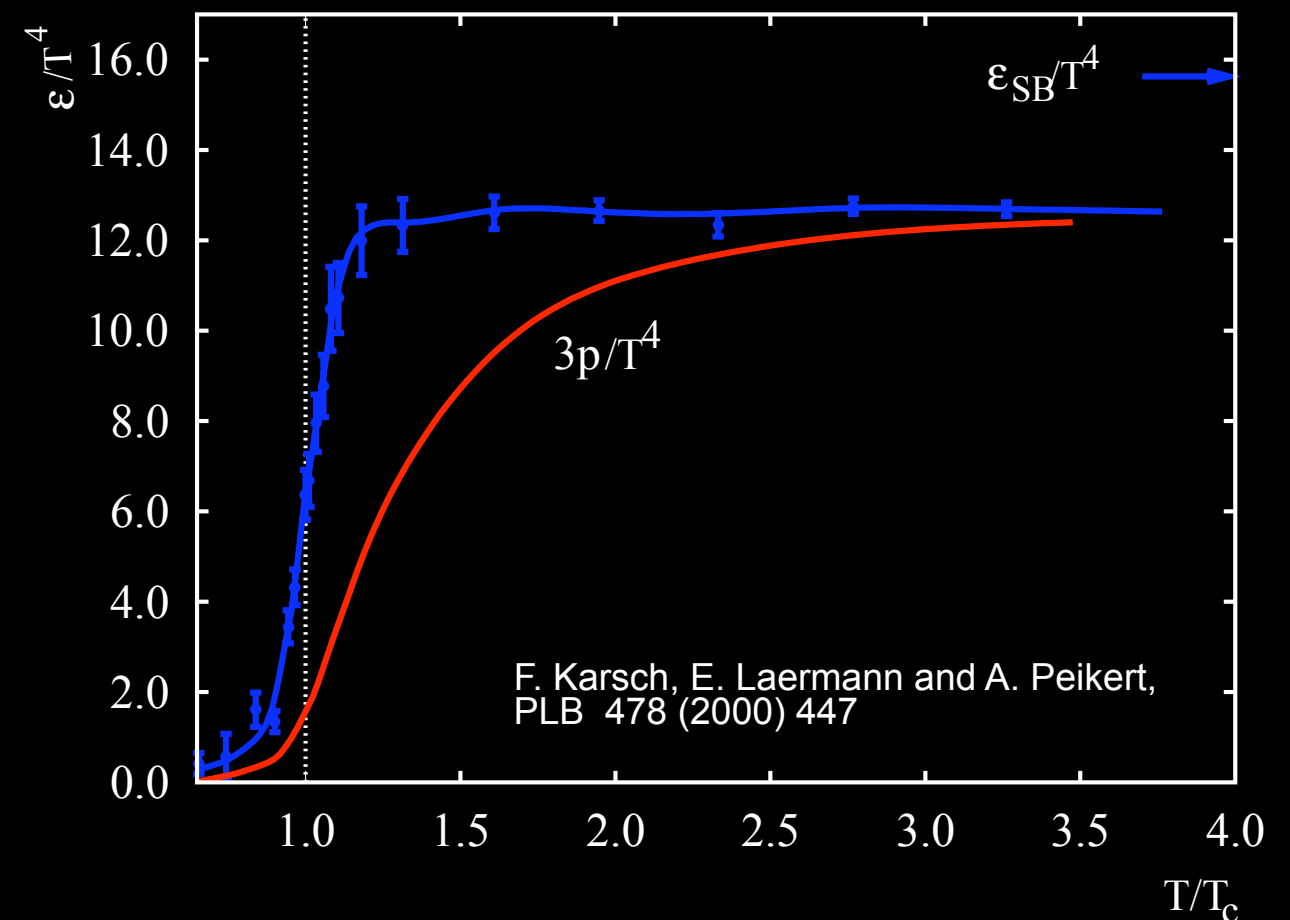


QCD and QGP

It would be interesting to explore new phenomena by distributing high energy or high nuclear density over a relatively large volume

T.D. Lee (1978)

Lattice QCD predicts a sharp rise in the number of degrees of freedom near T_c



The LHC

Running parameters

Collision system	$\sqrt{s_{NN}}$ (TeV)	\mathcal{L}_0 (cm ⁻² s ⁻¹)	$\langle \mathcal{L} \rangle / \mathcal{L}_0$ (%)	Run time (s/year)
pp	14.0	10^{34} *		10^7
PbPb	5.5	10^{27}	50	10^6 **
pPb	8.8	10^{29}		10^6
ArAr	6.3	10^{29}	65	10^6

* \mathcal{L}_{max} (ALICE) = 10^{31}

** \mathcal{L}_{int} (ALICE) ~ 0.5 nb⁻¹/year

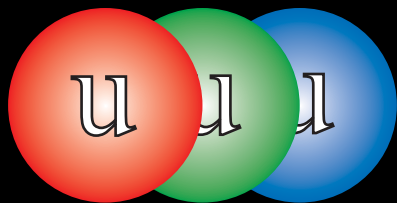
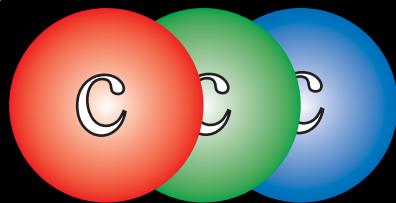
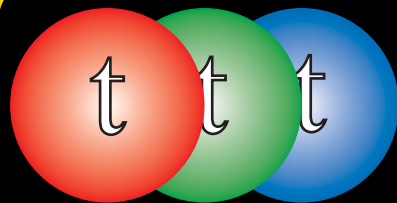
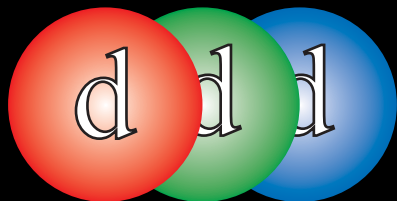
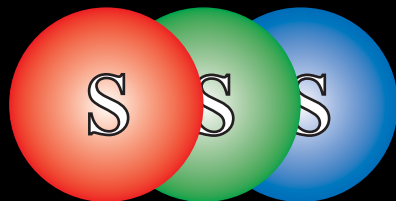
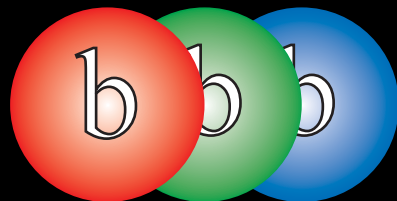



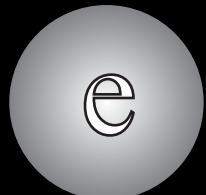
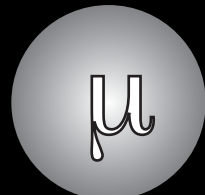
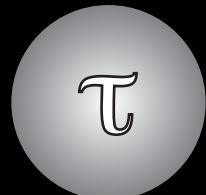
The Standard Model

(matter)

1st generation 2nd generation 3rd generation

quarks

leptons

	~3		1250		174300	2/3
	~6		120		4200	-1/3
	~0		~0		~0	0
	0.511		106		1770	-1
mass		mass		mass		charge

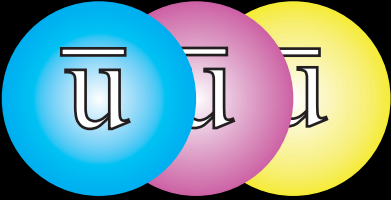
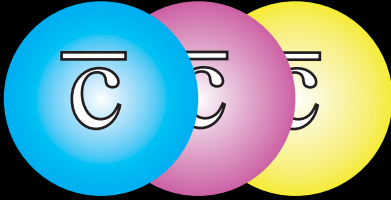
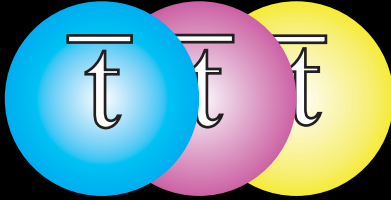
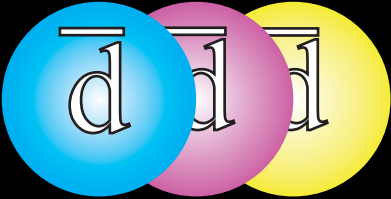
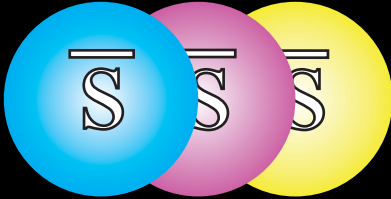
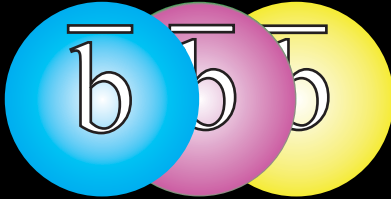



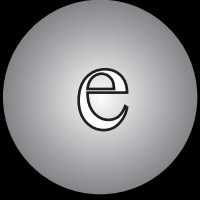
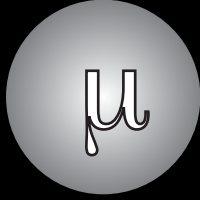
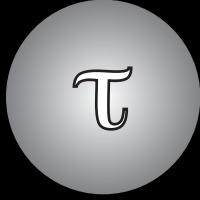
The Standard Model

(anti-matter)

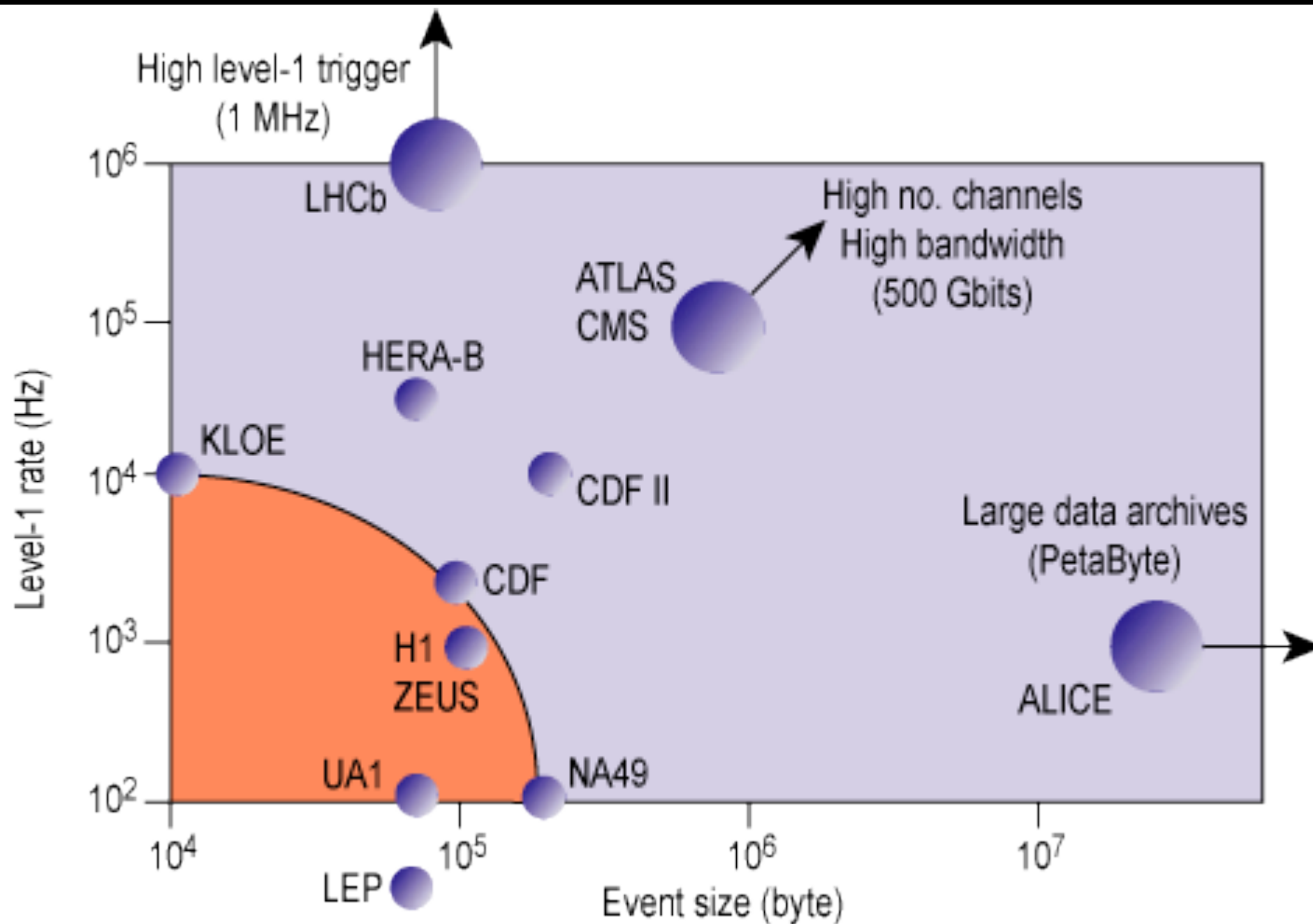
1st generation 2nd generation 3rd generation

anti-
quarks

anti-
leptons

anti-quarks		~3		1250		174300	-2/3
		~6		120		4200	1/3
anti-leptons		~0		~0		~0	0
		0.511		106		1770	+1
	mass		mass		mass		charge

Alice DAQ



Proton-proton physics with ALICE

- ❑ The ALICE detector works even better for pp collisions, because of the low occupancy (10^{-4} to 10^{-3}), even if there is a significant number of events overlapping.
- ❑ The first physics with ALICE will be proton-proton collisions, which correspond to a major part of the ALICE programme for **several reasons**:
 - to provide **“reference” data** to understand heavy ion collisions. In a new energy domain, each signal in HI has to be compared to pp;
 - For **genuine proton-proton physics** whenever ALICE is unique or competitive; note that ALICE can reach rather “high” p_T , up to ~ 100 GeV/c, ensuring overlap with other LHC experiments.
 - The **possibility of taking proton data at several center of mass energies (0.9 TeV, 2.4 TeV, perhaps 5.5 TeV, and 14 TeV)**, will provide ALICE with the possibility to understand the evolution of many of the properties of pp collisions as a function of the center of mass energy, and also to add to the measurements from previous experiments using proton-antiprotons.